

TWISTED DEFORMATIONS VS. COCYCLE DEFORMATIONS FOR QUANTUM GROUPS

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ABSTRACT. In this paper we study two deformation procedures for quantum groups: deformations by twists, that we call “*comultiplication twisting*”, as they modify the coalgebra structure, while keeping the algebra one — and deformations by 2-cocycle, that we call “*multiplication twisting*”, as they deform the algebra structure, but save the coalgebra one.

We deal with quantum universal enveloping algebras, in short QUEA’s, for which we accordingly consider those arising from twisted deformations (in short TwQUEA’s) and those arising from 2-cocycle deformations, usually called multiparameter QUEA’s (in short MpQUEA’s). Up to technicalities, we show that the two deformation methods are equivalent, in that they eventually provide isomorphic outputs, which are deformations (of either kinds) of the “canonical”, well-known one-parameter QUEA by Jimbo and Lusztig. It follows that the two notions of TwQUEA’s and of MpQUEA’s — which, in Hopf algebra theoretical terms are naturally dual to each other — actually coincide; thus, that there exists in fact only one type of “pluriparametric deformation” for QUEA’s. In particular, the link between the realization of any such QUEA as a MpQUEA and that as a TwQUEA is just a (very simple, and rather explicit) change of presentation.

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1. INTRODUCTION

Roughly speaking, quantum groups — in the form of quantized universal enveloping algebras — are Hopf algebra deformations of the universal enveloping algebra $U(\mathfrak{g})$ of some Lie algebra \mathfrak{g} . From this deformation, \mathfrak{g} itself inherits (as “semiclassical limit” of the deformed coproduct) a Lie cobracket that makes it into a Lie bialgebra — the infinitesimal counterpart of a Poisson group whose tangent Lie algebra is \mathfrak{g} .

When \mathfrak{g} is a complex simple Lie algebra, a quantum group in this sense, depending on a single parameter, was introduced by Drinfeld [Dr] as a formal series deformation $U_{\hbar}(\mathfrak{g})$ defined over a ring of formal power series (in the formal parameter \hbar) and by Jimbo and Lusztig (see [Ji], [Lu]) as a deformation $U_q(\mathfrak{g})$ defined over a ring of rational series (in the formal parameter q). Indeed, Jimbo’s $U_q(\mathfrak{g})$ is actually a “polynomial version” of Drinfeld’s $U_{\hbar}(\mathfrak{g})$.

Later on, several authors (cf. [BGH], [BW1, BW2], [CM], [CV1], [Hay], [HLT], [HPR], [Ko], [KT], [Ma], [OY], [Re], [Su], [Ta], to name a few) introduced many types of deformations of $U(\mathfrak{g})$ depending on several parameters, usually referred to as “multiparameter quantum groups”. In turn, these richer deformations induce as semiclassical limits corresponding “multiparameter” bialgebra structures on \mathfrak{g} . The construction of these multiparameter deformations applies a general procedure, always available for Hopf algebras, following two patterns that we recall hereafter.

Let H be any Hopf algebra (in some braided tensor category). Among all possible deformations of the Hopf structure of H , we look at those in which only one of either the product or the coproduct is actually modified, while the other one is kept fixed. The general deformation will then be, somehow, an intermediate case between two such extremes. On the one hand, a *twist deformation* of H is a (new) Hopf algebra structure on H where the multiplicative structure is unchanged, whereas a new coproduct is defined by $\Delta^{\mathcal{F}}(x) := \mathcal{F} \Delta(x) \mathcal{F}^{-1}$ for $x \in H$: here \mathcal{F} is an invertible element in $H^{\otimes 2}$ satisfying suitable axioms, called a “twist” for H . On the other hand, a *2-cocycle deformation* of H is one where the coproduct is unchanged, while a new product is defined via a formula which only depends on the old product and on a 2-cocycle σ of H (as an algebra): again, this procedure can be read as a suitable “conjugation” of the old product map by the 2-cocycle.

Inasmuch as a meaningful notion of “duality” applies to the Hopf algebras one is dealing with, these two constructions of deformations are *dual to each other*, directly by definition. In detail, if H^* is a Hopf algebra dual to H with respect to a non-degenerate (skew) Hopf pairing, e.g. H and $H^* := H^\circ$ (i.e., Sweedler’s *restricted dual*), then the dual of the deformation by twist, resp. by 2-cocycle, of H is a deformation by 2-cocycle, resp. by twist, of H^* ; moreover, the 2-cocycle, resp. the twist, on H^* is uniquely determined by the twist, resp. the 2-cocycle, on H . In order to stress this duality between the two types of deformation procedures that we are dealing with, as well as the fact that both are in fact “conjugations” of some sort, we adopt the terminology “*comultiplication twisting*” and “*multiplication twisting*”, instead of “*deformation by twist*” and of “*deformation by 2-cocycle*”, respectively.

It so happens that the large majority of multiparameter quantizations of $U(\mathfrak{g})$ considered in literature actually occur as either comultiplication twistings or multiplication twistings of a one-parameter quantization of Drinfeld’s type or Jimbo-Lusztig’s type. Indeed, in both cases the twists and the 2-cocycles taken into account are of

special type, namely “toral” ones, in that (roughly speaking) they are defined only in terms of the (quantum) toral part of the one-parameter deformation of $U(\mathfrak{g})$.

Technically speaking, Drinfeld’s $U_h(\mathfrak{g})$ is better suited for comultiplication twistings, while Jimbo-Lusztig’s $U_q(\mathfrak{g})$ is typically used for multiplication twistings (see [Re], [Ma], [Su], [HPR], [HLT], [CV1], [Ta]). As we aim to compare both kinds of twistings, we focus on polynomial one-parameter quantum groups $U_q(\mathfrak{g})$, and we adapt the very notion of “twist deformation”, or “comultiplication twisting”, to them. Then we consider both comultiplication twistings and multiplication twistings (of “toral type”, in both cases) of $U_q(\mathfrak{g})$ — thus getting “twisted quantized universal enveloping algebras (=TwQUEA’s)” and “multiparameter quantized universal enveloping algebras (=MpQUEA’s)”, respectively — and compare them. Moreover, by natural reasons we restrict ourselves to twists and cocycles that are defined by a *rational datum*, i.e., a matrix with rational entries.

As a first result, we describe the link *twist* \longleftrightarrow *2-cocycle* under duality. Namely, quantum Borel (sub)groups $U_q(\mathfrak{b}_\pm)$ of opposite signs are in Hopf duality (in a proper sense): then we prove that any twisting on one side — of either comultiplication or multiplication — and the *dual* one on the other side — of either multiplication or comultiplication, respectively — are described by *the same rational datum*. Indeed, we provide an explicit bijection between the sets of toral twists and toral 2-cocycles.

As a second, more striking result (the core of our paper, indeed), we find that, in short, *twisted quantum groups and multiparameter quantum groups coincide*: namely, any TwQUEA can be realized as a MpQUEA, and viceversa. Even more precisely, the twist and the 2-cocycle involved in either realization are described by *the same (rational) datum*. This result is, in a sense, a side effect of the “autoduality” of quantum groups (in particular Borel ones). The proof is constructive, and quite explicit: indeed, switching from the realization as TwQUEA to that as MpQUEA and viceversa is a sheer *change of presentation*. We can shortly sketch the underlying motivation: any “standard” (=undeformed) quantum group is pointed (as a Hopf algebra); then any TwQUEA of “toral type” is pointed as well, and it is generated by the quantum torus and $(1, g)$ -skew primitive elements: these new “homogeneous” generators yield a new presentation, which realizes the TwQUEA as a MpQUEA.

The direct consequence of this result is that (roughly speaking, and within the borders of our restrictions) there exists only *one* type of multiparameter quantization of $U(\mathfrak{g})$, and consequently only one type of corresponding multiparameter Lie bialgebra structure on \mathfrak{g} arising as semiclassical limits, as in [GG1].

All the elements that lead us to the above mentioned results for TwQUEA’s and MpQUEA’s are also available for Hopf algebras that (like Borel quantum subgroups) are bosonizations of Nichols algebras of diagonal type; thus, we can replicate our work in that context too. In another direction, we extend further on this analysis in the framework of multiparametric *formal* QUEA’s, à la Drinfeld — cf. [GG2].

We finish with a few words on the structure of the paper.

In Section 2 we collect the material on Hopf algebras and their deformations that will be later applied to quantum groups. Section 3 is devoted to introduce quantum groups (both in Drinfeld’s version and in Jimbo-Lusztig’s one) and their comultiplication twistings (of rational, toral type), i.e., the TwQUEA’s: the part on *Drinfeld’s* quantum groups could be dropped, yet we present it to explain the deep-rooting motivations of our work. In Section 4, instead, we present the multiplication

twistings (of rational, toral type) of Jimbo-Lusztig’s quantum groups, hence the MpQUEA’s. Finally, in Section 5 we compare TwQUEA’s and MpQUEA’s, proving that — in a proper sense, under some finiteness assumption — they actually coincide.

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2. PRELIMINARIES

In this section we fix the basic material on Hopf algebras and combinatorial data that we shall need later on. In particular, $\mathbb{N} = \{0, 1, \dots\}$ and $\mathbb{N}_+ := \mathbb{N} \setminus \{0\}$.

2.1. The combinatorial tool box.

The definition of our multiparameter quantum groups requires a full lot of related material that we now present. First of all, \mathbb{k} will be a field of characteristic zero.

2.1.1. Root data and Lie algebras. Hereafter we fix $n \in \mathbb{N}_+$ and $I := \{1, \dots, n\}$. Let $A := (a_{ij})_{i,j \in I}$ be a generalized, symmetrisable Cartan matrix; then there exists a unique diagonal matrix $D := (d_i \delta_{ij})_{i,j \in I}$ with positive integral, pairwise coprime entries such that DA is symmetric. Let $\mathfrak{g} = \mathfrak{g}_A$ be the Kac-Moody algebra over \mathbb{C} associated with A (cf. [Ka]); we consider a split integral \mathbb{Z} -form of \mathfrak{g} , and its scalar extension \mathfrak{g}_R from \mathbb{Z} to any ring R : when this ring is \mathbb{k} , by abuse of notation *the resulting Lie algebra over \mathbb{k} will be denoted by \mathfrak{g} again*.

Let Φ be the root system of \mathfrak{g} , with $\Pi = \{\alpha_i \mid i \in I\}$ as a set of simple roots, $Q = \bigoplus_{i \in I} \mathbb{Z} \alpha_i$ the associated root lattice, Φ^+ the set of positive roots with respect to Π , $Q^+ = \bigoplus_{i \in I} \mathbb{N} \alpha_i$ the positive root (semi)lattice.

Fix a Cartan subalgebra \mathfrak{h} of \mathfrak{g} , whose associated set of roots identifies with Φ (so $\mathbb{k}Q \subseteq \mathfrak{h}^*$); then for all $\alpha \in \Phi$ we call \mathfrak{g}_α the corresponding root space. Now set $\mathfrak{h}' := \mathfrak{g}' \cap \mathfrak{h}$ where $\mathfrak{g}' := [\mathfrak{g}, \mathfrak{g}]$ is the derived Lie subalgebra of \mathfrak{g} : then $(\mathfrak{h}')^* = \mathbb{k}Q \subseteq \mathfrak{h}^*$. We fix a \mathbb{k} -basis $\Pi^\vee := \{h_i := \alpha_i^\vee\}_{i \in I}$ of \mathfrak{h}' so that $(\mathfrak{h}, \Pi, \Pi^\vee)$ is a *realization* of A , as in [Ka, Chapter 1]; in particular, $\alpha_i(\mathfrak{h}_j) = a_{ji}$ for all $i, j \in I$.

Let \mathfrak{h}'' be any vector space complement of \mathfrak{h}' inside \mathfrak{h} . Then there exists a unique symmetric \mathbb{k} -bilinear pairing on \mathfrak{h} , denoted $(\ , \)$, such that $(h_i, h_j) = a_{ji} d_i^{-1}$, $(h_i, h_2'') = \alpha_i(h_2'') d_i^{-1}$ and $(h_1'', h_2'') = 0$, for all $i, j \in I$, $h_1'', h_2'' \in \mathfrak{h}''$; in addition, this pairing is invariant and non-degenerate (cf. [Ka, Lemma 2.1]). By non-degeneracy, this pairing induces a \mathbb{k} -linear isomorphism $t : \mathfrak{h}^* \xrightarrow{\cong} \mathfrak{h}$, with $t(\alpha_i) = d_i \alpha_i^\vee = d_i h_i$ for all $i \in I$, and this in turn defines a similar pairing on \mathfrak{h}^* — again denoted $(\ , \)$ — via pull-back, namely $(t^{-1}(h_1), t^{-1}(h_2)) := (h_1, h_2)$; in particular, on simple roots this gives $(\alpha_i, \alpha_j) := d_i a_{ij}$ for all $i, j \in I$. In fact, this pairing on \mathfrak{h}^* does

restrict to a (symmetric, \mathbb{Z} -valued, \mathbb{Z} -bilinear) pairing on Q ; note that, in terms of the latter pairing on Q , one has $d_i = (\alpha_i, \alpha_i)/2$ and $a_{ij} = \frac{2(\alpha_i, \alpha_j)}{(\alpha_i, \alpha_i)}$ for all $i, j \in I$. We also notice that $t : \mathfrak{h}^* \xrightarrow{\cong} \mathfrak{h}$ restricts to another isomorphism $t' : (\mathfrak{h}')^* \xrightarrow{\cong} \mathfrak{h}'$ for which we use notation $t_\alpha := t'(\alpha) = t(\alpha)$.

When A is of *finite type* — which is equivalent to saying that $\mathfrak{h}' = \mathfrak{h}$ — we denote by P its associated *weight lattice*, with basis $\{\omega_i\}_{i \in I}$ dual to $\{\alpha_j\}_{j \in I}$, so that $\omega_i(\alpha_j) = \delta_{ij}$ for $i, j \in I$. If we identify \mathfrak{h}^* with \mathfrak{h} via the isomorphism $t : \mathfrak{h}^* \xrightarrow{\cong} \mathfrak{h}$ as above, the root lattice Q identifies with a suitable sublattice of P : then there is also an identification $\mathbb{Q}P = \mathbb{Q}Q$ — where hereafter we use such notation as $\mathbb{Q}Q := \mathbb{Q} \otimes_{\mathbb{Z}} Q$, etc. — and we have such identities as $\alpha_i = \sum_{j \in I} a_{ji} \omega_j$ and $(\omega_i, \alpha_j) := d_i \delta_{ij}$ for all $i, j \in I$.

According to our choice of positive and negative roots, let \mathfrak{b}_+ , resp. \mathfrak{b}_- , be the Borel subalgebra in \mathfrak{g} containing \mathfrak{h} and all positive, resp. negative, root spaces. There is a canonical, non-degenerate pairing between \mathfrak{b}_+ and \mathfrak{b}_- , using which one can construct a *Manin double* $\mathfrak{g}_D = \mathfrak{b}_+ \oplus \mathfrak{b}_-$, that is automatically endowed with a structure of Lie bialgebra — roughly, \mathfrak{g}_D is like \mathfrak{g} but with *two copies of \mathfrak{h}* inside it (cf. [CP], §1.4), namely $\mathfrak{h}_+ := \mathfrak{h} \oplus 0$ inside \mathfrak{b}_+ and $\mathfrak{h}_- := 0 \oplus \mathfrak{h}$ inside \mathfrak{b}_- ; accordingly, we set also $\mathfrak{h}'_+ := \mathfrak{h}' \oplus 0$ and $\mathfrak{h}'_- := 0 \oplus \mathfrak{h}'$. By construction both \mathfrak{b}_+ and \mathfrak{b}_- lies in \mathfrak{g}_D as Lie sub-bialgebras. Moreover, there exists a Lie bialgebra epimorphism $\pi_{\mathfrak{g}_D} : \mathfrak{g}_D \longrightarrow \mathfrak{g}$ which maps the copy of \mathfrak{b}_\pm inside \mathfrak{g}_D identically onto its copy in \mathfrak{g} .

For later use we fix generators e_i, h_i, f_i ($i \in I$) in \mathfrak{g} as in the usual Serre's presentation of \mathfrak{g} . Moreover, for the corresponding elements inside $\mathfrak{g}_D = \mathfrak{b}_+ \oplus \mathfrak{b}_-$ we adopt notation $e_i := (e_i, 0)$, $h_i^+ := (h_i, 0)$, $h_i^- := (0, h_i)$ and $f_i := (0, f_i)$, for all $i \in I$. Notice that we have by construction

$$e_i \in \mathfrak{g}_{+\alpha_i} \quad , \quad h_i = d_i^{-1} t_{\alpha_i} \in \mathfrak{h} \quad , \quad f_i \in \mathfrak{g}_{-\alpha_i} \quad \forall i \in I \quad (2.1)$$

2.1.2. Root twisting. Applying the twisting procedure to quantized universal enveloping algebras, we shall eventually be lead to consider an operation of “root twisting”, in some sort, that we formalize hereafter.

Let \hbar be a formal variable, and $\mathbb{k}((\hbar))$ the corresponding field of Laurent formal series with coefficients in \mathbb{k} ; for simplicity, we eventually will take a field \mathbb{F} of characteristic zero containing $\mathbb{k}((\hbar))$. Fix a subring \mathcal{R} of $\mathbb{k}((\hbar))$ containing $\mathbb{Q}[[\hbar]]$, let $\mathcal{R}Q$ be the scalar extension of Q by \mathcal{R} , and fix an $(n \times n)$ -matrix $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathcal{R})$. We define the endomorphisms $\psi_\pm : \mathcal{R}Q \longrightarrow \mathcal{R}Q$ given by

$$\psi_+(\alpha_\ell) = \zeta_\ell^+ := \sum_{i,j \in I} \psi_{ij} a_{j\ell} d_i^{-1} \alpha_i \quad , \quad \psi_-(\alpha_\ell) = \zeta_\ell^- := \sum_{i,j \in I} \psi_{ji} a_{j\ell} d_i^{-1} \alpha_i \quad \forall \ell \in I \quad (2.2)$$

that in matrix notation reads

$$(\psi_+(\alpha_\ell) = \zeta_\ell^+)_{\ell \in I} := A^T \Psi^T D^{-1} (\alpha_k)_{k \in I} \quad , \quad (\psi_-(\alpha_\ell) = \zeta_\ell^-)_{\ell \in I} := A^T \Psi D^{-1} (\alpha_k)_{k \in I} \quad (2.3)$$

where $(\alpha_k)_{k \in I} = (\alpha_1, \dots, \alpha_n)^T$ is thought of as a column vector, and likewise for $(\psi_\pm(\alpha_\ell))_{\ell \in I} = (\zeta_\ell^\pm)_{\ell \in I}$. Now, borrowing notation from §2.1.1 we fix the \mathcal{R} -integral form $\mathfrak{h}'_{\mathcal{R}}$ of \mathfrak{h}' in $\mathfrak{g}_{\mathcal{R}}$ spanned by the simple coroots h_i ($i \in I$), and the corresponding isomorphism $t' : (\mathfrak{h}'_{\mathcal{R}})^* \xrightarrow{\cong} \mathfrak{h}'_{\mathcal{R}}$ ($\alpha \mapsto t_\alpha$) — this does make sense, because the original isomorphism $t' : (\mathfrak{h}')^* \xrightarrow{\cong} \mathfrak{h}'$ in §2.1.1 is actually well-defined over \mathbb{Q} .

Then we define endomorphisms $\psi_{\pm}^{\mathfrak{h}'}$ of $\mathfrak{h}'_{\mathcal{R}}$ as $\psi_{\pm}^{\mathfrak{h}'} := t' \circ \psi_{\pm} \circ (t')^{-1}$; if we set $T_{\ell} := t_{\alpha_{\ell}}$ and $H_{\ell} := d_{\ell}^{-1} T_{\ell}$ for $\ell \in I$, these are obviously described by

$$\psi_{+}^{\mathfrak{h}'}(T_{\ell}) := \sum_{i,j \in I} \psi_{ij} a_{j\ell} d_i^{-1} T_i := \sum_{i,j \in I} \psi_{ij} a_{j\ell} H_i \quad \forall \ell \in I$$

$$\psi_{-}^{\mathfrak{h}'}(T_{\ell}) := \sum_{i,j \in I} \psi_{ji} a_{j\ell} d_i^{-1} T_i := \sum_{i,j \in I} \psi_{ji} a_{j\ell} H_i \quad \forall \ell \in I$$

that in matrix notation reads

$$\begin{aligned} (\psi_{+}^{\mathfrak{h}'}(T_{\ell}))_{\ell \in I} &:= A^T \Psi^T D^{-1} (T_k)_{k \in I} = A^T \Psi^T (H_k)_{k \in I} \\ (\psi_{-}^{\mathfrak{h}'}(T_{\ell}))_{\ell \in I} &:= A^T \Psi D^{-1} (T_k)_{k \in I} = A^T \Psi (H_k)_{k \in I} \end{aligned}$$

Note that, by definition, we have $\psi_{+} = \psi_{-}$, or equivalently $\psi_{+}^{\mathfrak{h}'} = \psi_{-}^{\mathfrak{h}'}$, if and only if $\Psi^T = \Psi$, i.e., Ψ is symmetric.

Finally we introduce the following elements of $(\mathfrak{h}'_{\mathcal{R}})_{+} \oplus (\mathfrak{h}'_{\mathcal{R}})_{-}$, for all $i \in I$:

$$T_{i,+}^{\Psi} := (\text{id}_{\mathfrak{h}'_{+}} + \psi_{+}^{\mathfrak{h}'}) (T_i^{+}) - \psi_{-}^{\mathfrak{h}'} (T_i^{-}) \quad , \quad T_{i,-}^{\Psi} := (\text{id}_{\mathfrak{h}'_{-}} + \psi_{-}^{\mathfrak{h}'}) (T_i^{-}) - \psi_{+}^{\mathfrak{h}'} (T_i^{+})$$

The following two results will be of use later on:

Lemma 2.1.3.

(a) The maps $\pm(\psi_{+} - \psi_{-}) : (\mathfrak{h}'_{\mathcal{R}})^{*} \longrightarrow (\mathfrak{h}'_{\mathcal{R}})^{*}$ are antisymmetric with respect to the (symmetric) bilinear product $(\ , \)$ on $(\mathfrak{h}'_{\mathcal{R}})^{*}$;

(b) The maps $\pm(\psi_{+}^{\mathfrak{h}'} - \psi_{-}^{\mathfrak{h}'}) : \mathfrak{h}'_{\mathcal{R}} \longrightarrow \mathfrak{h}'_{\mathcal{R}}$ are antisymmetric with respect to the (symmetric) bilinear product $(\ , \)$ on $\mathfrak{h}'_{\mathcal{R}}$.

Proof. Both claims in the statement follow by sheer computation. Namely, for the map $(\psi_{+} - \psi_{-})$ this gives, for all $h, k \in I$,

$$\begin{aligned} & \left((\psi_{+} - \psi_{-})(\alpha_h), \alpha_k \right) + \left(\alpha_h, (\psi_{+} - \psi_{-})(\alpha_k) \right) = \\ & = \sum_{i,j \in I} \left((\psi_{ij} a_{jh} d_i^{-1} \alpha_i, \alpha_k) - (\psi_{ji} a_{jh} d_i^{-1} \alpha_i, \alpha_k) \right) + \\ & \quad + \sum_{i,j \in I} \left((\alpha_h, \psi_{ij} a_{jk} d_i^{-1} \alpha_i) - (\alpha_h, \psi_{ji} a_{jk} d_i^{-1} \alpha_i) \right) = \\ & = \sum_{i,j \in I} \left(\psi_{ij} a_{jh} a_{ik} - \psi_{ji} a_{jh} a_{ik} \right) + \sum_{i,j \in I} \left(\psi_{ij} a_{jk} a_{ih} - \psi_{ji} a_{jk} a_{ih} \right) = \\ & = \left(A^T \Psi A \right)_k^h - \left(A^T \Psi^T A \right)_k^h + \left(A^T \Psi^T A \right)_k^h - \left(A^T \Psi A \right)_k^h = 0 \end{aligned}$$

where M_k^h always denotes the (k, h) -entry of a matrix M . So $\pm(\psi_{+} - \psi_{-})$ is antisymmetric. The proof for $\pm(\psi_{+}^{\mathfrak{h}'} - \psi_{-}^{\mathfrak{h}'})$ is analogous. \square

Lemma 2.1.4. Assume that the Cartan matrix A is of finite type.

(a) The maps $(\text{id}_{\mathfrak{h}'} \pm (\psi_{+}^{\mathfrak{h}'} - \psi_{-}^{\mathfrak{h}'})) : \mathfrak{h}'_{\mathcal{R}} \longrightarrow \mathfrak{h}'_{\mathcal{R}}$ are bijective.

(b) The maps $(\mathfrak{h}'_{\mathcal{R}})_{\pm} \longrightarrow (\mathfrak{h}'_{\mathcal{R}})_{+} \oplus (\mathfrak{h}'_{\mathcal{R}})_{-}$ defined by $T_{\ell}^{\pm} \mapsto T_{\ell,\pm}^{\Psi}$ are injective.

Proof. Set $\phi = (\psi_{+}^{\mathfrak{h}'} - \psi_{-}^{\mathfrak{h}'})$: this is antisymmetric, hence $(\text{id}_{\mathfrak{h}'} - \phi)(\text{id}_{\mathfrak{h}'} + \phi) = (\text{id}_{\mathfrak{h}'} - \phi^2) = (\text{id}_{\mathfrak{h}'} + \phi \phi^t)$. So the claim in (a) is the same as claiming that

$(\text{id}_{\mathfrak{h}'} + \phi \phi^t)$ is non-singular, which in turn is the same as stating that -1 is not an eigenvalue of $\phi \phi^t$: the claim then follows since the latter always holds true.

The claim in (b) is a direct consequence of (a). \square

2.1.5. q -numbers. Throughout the paper we shall consider several kinds of “ q -numbers”. Let $\mathbb{Z}[q, q^{-1}]$ be the ring of Laurent polynomials with integral coefficients in the indeterminate q . For every $n \in \mathbb{N}$ we define

$$\begin{aligned} (0)_q &:= 1, & (n)_q &:= \frac{q^n - 1}{q - 1} = 1 + q + \cdots + q^{n-1} = \sum_{s=0}^{n-1} q^s & (\in \mathbb{Z}[q]) \\ (n)_q! &:= (0)_q(1)_q \cdots (n)_q = \prod_{s=0}^n (s)_q, & \binom{n}{k}_q &:= \frac{(n)_q!}{(k)_q!(n-k)_q!} & (\in \mathbb{Z}[q]) \\ [0]_q &:= 1, & [n]_q &:= \frac{q^n - q^{-n}}{q - q^{-1}} = q^{-(n-1)} + \cdots + q^{n-1} = \sum_{s=0}^{n-1} q^{2s-n+1} & (\in \mathbb{Z}[q, q^{-1}]) \\ [n]_q! &:= [0]_q[1]_q \cdots [n]_q = \prod_{s=0}^n [s]_q, & \left[\begin{matrix} n \\ k \end{matrix} \right]_q &:= \frac{[n]_q!}{[k]_q![n-k]_q!} & (\in \mathbb{Z}[q, q^{-1}]) \end{aligned}$$

In particular, we have the identities

$$(n)_{q^2} = q^{n-1}[n]_q, \quad (n)_{q^2}! = q^{\frac{n(n-1)}{2}}[n]_q!, \quad \binom{n}{k}_{q^2} = q^{k(n-k)} \left[\begin{matrix} n \\ k \end{matrix} \right]_q.$$

Furthermore, thinking of Laurent polynomials as functions on \mathbb{F}^\times , for any $q \in \mathbb{F}^\times$ we shall read every symbol above as representing the corresponding element in \mathbb{F} .

2.2. Multiparameters.

The main objects of our study will depend on “multiparameters”, i.e., suitable collections of parameters. Hereafter we introduce these gadgets and the technical results about them that we shall need later on.

2.2.1. Multiplicative multiparameters. Let \mathbb{F} be a fixed ground field, and let $I := \{1, \dots, n\}$ be as in §2.1.1 above. We fix a matrix $\mathbf{q} := (q_{ij})_{i,j \in I}$, whose entries belong to \mathbb{F}^\times and will play the role of “parameters” of our quantum groups. Then inside the lattice $\Gamma := \mathbb{Z}^n$ one has a “generalized root system” associated with the diagonal braiding given by \mathbf{q} , in which the vectors in the canonical basis of $\Gamma := \mathbb{Z}^n$ are taken as (positive) simple roots α_i ($i = 1, \dots, n$).

We shall say that the matrix \mathbf{q} is of *Cartan type* if there is a symmetrisable generalized Cartan matrix $A = (a_{ij})_{i,j \in I}$ such that

$$q_{ij} q_{ji} = q_{ii}^{a_{ij}} \quad \forall i, j \in I \quad (2.4)$$

In such a case, to avoid some irrelevant technicalities we shall assume that the Cartan matrix A is *indecomposable*: that is, A is not expressible, after any permutation of indices, as a block-diagonal matrix with more than one block (although this is *not* necessary for the theory; it only helps in having the main results look better, say). In particular, this implies that for all $i, j \in I$, there exists a sequence $i = k_1, \dots, k_\ell = j$ in I such that $a_{k_s, k_{s+1}} \neq 0$ for all $1 \leq s < \ell$. Moreover — still in the Cartan case — for later use we fix in \mathbb{F} some “square roots” of all the q_{ii} ’s, as follows. From relations (2.4) one easily finds — because the Cartan matrix A is indecomposable

— that there exists an index $j_0 \in I$ such that $q_{ii} = q_{j_0 j_0}^{e_i}$ for some $e_i \in \mathbb{N}$, for all $i \in I$. Now we assume hereafter that \mathbb{F} contains a square root of $q_{j_0 j_0}$, which we fix throughout and denote by $q_{j_0} := \sqrt{q_{j_0 j_0}}$, and also by $q := \sqrt{q_{j_0 j_0}}$ ($= q_{j_0}$). Then we set also $q_i := q_{j_0}^{e_i}$ (a square root of q_{ii}) for all $i \in I$.

As recorded in §2.1.1 above, we fix positive, relatively prime integers d_1, \dots, d_n such that the diagonal matrix $D = \text{diag}(d_1, \dots, d_n)$ symmetrizes A , i.e., DA is symmetric; in fact, these d_i 's coincide with the exponents e_i mentioned above.

We introduce now two special cases of Cartan type multiparameter matrices.

Integral type: We say that $\mathbf{q} := (q_{ij})_{i,j \in I}$ is of *integral type* if it is of Cartan type with associated Cartan matrix $A = (a_{ij})_{i,j \in I}$, and there exist $p \in \mathbb{F}^\times$ and $b_{ij} \in \mathbb{Z}$ ($i, j \in I$) such that $b_{ii} = 2d_i$ and $q_{ij} = p^{b_{ij}}$ for $i \neq j \in I$. The Cartan condition (2.4) yields $b_{ij} + b_{ji} = 2d_i a_{ij}$, for $i, j \in I$ (with d_i 's as above). To be precise, we say also that \mathbf{q} is “of integral type (p, B) ”, with $B := (b_{ij})_{i,j \in I}$.

Canonical multiparameter: Given $q \in \mathbb{F}^\times$ and a symmetrisable (generalised) Cartan matrix $A = (a_{ij})_{i,j \in I}$, consider

$$\check{q}_{ij} := q^{d_i a_{ij}} \quad \forall i, j \in I \quad (2.5)$$

with d_i ($i \in I$) given as above. Then these special values of the $q_{ij} = \check{q}_{ij}$'s do satisfy condition (2.4), hence they provide a special example of matrix $\mathbf{q} = \check{\mathbf{q}}$ of Cartan type, to which we shall refer to hereafter as *the “ q -canonical” case*.

Note also that $\check{\mathbf{q}}$ is of integral type (q, DA) .

By the way, when the multiparameter matrix $\mathbf{q} := (q_{ij})_{i,j \in I}$ is *symmetric*, i.e., $q_{ij} = q_{ji}$ (for all $i, j \in I$), then the “Cartan conditions” $q_{ij} q_{ji} = q_{ii}^{a_{ij}}$ read $q_{ij}^2 = q^{2d_i a_{ij}}$, hence $q_{ij} = \pm q^{d_i a_{ij}}$ (for all $i, j \in I$). Thus every symmetric multiparameter is “almost the q -canonical” one, as indeed it is the q -canonical one “up to sign(s)”.

2.2.2. Equivalence and group action for multiparameters. Let $\mathbf{q} := (q_{ij})_{i,j \in I}$ be a multiparameter matrix. The *generalized Dynkin diagram* $D(\mathbf{q})$ associated with \mathbf{q} is a *labelled graph* whose set of vertices is I , where the i -th vertex is labelled with q_{ii} , and there exists an edge between the vertices i and j only if $\widetilde{q}_{ij} := q_{ij} q_{ji} \neq 1$, in which case the edge is decorated by \widetilde{q}_{ij} . An automorphism of a generalized Dynkin diagram $D(\mathbf{q})$ is an automorphisms as a *labelled graph*. In the set $M_n(\mathbb{F}^\times)$ of all \mathbb{F} -valued multiparameters, we consider the relation \sim defined by

$$\mathbf{q}' \sim \mathbf{q}'' \iff q'_{ij} q'_{ji} = q''_{ij} q''_{ji}, \quad q'_{ii} = q''_{ii} \quad \forall i, j \in I := \{1, \dots, n\}. \quad (2.6)$$

This is obviously an equivalence, which is known in the literature as *twist equivalence*. Indeed, \sim is nothing but the equivalence relation associated with the map $\mathbf{q} \mapsto D(\mathbf{q})$, which to every multiparameter \mathbf{q} associates its unique generalized Dynkin diagram $D(\mathbf{q})$. In particular, if θ is an automorphism of $D(\mathbf{q})$ — in the obvious sense — then we have $\theta(\mathbf{q}) := (q_{\theta(i), \theta(j)})_{i,j \in I} \sim \mathbf{q}$.

Observe also that $M_n(\mathbb{F}^\times)$ is a group (isomorphic to the direct product of n^2 copies of F^\times) with the Hadamard product, i.e., the operation — denoted \odot — is the componentwise multiplication. For $\mathbf{q} = (q_{ij})_{i,j \in I} \in M_n(\mathbb{F}^\times)$, we have that

$\mathbf{q}^{-1} = (q_{ij}^{-1})_{i,j \in I}$; moreover, we set $\mathbf{q}^T = (q_{ji})_{i,j \in I}$ and $\mathbf{q}^{-T} := (\mathbf{q}^{-1})^T = (\mathbf{q}^T)^{-1} = (q_{ji}^{-1})_{i,j \in I}$. The group $(M_n(\mathbb{F}^\times); \odot)$ acts on itself by the adjoint action

$$\boldsymbol{\nu} \cdot \mathbf{q} := \boldsymbol{\nu} \odot \mathbf{q} \odot \boldsymbol{\nu}^{-T} \quad \forall \boldsymbol{\nu}, \mathbf{q} \in M_n(\mathbb{F}^\times) \quad (2.7)$$

In down-to-earth terms, this action is given — for all $\boldsymbol{\nu} = (\nu_{ij})_{i,j \in I}$, $\mathbf{q} = (q_{ij})_{i,j \in I} \in M_n(\mathbb{F}^\times)$ — by

$$\boldsymbol{\nu} \cdot \mathbf{q} := \mathbf{q}^\nu = (q_{ij}^\nu := \nu_{ij} q_{ij} \nu_{ji}^{-1})_{i,j \in I} \quad (2.8)$$

The following lemma shows that the equivalence in $M_n(\mathbb{F}^\times)$ induced by this action is just the twist equivalence introduced above (the proof is a trivial calculation):

Lemma 2.2.3. *Let $\tilde{\mathbf{q}} = (\tilde{q}_{ij})_{i,j \in I}$, $\hat{\mathbf{q}} = (\hat{q}_{ij})_{i,j \in I} \in M_n(\mathbb{F}^\times)$. Then*

$$(a) \quad \tilde{\mathbf{q}} \sim \hat{\mathbf{q}} \iff \exists \boldsymbol{\nu} \in M_n(\mathbb{F}^\times) : \hat{\mathbf{q}} = \boldsymbol{\nu} \cdot \tilde{\mathbf{q}}$$

It follows then that the \sim -equivalence classes coincide with the orbits of the $M_n(\mathbb{F}^\times)$ -action on itself. In particular, if $\tilde{\mathbf{q}} \sim \hat{\mathbf{q}}$ then:

$$(b) \text{ the } \boldsymbol{\nu} \text{'s in } M_n(\mathbb{F}^\times) \text{ satisfying } \hat{\mathbf{q}} = \boldsymbol{\nu} \cdot \tilde{\mathbf{q}} \text{ are those such that } \nu_{ij} \nu_{ji}^{-1} = \hat{q}_{ij} \tilde{q}_{ij}^{-1};$$

$$(c) \text{ an explicit } \boldsymbol{\nu} \text{ such that } \hat{\mathbf{q}} = \boldsymbol{\nu} \cdot \tilde{\mathbf{q}} \text{ is given by } \boldsymbol{\nu} = \begin{cases} \hat{q}_{ij} \tilde{q}_{ij}^{-1} & \forall i \leq j \\ 1 & \forall i > j \end{cases}$$

(d) *if \mathbb{F} contains all square roots of the \tilde{q}_{ij} 's and \hat{q}_{ij} 's, and we fix them so that for the multiparameters formed by these square roots we have $\tilde{\mathbf{q}}^{1/2} \sim \hat{\mathbf{q}}^{1/2}$ (as it is always possible) then another explicit $\boldsymbol{\nu}$ such that $\hat{\mathbf{q}} = \boldsymbol{\nu} \cdot \tilde{\mathbf{q}}$ is given by $\boldsymbol{\nu} = \hat{\mathbf{q}}^{1/2} \tilde{\mathbf{q}}^{-1/2}$. \square*

For multiparameters of Cartan type we also have the following:

Proposition 2.2.4. *Each multiparameter \mathbf{q} of Cartan type is \sim -equivalent to some multiparameter $\tilde{\mathbf{q}}$ of canonical type.*

Proof. Let $\mathbf{q} = (q_{ij})_{i,j \in I} \in M_n(\mathbb{F}^\times)$ be a multiparameter of Cartan type. By definition, — cf. §2.2.1 — this means that there exists $q \in \mathbb{F}^\times$ and a (generalised) symmetrisable Cartan matrix $A = (a_{ij})_{i,j \in I}$ such that $q_{ij} q_{ji} = q_{ii}^{a_{ij}}$ — see (2.4) — and $q_{ii} = q^{2d_i}$ for all $i, j \in I$; altogether these imply

$$q_{ij} q_{ji} = q_{ii}^{a_{ij}} = q^{2d_i a_{ij}} = q^{d_i a_{ij}} q^{d_j a_{ji}} = \tilde{q}_{ij} \tilde{q}_{ji} \quad \forall i, j \in I$$

while $q_{ii} = q^{2d_i}$ also reads $q_{ii} = q^{2d_i} = \tilde{q}_{ii}$, so that in the end we have $\mathbf{q} \sim \tilde{\mathbf{q}}$ with $\tilde{\mathbf{q}} := (q^{d_i a_{ij}})_{i,j \in I}$ a multiparameter of canonical type as in the claim. \square

2.3. Conventions for Hopf algebras.

Our main references for the theory of Hopf algebras are [Mo] and [Ra]. We use standard notation: the comultiplication is denoted Δ and the antipode \mathcal{S} . For the first, we use the Sweedler-Heyneman notation but with the summation sign suppressed, namely we write the coproduct as $\Delta(x) = x_{(1)} \otimes x_{(2)}$.

Although in most parts of the paper it is not needed, we assume that the antipode \mathcal{S} is bijective and denote by \mathcal{S}^{-1} its composition inverse.

Hereafter by \mathbb{k} we denote the ground ring of our algebras, coalgebras, etc.

In any coalgebra C , the set of group-like elements is denoted by $G(C)$; also, we denote by $C^+ := \text{Ker}(\epsilon)$ the augmentation ideal, where $\epsilon : C \rightarrow \mathbb{k}$ is the

count map. If $g, h \in G(C)$, the set of (g, h) -primitive elements is defined as $P_{g,h}(C) := \{x \in C \mid \Delta(x) = x \otimes g + h \otimes x\}$.

If H is a Hopf algebra (or just a bialgebra), we write H^{op} , resp. H^{cop} , for the Hopf algebra (or bialgebra) given by taking in H the opposite product, resp. coproduct.

Finally, we recall the notion of *skew-Hopf pairing* between two Hopf algebras (taken from [AY], §2.1, but essentially standard):

Definition 2.3.1. Given two Hopf algebras H and K with bijective antipode over the ring \mathbb{k} , a \mathbb{k} -linear map $\eta : H \otimes_{\mathbb{k}} K \longrightarrow \mathbb{k}$ is called a *skew-Hopf pairing* (between H and K) if, for all $h \in H$, $k \in K$, one has

$$\begin{aligned} \eta(h, k' k'') &= \eta(h_{(1)}, k') \eta(h_{(2)}, k'') \quad , \quad \eta(h' h'', k) = \eta(h', k_{(2)}) \eta(h'', k_{(1)}) \\ \eta(h, 1) &= \epsilon(h) \quad , \quad \eta(1, k) = \epsilon(k) \quad , \quad \eta(\mathcal{S}^{\pm 1}(h), k) = \eta(h, \mathcal{S}^{\mp 1}(k)) \end{aligned}$$

Recall that, given two Hopf algebras H and K over \mathbb{k} , and a skew-Hopf pairing, say $\eta : H \otimes_{\mathbb{k}} K \longrightarrow \mathbb{k}$, the *Drinfeld double* $D(H, K, \eta)$ is the quotient algebra $T(H \oplus K)/\mathcal{I}$ where \mathcal{I} is the (two-sided) ideal generated by the relations

$$\begin{aligned} 1_H = 1 &= 1_K \quad , \quad a \otimes b = ab \quad \forall a, b \in H \quad \text{or} \quad a, b \in K \quad , \\ x_{(1)} \otimes y_{(1)} \eta(y_{(2)}, x_{(2)}) &= \eta(y_{(1)}, x_{(1)}) y_{(2)} \otimes x_{(2)} \quad \forall x \in K, y \in H \quad ; \end{aligned}$$

such a quotient \mathbb{k} -algebra is also endowed with a standard Hopf algebra structure, which is *consistent*, in that both H and K are Hopf \mathbb{k} -subalgebras of it.

2.4. Hopf algebra deformations.

There exist two standard methods to deform Hopf algebras, usually called “2-cocycle deformations” and “twist deformations”: we shortly recall both, giving them new names; then later on in the paper we shall apply them to quantum groups.

2.4.1. Multiplication twistings (or “2-cocycle deformations”). Let us consider a bialgebra $(H, m, 1, \Delta, \epsilon)$ over a ring \mathbb{k} . A convolution invertible linear map σ in $\text{Hom}_{\mathbb{k}}(H \otimes H, \mathbb{k})$ is called a *normalized multiplicative (or Hopf) 2-cocycle* if

$$\sigma(b_{(1)}, c_{(1)}) \sigma(a, b_{(2)} c_{(2)}) = \sigma(a_{(1)}, b_{(1)}) \sigma(a_{(2)} b_{(2)}, c)$$

and $\sigma(a, 1) = \epsilon(a) = \sigma(1, a)$ for all $a, b, c \in H$, see [Mo, Sec. 7.1]. We will simply call it a 2-cocycle if no confusion arises.

Using a 2-cocycle σ it is possible to define a new algebra structure on H by deforming the multiplication. Indeed, define $m_{\sigma} = \sigma * m * \sigma^{-1} : H \otimes H \longrightarrow H$ by

$$m_{\sigma}(a, b) = a \cdot_{\sigma} b = \sigma(a_{(1)}, b_{(1)}) a_{(2)} b_{(2)} \sigma^{-1}(a_{(3)}, b_{(3)}) \quad \forall a, b \in H$$

If in addition H is a Hopf algebra with antipode \mathcal{S} , then define also $\mathcal{S}_{\sigma} : H \longrightarrow H$ as $\mathcal{S}_{\sigma} : H \longrightarrow H$ where

$$\mathcal{S}_{\sigma}(a) = \sigma(a_{(1)}, \mathcal{S}(a_{(2)})) \mathcal{S}(a_{(3)}) \sigma^{-1}(\mathcal{S}(a_{(4)}), a_{(5)}) \quad \forall a \in H$$

It is known that $(H, m_{\sigma}, 1, \Delta, \epsilon)$ is in turn a bialgebra, and $(H, m_{\sigma}, 1, \Delta, \epsilon, \mathcal{S}_{\sigma})$ is a Hopf algebra (see [DT] for details): this new bialgebra or Hopf algebra structure on H , graphically denoted by H_{σ} , is usually called “cocycle deformation” of the old one; we adopt instead hereafter the terminology “multiplication twisting”.

2.4.2. Comultiplication twistings (or “twist deformations”). Let H be a Hopf algebra (over a commutative ring), and let $\mathcal{F} \in H \otimes H$ be an invertible element in $H^{\otimes 2}$ (later called a “twist”, or “twisting element”) such that

$$\mathcal{F}_{12}(\Delta \otimes \text{id})(\mathcal{F}) = \mathcal{F}_{23}(\text{id} \otimes \Delta)(\mathcal{F}) \quad , \quad (\epsilon \otimes \text{id})(\mathcal{F}) = 1 = (\text{id} \otimes \epsilon)(\mathcal{F})$$

Then H bears a second Hopf algebra structure, denoted $H^{\mathcal{F}}$, with the old product, unit and counit, but with new “twisted” coproduct $\Delta^{\mathcal{F}}$ and antipode $\mathcal{S}^{\mathcal{F}}$ given by

$$\Delta^{\mathcal{F}}(x) := \mathcal{F} \Delta(x) \mathcal{F}^{-1} \quad , \quad \mathcal{S}^{\mathcal{F}}(x) := v \mathcal{S}(x) v^{-1} \quad \forall x \in H \quad (2.9)$$

where $v := \sum_{\mathcal{F}} \mathcal{S}(f'_1) f'_2$ — with $\sum_{\mathcal{F}} f'_1 \otimes f'_2 = \mathcal{F}^{-1}$ — is invertible in H ; see [CP], §4.2.E, and references therein, for further details. Note also that if H is just a bialgebra then the procedure applies as well, with $H^{\mathcal{F}}$ being just a bialgebra.

The Hopf algebra (or just bialgebra) $H^{\mathcal{F}}$ is usually said to be a “twist deformation” of H : we adopt here instead the terminology “comultiplication twisting” to denote both this deformation procedure and its final outcome $H^{\mathcal{F}}$.

2.4.3. Deformations and duality. The two notions of “2-cocycle” and of “twist” are dual to each other with respect to Hopf duality. In detail, we have the following:

Proposition 2.4.4. *Let H be a Hopf algebra over a field, H^* be its linear dual, and H° its Sweedler dual.*

(a) *Let \mathcal{F} be a twist for H , and $\sigma_{\mathcal{F}}$ the image of \mathcal{F} in $(H^* \otimes H^*)^*$ for the natural composed embedding $H \otimes H \hookrightarrow H^{**} \otimes H^{**} \hookrightarrow (H^* \otimes H^*)^*$. Then $\sigma_{\mathcal{F}}$ is a 2-cocycle for H° , and there exists a canonical isomorphism $(H^\circ)_{\sigma_{\mathcal{F}}} \cong (H^{\mathcal{F}})^\circ$.*

(b) *Let σ be a 2-cocycle for H ; assume that H is finite-dimensional, and let \mathcal{F}_σ be the image of σ in the natural identification $(H \otimes H)^* = H^* \otimes H^*$. Then \mathcal{F}_σ is a twist for H^* , and there exists a canonical isomorphism $(H^*)^{\mathcal{F}_\sigma} \cong (H_\sigma)^*$.*

Proof. The proof is an exercise in Hopf duality theory, left to the reader — a matter of reading the identities that characterize a twist or a 2-cocycle in dual terms. \square

As a final remark, let us mention that the notions of 2-cocycle and twist and of the associated deformations (by “twisting”, as precised above) can also be extended to more general contexts where “Hopf algebra” has a broader meaning — essentially, taking place in more general tensor categories than vector spaces over a field with algebraic tensor product. In these cases, if a suitable notion of Hopf “duality” is established (possibly involving different tensor categories), then Proposition 2.4.4 still makes sense, with no need of the finite-dimensionality assumption in claim (b). For instance, this applies to Drinfeld’s “quantized universal enveloping algebras” $U_{\hbar}(\mathfrak{g})$ and their dual “quantized formal series Hopf algebras” $F_{\hbar}[G]$, that both are topological Hopf algebras (with respect to different topologies) over $\mathbb{k}[[\hbar]]$.

3. QUANTUM GROUPS (AS QUEA’S) AND THEIR COMULTIPLICATION TWISTINGS

In this section we briefly recall the notion of formal and polynomial “quantized universal enveloping algebra” (or QUEA in short), following Drinfeld and Jimbo-Lusztig. Then we discuss their comultiplication twistings (via so-called “toral twists”), and we provide an alternative presentation of these twisted QUEA’s.

3.1. Formal QUEA's (“à la Drinfeld”).

We begin by the description of formal quantized universal enveloping algebra following Drinfeld and others, and some related tools.

3.1.1. The formal QUEA $U_{\hbar}(\mathfrak{g})$. Let $\mathbb{k}[[\hbar]]$ be the ring of formal power series in \hbar . Let $A := (a_{ij})_{i,j \in I}$ be a generalized symmetrisable Cartan matrix. The quantized universal enveloping algebra, or QUEA in short, $U_{\hbar}(\mathfrak{g}'(A)) = U_{\hbar}(\mathfrak{g})$ is the associative, unital, topologically complete $\mathbb{k}[[\hbar]]$ -algebra with generators H_i , E_i and F_i ($i \in I := \{1, \dots, n\}$) and relations (for all $i, j \in I$)

$$\begin{aligned} H_i H_j - H_j H_i &= 0, & H_i E_j - E_j H_i &= +a_{ij} E_j, & H_i F_j - F_j H_i &= -a_{ij} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{e^{+\hbar d_i H_i} - e^{-\hbar d_i H_i}}{e^{+\hbar d_i} - e^{-\hbar d_i}} \\ \sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1-a_{ij} \\ k \end{bmatrix}_{q_i} E_i^{1-a_{ij}-k} E_j E_i^k &= 0 & (i \neq j) \end{aligned} \quad (3.1)$$

$$\sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1-a_{ij} \\ k \end{bmatrix}_{q_i} F_i^{1-a_{ij}-k} F_j F_i^k = 0 \quad (i \neq j) \quad (3.2)$$

where hereafter we use shorthand notation $e^X := \exp(X)$, $q := e^{\hbar}$, $q_i := q^{d_i} = e^{\hbar d_i}$. It is known that $U_{\hbar}(\mathfrak{g})$ has a structure of (topological) Hopf algebra, given by

$$\begin{aligned} \Delta(E_i) &:= E_i \otimes 1 + e^{+\hbar d_i H_i} \otimes E_i, & \mathcal{S}(E_i) &:= -e^{-\hbar d_i H_i} E_i, & \epsilon(E_i) &:= 0 \\ \Delta(H_i) &:= H_i \otimes 1 + 1 \otimes H_i, & \mathcal{S}(H_i) &:= -H_i, & \epsilon(H_i) &:= 0 \\ \Delta(F_i) &:= F_i \otimes e^{-\hbar d_i H_i} + 1 \otimes F_i, & \mathcal{S}(F_i) &:= -F_i e^{+\hbar d_i H_i}, & \epsilon(F_i) &:= 0 \end{aligned}$$

(for all $i \in I$) where the coproduct takes values in the \hbar -adic completion $U_{\hbar}(\mathfrak{g})^{\widehat{\otimes} 2}$ of the algebraic tensor square $U_{\hbar}(\mathfrak{g})^{\otimes 2}$ — see, e.g., [CP] (and references therein) for details, taking into account that we adopt slightly different normalizations.

Finally, the semiclassical “limit” $U_{\hbar}(\mathfrak{g}) / \hbar U_{\hbar}(\mathfrak{g})$ of $U_{\hbar}(\mathfrak{g})$ is isomorphic, as a Hopf algebra, to $U(\mathfrak{g})$ — via $E_i \mapsto e_i$, $H_i \mapsto h_i$, $F_i \mapsto f_i$ (notation of §2.1.1) — hence $U_{\hbar}(\mathfrak{g})$ is a deformation quantization of $U(\mathfrak{g}'(A))$, the universal enveloping algebra of the derived algebra of the Kac-Moody Lie algebra $\mathfrak{g}(A)$. Then the latter is endowed with a structure of *co-Poisson* Hopf algebra, which makes $\mathfrak{g}'(A)$ into a Lie bialgebra.

Remark 3.1.2. Clearly, one can slightly modify the definition of $U_{\hbar}(\mathfrak{g}'(A)) = U_{\hbar}(\mathfrak{g})$ by allowing additional generators corresponding to the elements of the \mathbb{k} -basis of any vector space complement \mathfrak{h}'' of \mathfrak{h}' inside \mathfrak{g} (notation of §2.1.1), along with the obvious, natural relations: this provides a larger (if $\mathfrak{g}' \subsetneq \mathfrak{g}$) quantum Hopf algebra which is a formal quantization of $U_{\hbar}(\mathfrak{g}(A))$ and so endows the latter with a structure of *co-Poisson* Hopf algebra, thus making $\mathfrak{g}(A)$ into a Lie bialgebra for whom $\mathfrak{g}'(A)$ is then a Lie sub-bialgebra. *For this larger quantum Hopf algebra one can repeat, up to minimal changes, all what we do from now on.*

3.1.3. Quantum Borel (sub)algebras and their Drinfeld double. We denote hereafter by $U_{\hbar}(\mathfrak{h})$, resp. $U_{\hbar}(\mathfrak{b}_+)$, resp. $U_{\hbar}(\mathfrak{b}_-)$, the \hbar -adically complete subalgebra of $U_{\hbar}(\mathfrak{g})$ generated by all the H_i 's, resp. the H_i 's and the E_i 's, resp. the H_i 's and the F_i 's. We refer to $U_{\hbar}(\mathfrak{h})$ as *quantum Cartan (sub)algebra* and to $U_{\hbar}(\mathfrak{b}_+)$, resp.

$U_{\hbar}(\mathfrak{b}_-)$, as *quantum positive*, resp. *negative*, Borel (sub)algebra. It follows directly from definitions that $U_{\hbar}(\mathfrak{h})$, $U_{\hbar}(\mathfrak{b}_+)$ and $U_{\hbar}(\mathfrak{b}_-)$ all are Hopf subalgebras of $U_{\hbar}(\mathfrak{g})$.

It is also known that the quantum Borel subalgebras are related via a skew-Hopf pairing $\eta : U_{\hbar}(\mathfrak{b}_+) \otimes_{\mathbb{k}[[\hbar]]} U_{\hbar}(\mathfrak{b}_-) \longrightarrow \mathbb{k}[[\hbar]]$ given by

$$\eta(H_i, H_j) = -d_i a_{ij}, \quad \eta(E_i, F_j) = \delta_{ij} \frac{1}{q_i^{-1} - q_i}, \quad \eta(E_i, H_j) = 0 = \eta(H_i, F_j) \quad (3.3)$$

Using this skew-Hopf pairing, one considers as in §2.3 the corresponding Drinfeld double $D(U_{\hbar}(\mathfrak{b}_+), U_{\hbar}(\mathfrak{b}_-), \eta)$; in the sequel we denote the latter by $U_{\hbar}(\mathfrak{g}_D)$.

By construction, there is a Hopf algebra epimorphism $\pi_{\mathfrak{g}} : U_{\hbar}(\mathfrak{g}_D) \longrightarrow U_{\hbar}(\mathfrak{g})$. In order to describe it, we use the identification

$$U_{\hbar}(\mathfrak{g}_D) = D(U_{\hbar}(\mathfrak{b}_+), U_{\hbar}(\mathfrak{b}_-), \eta) \cong U_{\hbar}(\mathfrak{b}_+) \otimes_{\mathbb{k}[[\hbar]]} U_{\hbar}(\mathfrak{b}_-)$$

as coalgebras, and adopt such shorthand notation as $E_i = E_i \otimes 1$, $H_i^+ = H_i \otimes 1$, $H_i^- = 1 \otimes H_i$, $F_i = 1 \otimes F_i$ (for all $i \in I$); then the projection $\pi_{\mathfrak{g}}$ is determined by

$$\pi_{\mathfrak{g}} : \quad E_i \mapsto E_i, \quad H_i^+ \mapsto H_i, \quad H_i^- \mapsto H_i, \quad F_i \mapsto F_i \quad \forall i \in I \quad (3.4)$$

Furthermore, $U_{\hbar}(\mathfrak{g}_D)$ can be explicitly described as follows: it is the associative, unital, topologically complete $\mathbb{k}[[\hbar]]$ -algebra with generators E_i , H_i^+ , H_i^- and F_i ($i \in I := \{1, \dots, n\}$) satisfying the relations (3.1), (3.2) and the following:

$$\begin{aligned} H_i^+ H_j^+ - H_j^+ H_i^+ &= 0, & H_i^+ H_j^- - H_j^- H_i^+ &= 0, & H_i^- H_j^- - H_j^- H_i^- &= 0 \\ H_i^{\pm} E_j - E_j H_i^{\pm} &= \pm a_{ij} E_j, & H_i^{\pm} F_j - F_j H_i^{\pm} &= \mp a_{ij} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{e^{+\hbar d_i H_i^+} - e^{-\hbar d_i H_i^-}}{e^{+\hbar d_i} - e^{-\hbar d_i}} \end{aligned}$$

In addition, the structure of (topological) Hopf algebra of $U_{\hbar}(\mathfrak{g}_D)$ is given by

$$\begin{aligned} \Delta(E_i) &:= E_i \otimes 1 + e^{+\hbar d_i H_i^+} \otimes E_i, & \mathcal{S}(E_i) &:= -e^{-\hbar d_i H_i^+} E_i, & \epsilon(E_i) &:= 0 \\ \Delta(H_i^{\pm}) &:= H_i^{\pm} \otimes 1 + 1 \otimes H_i^{\pm}, & \mathcal{S}(H_i^{\pm}) &:= -H_i^{\pm}, & \epsilon(H_i^{\pm}) &:= 0 \\ \Delta(F_i) &:= F_i \otimes e^{-\hbar d_i H_i^-} + 1 \otimes F_i, & \mathcal{S}(F_i) &:= -F_i e^{+\hbar d_i H_i^-}, & \epsilon(F_i) &:= 0 \end{aligned}$$

By construction, both $U_{\hbar}(\mathfrak{b}_+)$ and $U_{\hbar}(\mathfrak{b}_-)$ are Hopf subalgebras of $U_{\hbar}(\mathfrak{g}_D)$: the natural embeddings $U_{\hbar}(\mathfrak{b}_+) \hookrightarrow U_{\hbar}(\mathfrak{g}_D)$ and $U_{\hbar}(\mathfrak{b}_-) \hookrightarrow U_{\hbar}(\mathfrak{g}_D)$ are described by

$$E_i \mapsto E_i, \quad H_i \mapsto H_i^+ \quad \text{and} \quad F_i \mapsto F_i, \quad H_i \mapsto H_i^- \quad \forall i \in I$$

Like for $U_{\hbar}(\mathfrak{g})$, if we look at the semiclassical limit of $U_{\hbar}(\mathfrak{h})$, resp. of $U_{\hbar}(\mathfrak{b}_+)$, resp. of $U_{\hbar}(\mathfrak{b}_-)$, resp. of $U_{\hbar}(\mathfrak{g}_D)$, we find $U(\mathfrak{h}'(A))$, resp. $U(\mathfrak{b}'_+(A))$, resp. $U(\mathfrak{b}'_-(A))$, resp. $U(\mathfrak{g}'_D(A))$, where $\mathfrak{g}'_D(A)$ is the Manin double (see §2.1.1). This entails that all these universal enveloping algebras also are *co-Poisson* Hopf algebras, whose co-Poisson structure in the first three cases is just the restriction of that of $U(\mathfrak{g}'(A))$; in particular, all of $\mathfrak{h}'(A)$, $\mathfrak{b}'_+(A)$ and $\mathfrak{b}'_-(A)$ are Lie sub-bialgebras of $\mathfrak{g}'(A)$.

Finally, a suitably modified version of Remark 3.1.2 applies to $U_{\hbar}(\mathfrak{g}_D)$ as well.

3.2. Comultiplication twistings of formal QUEA's.

Following an idea of Reshetikhin (cf. [Re]), we shall consider special twisting elements for $U_{\hbar}(\mathfrak{g})$ and use them to provide a new, twisted QUEA denoted $U_{\hbar}^{\Psi}(\mathfrak{g})$. Then we shall extend the same method to $U_{\hbar}(\mathfrak{g}_D)$ as well.

3.2.1. Comultiplication twistings of $U_{\hbar}(\mathfrak{g})$. Let us consider an $(n \times n)$ -matrix $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{k}[[\hbar]])$. A straightforward check shows that the element

$$\mathcal{F}_{\Psi} := \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i \otimes H_j\right) \quad (3.5)$$

is actually a *twist* of $U_{\hbar}(\mathfrak{g})$ in the sense of §2.4.2. Indeed, since $U_{\hbar}(\mathfrak{h}_D)$ is commutative, all computations follow from the simple fact that $\exp(\hbar a)\exp(\hbar b) = \exp(\hbar(a+b))$ for all commuting elements a and b ; for instance, this implies that \mathcal{F}_{Ψ} is invertible in $U_{\hbar}(\mathfrak{g}_D)^{\widehat{\otimes}^2}$ with inverse $\mathcal{F}_{\Psi}^{-1} := \exp\left(-\hbar \sum_{i,j=1}^n \psi_{ij} H_i \otimes H_j\right)$. Also, since ϵ is an algebra map (which is continuous in the \hbar -adic topology) we have that

$$(\epsilon \otimes \text{id})(\mathcal{F}_{\Psi}) = \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} \epsilon(H_i) \otimes H_j\right) = 1$$

and

$$(\text{id} \otimes \epsilon)(\mathcal{F}_{\Psi}) = \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i \otimes \epsilon(H_j)\right) = 1$$

Finally, let us check that \mathcal{F}_{Ψ} satisfies the cocycle condition. Since Δ is also an algebra map (again, continuous in the \hbar -adic topology), we have

$$\begin{aligned} (\mathcal{F}_{\Psi})_{12} (\Delta \otimes \text{id})(\mathcal{F}_{\Psi}) &= \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i \otimes H_j \otimes 1\right) \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} \Delta(H_i) \otimes H_j\right) \\ &= \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} (H_i \otimes H_j \otimes 1 + H_i \otimes 1 \otimes H_j + 1 \otimes H_i \otimes H_j)\right) \\ (\mathcal{F}_{\Psi})_{23} (\text{id} \otimes \Delta)(\mathcal{F}_{\Psi}) &= \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} 1 \otimes H_i \otimes H_j\right) \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i \otimes \Delta(H_j)\right) \\ &= \exp\left(\hbar \sum_{i,j=1}^n \psi_{ij} (1 \otimes H_i \otimes H_j + H_i \otimes H_j \otimes 1 + H_i \otimes 1 \otimes H_j)\right) \end{aligned}$$

so that $(\mathcal{F}_{\Psi})_{12} (\Delta \otimes \text{id})(\mathcal{F}_{\Psi}) = (\mathcal{F}_{\Psi})_{23} (\text{id} \otimes \Delta)(\mathcal{F}_{\Psi})$ as expected.

Therefore, the recipe in §2.4.2 endows $U_{\hbar}(\mathfrak{g})$ with a new, “twisted” coproduct $\Delta^{(\Psi)} := \Delta^{\mathcal{F}_{\Psi}}$ which gives us a new Hopf structure $U_{\hbar}(\mathfrak{g})^{\mathcal{F}_{\Psi}}$. An easy computation proves that the new coproduct reads as follows on generators (for $\ell \in I$):

$$\begin{aligned} \Delta^{(\Psi)}(E_{\ell}) &= E_{\ell} \otimes e^{+\hbar \sum_{i,j \in I} \psi_{ij} a_{i\ell} H_j} + e^{+\hbar d_{\ell} H_{\ell} + \hbar \sum_{i,j \in I} \psi_{ij} a_{j\ell} H_i} \otimes E_{\ell} \\ \Delta^{(\Psi)}(H_{\ell}) &= H_{\ell} \otimes 1 + 1 \otimes H_{\ell} \\ \Delta^{(\Psi)}(F_{\ell}) &= F_{\ell} \otimes e^{-\hbar d_{\ell} H_{\ell} - \hbar \sum_{i,j \in I} \psi_{ij} a_{i\ell} H_j} + e^{-\hbar \sum_{i,j \in I} \psi_{ij} a_{j\ell} H_i} \otimes F_{\ell} \end{aligned}$$

Using notation of §2.1.2 with $\mathcal{R} := \mathbb{k}[[\hbar]]$ and $T_\ell = d_\ell H_\ell$ for all $\ell \in I$, these formulas read (for $\ell \in I$)

$$\begin{aligned}\Delta^{(\Psi)}(E_\ell) &= E_\ell \otimes e^{+\hbar\psi_-^{b'}(T_\ell)} + e^{+\hbar(\text{id}_{\mathfrak{b}'} + \psi_+^{b'})}(T_\ell) \otimes E_\ell \\ \Delta^{(\Psi)}(H_\ell) &= H_\ell \otimes 1 + 1 \otimes H_\ell \\ \Delta^{(\Psi)}(F_\ell) &= F_\ell \otimes e^{-\hbar(\text{id}_{\mathfrak{b}'} + \psi_-^{b'})}(T_\ell) + e^{-\hbar\psi_+^{b'}(T_\ell)} \otimes F_\ell\end{aligned}$$

Similarly, the “twisted” antipode $\mathcal{S}^{(\Psi)} := \mathcal{S}^{\mathcal{F}^\Psi}$ is expressed by (for all $\ell \in I$)

$$\begin{aligned}\mathcal{S}^{(\Psi)}(E_\ell) &= -e^{-\hbar(\text{id}_{\mathfrak{b}'} + \psi_+^{b'})}(T_\ell) E_\ell e^{-\hbar\psi_-^{b'}(T_\ell)} \\ \mathcal{S}^{(\Psi)}(H_\ell) &= -H_\ell \\ \mathcal{S}^{(\Psi)}(F_\ell) &= -e^{+\hbar\psi_+^{b'}(T_\ell)} F_\ell e^{+\hbar(\text{id}_{\mathfrak{b}'} + \psi_-^{b'})}(T_\ell)\end{aligned}$$

and the (untwisted!) counit by $\epsilon(E_\ell) = 0$, $\epsilon(H_\ell) = 0$, $\epsilon(F_\ell) = 0$ again (for $\ell \in I$).

In the sequel we shall use the shorter notation $U_\hbar^\Psi(\mathfrak{g})$ to denote $U_\hbar(\mathfrak{g})^{\mathcal{F}^\Psi}$.

The previous formulas show that the quantum Borel subalgebras $U_\hbar(\mathfrak{b}_+)$ and $U_\hbar(\mathfrak{b}_-)$ are still Hopf subalgebras inside $U_\hbar^\Psi(\mathfrak{g})$ as well. In fact, the element \mathcal{F}^Ψ can also be seen as a twisting element for both $U_\hbar(\mathfrak{b}_+)$ and $U_\hbar(\mathfrak{b}_-)$, and then the corresponding twisted Hopf algebras obviously sit as Hopf subalgebras inside $U_\hbar^\Psi(\mathfrak{g})$.

On the other hand, the previous formulas suggest a different presentation of the quantum Borel subalgebras. *Hereafter, we assume for simplicity that the Cartan matrix A be of finite type*; actually, this assumption can be lifted — indeed, we do it in [GG2] — but it makes our discussion definitely simpler.

Namely, let us consider the “twisted generators”

$$E_i^\Psi := e^{-\hbar\psi_-^{b'}(T_i)} E_i \quad , \quad T_{i,+}^\Psi := (\text{id} + \psi_+^{b'} - \psi_-^{b'})(T_i) \quad \forall i \in I$$

in $U_\hbar^\Psi(\mathfrak{b}_+)$: since, by Lemma 2.1.4(a) — which exploits the finiteness assumption on A — the set $\{T_{i,+}^\Psi\}_{i \in I}$ is a $\mathbb{k}[[\hbar]]$ -basis of $\mathfrak{h}'_{\mathbb{k}[[\hbar]]}$, we have that it generates $U_\hbar^\Psi(\mathfrak{b}_+) = U_\hbar(\mathfrak{b}_+)$; in particular, the $\mathbb{k}[[\hbar]]$ -subalgebra generated (in topological sense) by these twisted generators coincide with $U_\hbar^\Psi(\mathfrak{b}_+)$. Analogously, consider the “twisted generators”

$$T_{i,-}^\Psi := (\text{id} + \psi_-^{b'} - \psi_+^{b'})(T_i) \quad , \quad F_i^\Psi := e^{+\hbar\psi_+^{b'}(T_i)} F_i \quad \forall i \in I$$

for $U_\hbar^\Psi(\mathfrak{b}_-)$. Then the above formulas for $\Delta^{(\Psi)}$ give (for all $\ell \in I$)

$$\begin{aligned}\Delta^{(\Psi)}(E_\ell^\Psi) &= E_\ell^\Psi \otimes 1 + e^{+\hbar T_{\ell,+}^\Psi} \otimes E_\ell^\Psi \\ \Delta^{(\Psi)}(T_{\ell,\pm}^\Psi) &= T_{\ell,\pm}^\Psi \otimes 1 + 1 \otimes T_{\ell,\pm}^\Psi \\ \Delta^{(\Psi)}(F_\ell^\Psi) &= F_\ell^\Psi \otimes e^{-\hbar T_{\ell,-}^\Psi} + 1 \otimes F_\ell^\Psi\end{aligned} \tag{3.6}$$

and those for $\epsilon^{(\Psi)} := \epsilon$ and $\mathcal{S}^{(\Psi)}$ yield (for all $\ell \in I$)

$$\begin{aligned}\epsilon^{(\Psi)}(E_\ell^\Psi) &= 0 \quad , \quad \epsilon^{(\Psi)}(T_{\ell,\pm}^\Psi) = 0 \quad , \quad \epsilon^{(\Psi)}(F_\ell^\Psi) = 0 \\ \mathcal{S}(E_\ell^\Psi) &= -e^{-\hbar T_{\ell,+}^\Psi} E_\ell^\Psi \quad , \quad \mathcal{S}(T_{\ell,\pm}^\Psi) = -T_{\ell,\pm}^\Psi \quad , \quad \mathcal{S}(F_\ell^\Psi) = -F_\ell^\Psi e^{+\hbar T_{\ell,-}^\Psi}\end{aligned} \tag{3.7}$$

Indeed, using these generators we may write down a complete presentation by generators and relations of $U_\hbar^\Psi(\mathfrak{g})$ and similarly also of $U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar^\Psi(\mathfrak{b}_-)$; in fact,

a complete set of relations is the following (for all $i, j \in I$)

$$\begin{aligned}
(\text{id} + \psi_+^{\mathfrak{h}'} - \psi_-^{\mathfrak{h}'})^{-1}(T_{i,+}^\Psi) &= (\text{id} + \psi_-^{\mathfrak{h}'} - \psi_+^{\mathfrak{h}'})^{-1}(T_{i,-}^\Psi) \\
T_{i,\pm}^\Psi T_{i,\pm}^\Psi &= T_{j,\pm}^\Psi T_{i,\pm}^\Psi \quad , \quad T_{i,\pm}^\Psi T_{j,\mp}^\Psi = T_{j,\mp}^\Psi T_{i,\pm}^\Psi \\
T_{i,+}^\Psi E_j^\Psi - E_j^\Psi T_{i,+}^\Psi &= +p_{i,j} E_j^\Psi \quad , \quad T_{i,-}^\Psi E_j^\Psi - E_j^\Psi T_{i,-}^\Psi = +p_{j,i} E_j^\Psi \\
T_{i,+}^\Psi F_j^\Psi - F_j^\Psi T_{i,+}^\Psi &= -p_{i,j} F_j^\Psi \quad , \quad T_{i,-}^\Psi F_j^\Psi - F_j^\Psi T_{i,-}^\Psi = -p_{j,i} F_j^\Psi \\
E_i^\Psi F_j^\Psi - F_j^\Psi E_i^\Psi &= \delta_{i,j} \frac{e^{+\hbar T_{i,+}^\Psi} - e^{-\hbar T_{i,-}^\Psi}}{e^{+\hbar p_{ii}/2} - e^{-\hbar p_{ii}/2}} \quad (3.8)
\end{aligned}$$

$$\begin{aligned}
\sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1 - a_{ij} \\ k \end{bmatrix}_{e^{\hbar p_{ii}/2}} (E_i^\Psi)^{1-a_{ij}-k} E_j^\Psi (E_i^\Psi)^k &= 0 \quad (i \neq j) \\
\sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1 - a_{ij} \\ k \end{bmatrix}_{e^{\hbar p_{ii}/2}} (F_i^\Psi)^{1-a_{ij}-k} F_j^\Psi (F_i^\Psi)^k &= 0 \quad (i \neq j)
\end{aligned}$$

with $p_{i,j} := \alpha_j((\text{id} + \psi_+^{\mathfrak{h}'} - \psi_-^{\mathfrak{h}'})(T_i)) = (DA + A^T(\Psi^T - \Psi)A)_i^j \in \mathbb{k}[[\hbar]]$ — where the very first relation above makes sense because of Lemma 2.1.4(a).

Note, by the way, that the matrix $P := (p_{i,j})_{i,j \in I} = DA + A^T(\Psi^T - \Psi)A$ is also described by $p_{i,j} = \alpha_j((\text{id} + \psi_-^{\mathfrak{h}'} - \psi_+^{\mathfrak{h}'})(T_i))$. Moreover, we notice that the matrix $P - DA = A^T(\Psi^T - \Psi)A$ is antisymmetric.

As we pointed out, thanks to Lemma 2.1.4(a), both sets $\{T_{i,+}^\Psi\}_{i \in I}$ and $\{T_{i,-}^\Psi\}_{i \in I}$ are $\mathbb{k}[[\hbar]]$ -basis of $\mathfrak{h}'_{\mathbb{k}[[\hbar]]}$, so each one of them is enough to generate $U_{\hbar}^\Psi(\mathfrak{h}) = U_{\hbar}(\mathfrak{h})$. It follows that $U_{\hbar}^\Psi(\mathfrak{g})$ can be generated by either one of the smaller sets $\{E_i^\Psi, T_{i,+}^\Psi, F_i^\Psi\}_{i \in I}$ or $\{E_i^\Psi, T_{i,-}^\Psi, F_i^\Psi\}_{i \in I}$; in both case, however, the presentation would be less good-looking than using the bigger set $\{E_i^\Psi, T_{i,+}^\Psi, T_{i,-}^\Psi, F_i^\Psi\}_{i \in I}$.

3.2.2. Comultiplication twistings of $U_{\hbar}(\mathfrak{g}_D)$. Let $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{k}[[\hbar]])$ as in §3.2.1 above. Again, a direct check shows that the element

$$\mathcal{F}_\Psi := \exp \left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i^+ \otimes H_j^- \right) \quad (3.9)$$

is actually a *twist* for $U_{\hbar}(\mathfrak{g}_D)$ in the sense of §2.4.2. Henceforth, we have analogous formulas for the new coproduct on generators similar to those in §3.2.1, namely

$$\begin{aligned}
\Delta^{(\Psi)}(E_\ell) &= E_\ell \otimes e^{+\hbar \sum_{i,j \in I} \psi_{ij} a_{i\ell} H_j^-} + e^{+\hbar d_\ell H_\ell^+ + \hbar \sum_{i,j \in I} \psi_{ij} a_{j\ell} H_i^+} \otimes E_\ell \\
\Delta^{(\Psi)}(H_\ell^\pm) &= H_\ell^\pm \otimes 1 + 1 \otimes H_\ell^\pm \\
\Delta^{(\Psi)}(F_\ell) &= F_\ell \otimes e^{-\hbar d_\ell H_\ell^- - \hbar \sum_{i,j \in I} \psi_{ij} a_{i\ell} H_j^-} + e^{-\hbar \sum_{i,j \in I} \psi_{ij} a_{j\ell} H_i^+} \otimes F_\ell
\end{aligned}$$

Using notation of §2.1.2 with $\mathcal{R} := \mathbb{k}[[\hbar]]$ and $\{T_i^\pm\}_{i \in I}$ thought of as a basis of \mathfrak{h}'_\pm , the formulas above read

$$\begin{aligned}
\Delta^{(\Psi)}(E_\ell) &= E_\ell \otimes e^{+\hbar \psi_-^{\mathfrak{h}'}(T_\ell^-)} + e^{+\hbar (\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'}) (T_\ell^+)} \otimes E_\ell \\
\Delta^{(\Psi)}(H_\ell) &= H_\ell \otimes 1 + 1 \otimes H_\ell \\
\Delta^{(\Psi)}(F_\ell) &= F_\ell \otimes e^{-\hbar (\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'}) (T_\ell^-)} + e^{-\hbar \psi_+^{\mathfrak{h}'}(T_\ell^+)} \otimes F_\ell
\end{aligned}$$

Similarly, the “twisted” antipode $\mathcal{S}^{(\Psi)} := \mathcal{S}^{\mathcal{F}^\Psi}$ and the (untwisted!) counit $\epsilon^{(\Psi)} := \epsilon$ are expressed by (for all $\ell \in I$)

$$\begin{aligned} \mathcal{S}^{(\Psi)}(E_\ell) &= -e^{-\hbar(\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'})^{(T_\ell^+)}} E_\ell e^{-\hbar\psi_-^{\mathfrak{h}'}(T_\ell^-)} \quad , \quad \epsilon(E_\ell) = 0 \\ \mathcal{S}^{(\Psi)}(H_\ell^\pm) &= -H_\ell^\pm \quad , \quad \epsilon(H_\ell^\pm) = 0 \\ \mathcal{S}^{(\Psi)}(F_\ell) &= -e^{+\hbar\psi_+^{\mathfrak{h}'}(T_\ell^+)} F_\ell e^{+\hbar(\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'})^{(T_\ell^-)}} \quad , \quad \epsilon(F_\ell) = 0 \end{aligned}$$

Once more, acting again like in §3.2.1 — and again assuming for simplicity that the Cartan matrix A be of finite type — we can consider the “twisted” generators (in topological sense, as usual)

$$E_i^\Psi := e^{-\hbar\psi_-^{\mathfrak{h}'}(T_i^-)} E_i \quad , \quad T_{i,+}^\Psi := (\text{id} + \psi_+^{\mathfrak{h}'})^{(T_i^+)} - \psi_-^{\mathfrak{h}'}(T_i^-) \quad \forall i \in I$$

for $U_{\hbar}^\Psi(\mathfrak{b}_+)$ and

$$T_{i,-}^\Psi := (\text{id} + \psi_-^{\mathfrak{h}'})^{(T_i^-)} - \psi_+^{\mathfrak{h}'}(T_i^+) \quad , \quad F_i^\Psi := e^{+\hbar\psi_+^{\mathfrak{h}'}(T_i^+)} F_i \quad \forall i \in I$$

for $U_{\hbar}^\Psi(\mathfrak{b}_-)$. Then the above formulas for the Hopf operations along with the commutation relations in $U_{\hbar}^\Psi(\mathfrak{g}_D) = U_{\hbar}(\mathfrak{g}_D)$ — an identity of algebras! — yield a presentation for $U_{\hbar}^\Psi(\mathfrak{g}_D)$, quite similar to that for $U_{\hbar}^\Psi(\mathfrak{g})$.

Finally, we remark that the epimorphism $\pi_{\mathfrak{g}} : U_{\hbar}(\mathfrak{g}_D) \longrightarrow U_{\hbar}(\mathfrak{g})$ (see §3.1.3) of (topological) Hopf algebras is also an epimorphism for the twisted Hopf structure on both sides, i.e., it is an epimorphism $\pi_{\mathfrak{g}}^\Psi := \pi_{\mathfrak{g}} : U_{\hbar}^\Psi(\mathfrak{g}_D) \longrightarrow U_{\hbar}^\Psi(\mathfrak{g})$. Indeed, this is a direct, easy consequence of the fact that $\pi_{\mathfrak{g}}^{\otimes 2}$ maps the twist element (3.9) of $U_{\hbar}^\Psi(\mathfrak{g}_D)$ onto the twist element (3.5) of $U_{\hbar}^\Psi(\mathfrak{g})$. As a matter of description, using either presentation — the one with the E_i ’s, the $T_{i,\pm}$ ’s and the F_i ’s or the one with the E_i^Ψ ’s, the $T_{i,\pm}^\Psi$ ’s and the F_i^Ψ ’s — for both algebras, it is clear that $\pi_{\mathfrak{g}}^\Psi$ maps each generator of $U_{\hbar}^\Psi(\mathfrak{g}_D)$ onto the same name generator of $U_{\hbar}^\Psi(\mathfrak{g})$.

3.2.3. Twisting Borel subalgebras in $U_{\hbar}(\mathfrak{g}_D)$. The formulas for the twisted coproduct in §3.2.2 above show that *the quantum Borel subalgebras $U_{\hbar}(\mathfrak{b}_+)$ and $U_{\hbar}(\mathfrak{b}_-)$ are no longer (in general) Hopf subalgebras inside $U_{\hbar}^\Psi(\mathfrak{g}_D)$* . Instead, let $\dot{U}_{\hbar}^\Psi(\mathfrak{b}_+)$, resp. $\dot{U}_{\hbar}^\Psi(\mathfrak{b}_-)$, be the complete, unital subalgebra of $U_{\hbar}^\Psi(\mathfrak{g}_D)$ generated by $\{E_i, (\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'})^{(T_i^+)}, \psi_-^{\mathfrak{h}'}(T_i^-)\}_{i \in I}$, resp. by $\{\psi_+^{\mathfrak{h}'}(T_i^+), (\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'})^{(T_i^-)}, F_i\}_{i \in I}$.

Then the same formulas yield

$$\dot{U}_{\hbar}^\Psi(\mathfrak{b}_-) \text{ and } \dot{U}_{\hbar}^\Psi(\mathfrak{b}_+) \text{ are Hopf subalgebras of } U_{\hbar}^\Psi(\mathfrak{g}_D).$$

Clearly, the semiclassical limits of these twisted subalgebras are

$$\dot{U}_{\hbar}^\Psi(\mathfrak{b}_+) / \hbar \dot{U}_{\hbar}^\Psi(\mathfrak{b}_+) \cong U(\dot{\mathfrak{b}}_+^\Psi) \quad \text{and} \quad \dot{U}_{\hbar}^\Psi(\mathfrak{b}_-) / \hbar \dot{U}_{\hbar}^\Psi(\mathfrak{b}_-) \cong U(\dot{\mathfrak{b}}_-^\Psi)$$

where (with notation of §2.1.1)

$$\dot{\mathfrak{b}}_+^\Psi := \text{Lie subalgebra of } \mathfrak{g}'_D(A) \text{ generated by } \{e_i, (\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'})^{(t_i^+)}, \psi_-^{\mathfrak{h}'}(t_i^-)\}_{i \in I}$$

$$\dot{\mathfrak{b}}_-^\Psi := \text{Lie subalgebra of } \mathfrak{g}'_D(A) \text{ generated by } \{\psi_+^{\mathfrak{h}'}(t_i^+), (\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'})^{(t_i^-)}, f_i\}_{i \in I}$$

— where $\psi_\pm^{\mathfrak{h}'}(t_i^\pm) := \psi_\pm^{\mathfrak{h}'}(T_i^\pm) \bmod \hbar \mathbb{K}[[\hbar]] \in \mathfrak{h}'^\pm \subseteq \mathfrak{b}'_\pm(A)$ — so that a more inspiring, self-explaining notation might be $U_{\hbar}(\dot{\mathfrak{b}}_\pm^\Psi) := \dot{U}_{\hbar}^\Psi(\mathfrak{b}_\pm)$.

Note, however, that \mathfrak{b}_+^Ψ and \mathfrak{b}_-^Ψ both have a larger ‘‘Cartan subalgebra’’ than \mathfrak{h}' , so they cannot be correctly thought of as ‘‘twisted’’ Borel subalgebras inside $\mathfrak{g}'_D(A)$.

On the other hand, let us consider the complete, unital subalgebras $U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar^\Psi(\mathfrak{b}_-)$ of $U_\hbar^\Psi(\mathfrak{g}_D)$ generated respectively by the ‘‘twisted generators’’

$$E_i^\Psi := e^{-\hbar\psi_-^{\mathfrak{h}'}(T_i^-)} E_i \quad , \quad T_{i,+}^\Psi := (\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'})(T_i^+) - \psi_-^{\mathfrak{h}'}(T_i^-) \quad \forall i \in I$$

and

$$F_i^\Psi := e^{+\hbar\psi_+^{\mathfrak{h}'}(T_i^+)} F_i \quad , \quad T_{i,-}^\Psi := (\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'})(T_i^-) - \psi_+^{\mathfrak{h}'}(T_i^+) \quad \forall i \in I$$

given in §3.2.2 — still assuming for simplicity that the matrix A be of finite type. Then the above formulas for $\Delta^{(\Psi)}$, $\mathcal{S}^{(\Psi)}$ and $\epsilon^{(\Psi)} := \epsilon$ altogether yield (for all $\ell \in I$)

$$\begin{aligned} \Delta^{(\Psi)}(E_\ell^\Psi) &= E_\ell^\Psi \otimes 1 + e^{+\hbar T_{\ell,+}^\Psi} \otimes E_\ell^\Psi \quad , & \Delta^{(\Psi)}(T_{\ell,+}^\Psi) &= T_{\ell,+}^\Psi \otimes 1 + 1 \otimes T_{\ell,+}^\Psi \\ \mathcal{S}^{(\Psi)}(E_\ell^\Psi) &= -e^{-\hbar T_{\ell,+}^\Psi} E_\ell^\Psi \quad , & \mathcal{S}^{(\Psi)}(T_{\ell,+}^\Psi) &= -T_{\ell,+}^\Psi \\ \epsilon(E_\ell^\Psi) &= 0 \quad , & \epsilon(T_{\ell,\pm}^\Psi) &= 0 \end{aligned}$$

on the generators of $U_\hbar^\Psi(\mathfrak{b}_+)$ and

$$\begin{aligned} \Delta^{(\Psi)}(F_\ell^\Psi) &= F_\ell^\Psi \otimes e^{-\hbar T_{\ell,-}^\Psi} + 1 \otimes F_\ell^\Psi \quad , & \Delta^{(\Psi)}(T_{\ell,-}^\Psi) &= T_{\ell,-}^\Psi \otimes 1 + 1 \otimes T_{\ell,-}^\Psi \\ \mathcal{S}^{(\Psi)}(F_\ell^\Psi) &= -F_\ell^\Psi e^{+\hbar T_{\ell,-}^\Psi} \quad , & \mathcal{S}^{(\Psi)}(T_{\ell,-}^\Psi) &= -T_{\ell,-}^\Psi \\ \epsilon(F_\ell^\Psi) &= 0 \quad , & \epsilon(T_{\ell,-}^\Psi) &= 0 \end{aligned}$$

(for all $\ell \in I$) on those of $U_\hbar^\Psi(\mathfrak{b}_-)$. Altogether, these formulas show that *both* $U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar^\Psi(\mathfrak{b}_-)$ are Hopf subalgebras of $U_\hbar^\Psi(\mathfrak{g}_D)$.

Now, at the semiclassical level, let us consider the elements

$$t_{i,+}^\Psi := (\text{id}_{\mathfrak{h}'_+} + \psi_+^{\mathfrak{h}'})(t_i^+) - \psi_-^{\mathfrak{h}'}(t_i^-) \quad , \quad t_{i,-}^\Psi := (\text{id}_{\mathfrak{h}'_-} + \psi_-^{\mathfrak{h}'})(t_i^-) - \psi_+^{\mathfrak{h}'}(t_i^+) \quad \forall i \in I$$

and the Lie subalgebras

$$\begin{aligned} \mathfrak{b}_+^\Psi &:= \text{Lie subalgebra of } \mathfrak{g}'_D(A) \text{ generated by } \{e_i, t_{i,+}^\Psi\}_{i \in I} \\ \mathfrak{b}_-^\Psi &:= \text{Lie subalgebra of } \mathfrak{g}'_D(A) \text{ generated by } \{t_{i,-}^\Psi, f_i\}_{i \in I} \end{aligned}$$

Notice that the toral part of \mathfrak{b}_\pm^Ψ , that is the Lie subalgebra generated by the new elements $t_{i,\pm}^\Psi$, is isomorphic to \mathfrak{h}' , thanks to Lemma 2.1.4(b).

Then one easily sees that the semiclassical limits of $U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar^\Psi(\mathfrak{b}_-)$ are

$$U_\hbar^\Psi(\mathfrak{b}_+) / \hbar U_\hbar^\Psi(\mathfrak{b}_+) \cong U(\mathfrak{b}_+^\Psi) \quad \text{and} \quad U_\hbar^\Psi(\mathfrak{b}_-) / \hbar U_\hbar^\Psi(\mathfrak{b}_-) \cong U(\mathfrak{b}_-^\Psi)$$

so that the self-explaining notation $U_\hbar(\mathfrak{b}_+^\Psi) := U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar(\mathfrak{b}_-^\Psi) := U_\hbar^\Psi(\mathfrak{b}_-)$ can be also adopted. All this, in turn, also reflects the fact that *both* \mathfrak{b}_+^Ψ and \mathfrak{b}_-^Ψ are Lie sub-bialgebras of the Lie bialgebra $\mathfrak{g}'_D(A)$ (and have the correct size for ‘‘Borel’’!).

Note that, like in §3.2.1 above, from the previous observations we can also find out a complete presentation by generators (the ‘‘twisted’’ ones) and relations for $U_\hbar^\Psi(\mathfrak{b}_+)$ and $U_\hbar^\Psi(\mathfrak{b}_-)$, easily deduced from that for $U_\hbar^\Psi(\mathfrak{g}_D)$: in comparison with those of the corresponding untwisted quantum algebras, in these presentations *the formulas for the coproduct will read the same*, while *the commutation relations will be deformed* instead — just the converse of what occurs in the original presentations.

Remark 3.2.4. Using the skew-Hopf pairing (3.3) between $U_{\hbar}(\mathfrak{b}_+)$ and $U_{\hbar}(\mathfrak{b}_-)$ and the formulas in §2.1.2, one may define a “twisted” skew-Hopf pairing η^Ψ between $U_{\hbar}^\Psi(\mathfrak{b}_+)$ and $U_{\hbar}^\Psi(\mathfrak{b}_-)$, in such a way that the corresponding Drinfeld double $D(U_{\hbar}^\Psi(\mathfrak{b}_+), U_{\hbar}^\Psi(\mathfrak{b}_-), \eta^\Psi)$ is canonically isomorphic to $U_{\hbar}^\Psi(\mathfrak{g}_D)$.

Remark 3.2.5. It is worth noting that, strictly speaking, $U_{\hbar}(\mathfrak{b}_{\pm}^\Psi)$ cannot be thought of as a comultiplication twisting of the isomorphic copy of $U_{\hbar}(\mathfrak{b}_{\pm})$ inside $U_{\hbar}(\mathfrak{g}_D)$; indeed, this occurs because the twisting element

$$\mathcal{F}_\Psi := \exp \left(\hbar \sum_{i,j=1}^n \psi_{ij} H_i^+ \otimes H_j^- \right)$$

that we used to deform $U_{\hbar}(\mathfrak{g}_D)$ — cf. §3.2.2 — does not belong to $U_{\hbar}(\mathfrak{b}_{\pm})$.

On the other hand, when the (generalized) Cartan matrix A is of finite type (which is equivalent to saying that A^{-1} exists!), let us write $(a'_{hk})_{h,k \in I} = A^{-1}$ and consider the elements $T_{\omega_i}^\pm := \sum_{j=1}^n a'_{ji} T_j^\pm$ ($i \in I$). Then by a straightforward calculation we can re-write \mathcal{F}_Ψ as

$$\mathcal{F}_\Psi := \exp \left(\hbar \sum_{k=1}^n d_k^{-1} T_{\omega_k}^+ \otimes \psi_-^{\hbar}(T_k^-) \right) \quad (3.10)$$

and also as

$$\mathcal{F}_\Psi := \exp \left(\hbar \sum_{k=1}^n \psi_+^{\hbar}(T_k^+) \otimes d_k^{-1} T_{\omega_k}^- \right) \quad (3.11)$$

Lead by these formulas, we define $\check{U}_{\hbar}(\mathfrak{b}_+)$ to be the complete, unital subalgebra of $U_{\hbar}^\Psi(\mathfrak{g}_D)$ generated by $\left\{ E_i, d_i^{-1} T_{\omega_i}^+, (\text{id}_{\mathfrak{b}_+} + \psi_+^{\hbar})(T_i^+), \psi_-^{\hbar}(T_i^-) \right\}_{i \in I}$, and $\check{U}_{\hbar}(\mathfrak{b}_-)$ the one generated by $\left\{ d_i^{-1} T_{\omega_i}^-, \psi_+^{\hbar}(T_i^+), (\text{id}_{\mathfrak{b}_+} + \psi_-^{\hbar})(T_i^-), F_i \right\}_{i \in I}$. Then $\check{U}_{\hbar}(\mathfrak{b}_{\pm})$ is a Hopf subalgebra of both $U_{\hbar}(\mathfrak{g}_D)$ and $U_{\hbar}^\Psi(\mathfrak{g}_D)$, hence we shall write $\check{U}_{\hbar}^\Psi(\mathfrak{b}_-)$ when we consider $\check{U}_{\hbar}(\mathfrak{b}_-)$ endowed with the Hopf (sub)algebra structure induced from the (twisted) Hopf algebra structure of $U_{\hbar}(\mathfrak{g}_D)$. Now, by (3.10) and (3.11) we have

$$\mathcal{F}_\Psi \in \check{U}_{\hbar}^\Psi(\mathfrak{b}_+) \widehat{\otimes} \check{U}_{\hbar}^\Psi(\mathfrak{b}_+) \quad \text{and} \quad \mathcal{F}_\Psi \in \check{U}_{\hbar}^\Psi(\mathfrak{b}_-) \widehat{\otimes} \check{U}_{\hbar}^\Psi(\mathfrak{b}_-)$$

so that \mathcal{F}_Ψ can indeed be rightfully called “twisting element” for $\check{U}_{\hbar}(\mathfrak{b}_{\pm})$ — as it does belong to $\check{U}_{\hbar}(\mathfrak{b}_{\pm}) \widehat{\otimes} \check{U}_{\hbar}(\mathfrak{b}_{\pm})$ and twisting $\check{U}_{\hbar}(\mathfrak{b}_{\pm})$ by it we actually get $\check{U}_{\hbar}^\Psi(\mathfrak{b}_{\pm})$.

3.3. Polynomial QUEA’s (“à la Jimbo-Lusztig”).

We shortly recall hereafter the “polynomial version” of the notion of QUEA, as introduced by Jimbo, Lusztig and others, as well as some related material.

Let \mathbb{k}_q be the subfield of $\mathbb{k}[[\hbar]]$ generated by $\mathbb{k} \cup \{ q^{\pm 1/m} := e^{\pm \hbar/m} \mid m \in \mathbb{N}_+ \}$; in particular $q^{\pm 1} := e^{\pm \hbar} \in \mathbb{k}_q$ and $q_i^{\pm 1} := q^{\pm d_i} \in \mathbb{k}_q$ for all $i \in I$. Note that \mathbb{k}_q is the injective limit of all the fields of rational functions $\mathbb{k}(q^{1/m})$, but in specific cases we shall be working with a specific bound on m , fixed from scratch, so in fact we can adopt as ground ring just one such ring $\mathbb{k}(q^{1/N})$ for a single, large enough N .

As a general matter of notation, hereafter by q^r for any $r \in \mathbb{Q}$ we shall always mean $q^r = q^{a/d} := (q^{1/d})^a$ if $r = a/d$ with $a \in \mathbb{Z}$ and $d \in \mathbb{N}_+$.

Remark 3.3.1. As a matter of fact, one could still work over subring(s) of $\mathbb{k}[[\hbar]]$ generated by elements of the set

$$\{q^c := \exp(\hbar c) \mid c \in \mathbb{k}[[\hbar]]\} = \{1 + \hbar c \mid c \in \mathbb{k}[[\hbar]]\}$$

This implies that one considers sublattices in \mathfrak{h}' which might generate different \mathbb{Q} -vector spaces than $\mathfrak{h}'_{\mathbb{Q}}$, but this makes no difference when it comes to *polynomial* QUEA's and their multiparameter versions. As a byproduct, later on we do not need to restrict ourselves to “*rational twists*”, i.e., (toral) twists coming from $\Psi \in M_n(\mathbb{Q})$.

3.3.2. The polynomial QUEA $U_q(\mathfrak{g})$. We introduce the “polynomial” QUEA for $\mathfrak{g}'(A)$, hereafter denoted $U_q(\mathfrak{g}'(A)) = U_q(\mathfrak{g})$, as being the unital \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{g})$ generated by $\{E_i, K_i^{\pm 1} := e^{\pm \hbar d_i H_i}, F_i \mid i \in I\}$. From this definition and from the presentation of $U_{\hbar}(\mathfrak{g})$ we deduce that $U_q(\mathfrak{g})$ can be presented as the associative, unital \mathbb{k}_q -algebra with generators $E_i, K_i^{\pm 1}$ and F_i ($i \in I := \{1, \dots, n\}$) and relations (for all $i, j \in I$)

$$\begin{aligned} K_i K_j &= K_j K_i \quad , \quad K_i^{\pm 1} K_i^{\mp 1} = 1 = K_i^{\mp 1} K_i^{\pm 1} \\ K_i E_j K_i^{-1} &= q_i^{+a_{ij}} E_j \quad , \quad K_i F_j K_i^{-1} = q_i^{-a_{ij}} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{K_i - K_i^{-1}}{q_i - q_i^{-1}} \\ \sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1-a_{ij} \\ k \end{bmatrix}_{q_i} E_i^{1-a_{ij}-k} E_j E_i^k &= 0 \quad (i \neq j) \end{aligned} \quad (3.12)$$

$$\sum_{k=0}^{1-a_{ij}} (-1)^k \begin{bmatrix} 1-a_{ij} \\ k \end{bmatrix}_{q_i} F_i^{1-a_{ij}-k} F_j F_i^k = 0 \quad (i \neq j) \quad (3.13)$$

The formulas for the coproduct, antipode and counit in $U_{\hbar}(\mathfrak{g})$ now give

$$\begin{aligned} \Delta(E_i) &:= E_i \otimes 1 + K_i \otimes E_i \quad , \quad \mathcal{S}(E_i) := -K_i^{-1} E_i \quad , \quad \epsilon(E_i) := 0 \\ \Delta(K_i^{\pm 1}) &:= K_i^{\pm 1} \otimes K_i^{\pm 1} \quad , \quad \mathcal{S}(K_i^{\pm 1}) := K_i^{\mp 1} \quad , \quad \epsilon(K_i^{\pm 1}) := 1 \\ \Delta(F_i) &:= F_i \otimes K_i^{-1} + 1 \otimes F_i \quad , \quad \mathcal{S}(F_i) := -F_i K_i \quad , \quad \epsilon(F_i) := 0 \end{aligned}$$

— for all $i \in I$ — so that $U_q(\mathfrak{g})$ is *actually* a (standard, i.e., non-topological) Hopf subalgebra inside $U_{\hbar}(\mathfrak{g})$; cf. [CP], and references therein for details, taking into account that we adopt slightly different normalizations.

3.3.3. Polynomial quantum Borel (sub)algebras and their double. We consider now the “polynomial version” of the quantum Cartan subalgebra $U_{\hbar}(\mathfrak{h})$, the quantum Borel subalgebras $U_{\hbar}(\mathfrak{b}_{\pm})$ and the “quantum double” $U_{\hbar}(\mathfrak{g}_D)$.

First, we define $U_q(\mathfrak{h})$ as being the unital \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{g})$ generated by $\{K_i^{\pm 1} \mid i \in I\}$; this is clearly a Hopf subalgebra of both $U_{\hbar}(\mathfrak{h})$ and $U_q(\mathfrak{g})$ — in fact, it coincides with $U_{\hbar}(\mathfrak{h}) \cap U_q(\mathfrak{g})$ — isomorphic to the group algebra over \mathbb{k}_q , with its canonical Hopf structure, of the Abelian group \mathbb{Z}^n .

Similarly, we define $U_q(\mathfrak{b}_+)$, resp. $U_q(\mathfrak{b}_-)$, as being the unital \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{g})$ generated by $\{E_i, K_i^{\pm 1} \mid i \in I\}$, resp. by $\{K_i^{\pm 1}, F_i \mid i \in I\}$. Then $U_q(\mathfrak{b}_{\pm})$ is a Hopf subalgebra of both $U_{\hbar}(\mathfrak{b}_{\pm})$ and $U_q(\mathfrak{g})$, coinciding with $U_{\hbar}(\mathfrak{b}_{\pm}) \cap U_q(\mathfrak{g})$.

From the presentation of $U_q(\mathfrak{g})$ one can easily deduce a similar presentation for $U_q(\mathfrak{h})$, $U_q(\mathfrak{b}_+)$ and $U_q(\mathfrak{b}_-)$ as well.

The skew-Hopf pairing $\eta : U_{\hbar}(\mathfrak{b}_+) \otimes_{\mathbb{k}[[\hbar]]} U_{\hbar}(\mathfrak{b}_-) \longrightarrow \mathbb{k}[[\hbar]]$ in §3.1.3 restricts to a similar skew-pairing $\eta : U_q(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_q(\mathfrak{b}_-) \longrightarrow \mathbb{k}_q$ between *polynomial* quantum Borel subalgebras, described by the formulas

$$\eta(K_i, K_j) = q^{-d_i a_{ij}}, \quad \eta(E_i, F_j) = \frac{\delta_{ij}}{q_i^{-1} - q_i}, \quad \eta(E_i, K_j) = 0 = \eta(K_i, F_j) \quad (3.14)$$

Using this pairing one constructs the corresponding Drinfeld double as in §2.3, in the sequel denoted also by $U_q(\mathfrak{g}_D) := D(U_q(\mathfrak{b}_+), U_q(\mathfrak{b}_-), \eta)$.

Tracking the whole construction, one realizes that $U_q(\mathfrak{g}_D)$ naturally identifies with the unital \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{g}_D)$ generated by

$$\{ E_i, K_{i,+} := e^{\pm d_i H_i^+} = e^{\pm T_i^+}, K_{i,-} := e^{\pm d_i H_i^-} = e^{\pm T_i^-}, F_i \mid i \in I \}$$

Then $U_q(\mathfrak{g}_D)$ can be presented as the associative, unital \mathbb{k}_q -algebra with generators $E_i, K_{i,\pm}$ and F_i — with $i \in I$ — satisfying (3.12) and (3.13), together with the relations (for all $i, j \in I$)

$$\begin{aligned} K_{i,+} K_{j,+} &= K_{j,+} K_{i,+}, & K_{i,+} K_{j,-} &= K_{j,-} K_{i,+}, & K_{i,-} K_{j,-} &= K_{j,-} K_{i,-} \\ K_{i,+}^{\pm 1} K_{i,+}^{\mp 1} &= 1 = K_{i,+}^{\mp 1} K_{i,+}^{\pm 1}, & K_{i,-}^{\pm 1} K_{i,-}^{\mp 1} &= 1 = K_{i,-}^{\mp 1} K_{i,-}^{\pm 1} \\ K_{i,\pm} E_j K_{i,\mp}^{-1} &= q^{\pm d_i a_{ij}} E_j, & K_{i,\pm} F_j K_{i,\pm}^{-1} &= q^{-d_i a_{ij}} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{K_{i,+} - K_{i,-}^{-1}}{q_i - q_i^{-1}} \end{aligned}$$

while coproduct, antipode and counit are described by the formulas

$$\begin{aligned} \Delta(E_i) &:= E_i \otimes 1 + K_{i,+} \otimes E_i, & \mathcal{S}(E_i) &:= -K_{i,+}^{-1} E_i, & \epsilon(E_i) &:= 0 \\ \Delta(K_{i,+}^{\pm 1}) &:= K_{i,+}^{\pm 1} \otimes K_{i,+}^{\pm 1}, & \mathcal{S}(K_{i,+}^{\pm 1}) &:= K_{i,+}^{\mp 1}, & \epsilon(K_{i,+}^{\pm 1}) &:= 1 \\ \Delta(K_{i,-}^{\pm 1}) &:= K_{i,-}^{\pm 1} \otimes K_{i,-}^{\pm 1}, & \mathcal{S}(K_{i,-}^{\pm 1}) &:= K_{i,-}^{\mp 1}, & \epsilon(K_{i,-}^{\pm 1}) &:= 1 \\ \Delta(F_i) &:= F_i \otimes K_{i,-}^{-1} + 1 \otimes F_i, & \mathcal{S}(F_i) &:= -F_i K_{i,-}, & \epsilon(F_i) &:= 0 \end{aligned}$$

for all $i \in I$ — so that $U_q(\mathfrak{g}_D)$ is *actually* a Hopf subalgebra inside $U_{\hbar}(\mathfrak{g}_D)$.

In terms of this description, and using the canonical identification

$$U_q(\mathfrak{g}_D) := D(U_q(\mathfrak{b}_+), U_q(\mathfrak{b}_-), \eta) \cong U_q(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_q(\mathfrak{b}_-)$$

as coalgebras, we have $E_i = E_i \otimes 1$, $K_{i,+}^{\pm 1} = K_{i,+}^{\pm 1} \otimes 1$, $K_{i,-}^{\pm 1} = 1 \otimes K_{i,-}^{\pm 1}$, $F_i = 1 \otimes F_i$ (for all $i \in I$). Moreover, the natural embeddings of $U_q(\mathfrak{b}_+)$ and $U_q(\mathfrak{b}_-)$ as Hopf subalgebras inside $U_q(\mathfrak{g}_D) \cong U_q(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_q(\mathfrak{b}_-)$ are described by

$$E_i \mapsto E_i, \quad K_{i,+}^{\pm 1} \mapsto K_{i,+}^{\pm 1} \quad \text{and} \quad F_i \mapsto F_i, \quad K_{i,-}^{\pm 1} \mapsto K_{i,-}^{\pm 1} \quad \forall i \in I$$

Also by construction, the projection $\pi_{\mathfrak{g}} : U_{\hbar}(\mathfrak{g}_D) \longrightarrow U_{\hbar}(\mathfrak{g})$ — an epimorphism of formal Hopf algebras over $\mathbb{k}[[\hbar]]$ — restricts to a Hopf \mathbb{k}_q -algebra epimorphism $\pi_{\mathfrak{g}} : U_q(\mathfrak{g}_D) \longrightarrow U_q(\mathfrak{g})$, which is explicitly described by

$$\pi_{\mathfrak{g}} : \quad E_i \mapsto E_i, \quad K_{i,+}^{\pm 1} \mapsto K_{i,+}^{\pm 1}, \quad K_{i,-}^{\pm 1} \mapsto K_{i,-}^{\pm 1}, \quad F_i \mapsto F_i \quad \forall i \in I \quad (3.15)$$

3.3.4. Larger tori for (polynomial) QUEA's. By definition, the “toral part” of a (polynomial) quantum group $U_q(\mathfrak{g})$ is its Cartan subalgebra $U_q(\mathfrak{h})$: the latter identifies with the group (Hopf) algebra of the Abelian group \mathbb{Z}^n , which in turn is identified with the free Abelian group \mathcal{K}_Q generated by the K_i 's via $K_i \mapsto e_i$ (the i -th element in the canonical \mathbb{Z} -basis of \mathbb{Z}^n). As \mathbb{Z}^n is isomorphic to the root lattice

$Q := \sum_{i \in I} \mathbb{Z} \alpha_i$ via $e_i \mapsto \alpha_i$, we have also an isomorphism $\Omega_Q : Q \xrightarrow{\cong} \mathcal{K}_Q$ given by $\alpha_i \mapsto K_i$; we then fix notation $K_\alpha := \Omega_Q(\alpha)$ ($\in \mathcal{K}_Q$) for every $\alpha \in Q$.

With this notation in use, note that the commutation relations between the $K_i^{\pm 1}$'s and the E_j 's or the F_j 's generalize to

$$K_\alpha E_j K_\alpha^{-1} = q^{+(\alpha, \alpha_j)} E_j \quad , \quad K_\alpha F_j K_\alpha^{-1} = q^{-(\alpha, \alpha_j)} F_j \quad \forall \alpha \in Q, j \in I$$

where (\cdot, \cdot) is the symmetric bilinear pairing on $\mathbb{Q}Q \times \mathbb{Q}Q$ introduced in §2.1.1. Note that $\pm(\alpha, \alpha_j) \in \mathbb{Z}$ so that $q^{\pm(\alpha, \alpha_j)} \in \mathbb{k}[q, q^{-1}] \subseteq \mathbb{k}_q$: in particular, $U_q(\mathfrak{g})$ is *actually defined over the smaller ring $\mathbb{k}[q, q^{-1}]$ too* — no need of the whole \mathbb{k}_q .

Let Γ be a sublattice of $\mathbb{Q}Q$ of rank n with $Q \subseteq \Gamma$. For any basis $\{\gamma_1, \dots, \gamma_n\}$ of Γ , let $C := (c_{ij})_{i=1, \dots, n; j=1, \dots, n}$ be the matrix of integers such that $\alpha_i = \sum_{j=1}^n c_{ij} \gamma_j$ for every $i \in I = \{1, \dots, n\}$. Write $c := |\det(C)| \in \mathbb{N}_+$: as it is known, it equals the index (as a subgroup) of Q in Γ ; in particular, it is independent of any choice of bases. Write $C^{-1} = (c'_{ij})_{i=1, \dots, n; j=1, \dots, n}$ for the inverse matrix to C : then $\gamma_i = \sum_{j=1}^n c'_{ij} \alpha_j$ for each $i \in I = \{1, \dots, n\}$, where $c'_{ij} \in c^{-1} \mathbb{Z}$, for all $i, j \in I$.

We denote by $U_{q, \Gamma}(\mathfrak{h})$ the group algebra over \mathbb{k}_q of the lattice Γ , with its canonical structure of Hopf algebra. So, each element $\gamma \in \Gamma$ corresponds to an element $K_\gamma \in U_{q, \Gamma}(\mathfrak{h})$. Denote by \mathcal{K}_Γ the subgroup of $U_{q, \Gamma}(\mathfrak{h})$ generated by all these K_γ 's.

This clearly yields a group isomorphism $\Omega_\Gamma : \Gamma \xrightarrow{\cong} \mathcal{K}_\Gamma$ given by $\gamma \mapsto K_\gamma$ that extends the Ω_Q given above. Moreover, as the group Q embeds into Γ we have a corresponding Hopf algebra embedding $U_q(\mathfrak{h}) = U_{q, Q}(\mathfrak{h}) \hookrightarrow U_{q, \Gamma}(\mathfrak{h})$. In particular, each $K_i = K_{\alpha_i}$ is expressed by the formula $K_{\alpha_i} = \prod_{j \in I} K_{\gamma_j}^{c_{ij}}$.

Note that each of these extended quantum Cartan (sub)algebras still embeds, in a natural way, inside $U_h(\mathfrak{h})$. To see it, we use again notation from §2.1.1: consider the \mathbb{Q} -span of the H_i 's as a \mathbb{Q} -integral form of \mathfrak{h} , take the associated isomorphism $t : \mathfrak{h}_\mathbb{Q}^* \xrightarrow{\cong} \mathfrak{h}_\mathbb{Q}$ and look at the elements $T_i := t_{\alpha_i} = d_i H_i$ for all $i \in I$. By construction, we have

$$K_{\alpha_i} \equiv K_i := e^{h d_i H_i} = e^{h T_i} \quad \forall i \in I$$

and more in general for $\alpha \in Q$ written as $\alpha = \sum_{j \in I} z_j \alpha_j$ (with $z_j \in \mathbb{Z}$) we have

$$K_\alpha := \prod_{j \in I} K_j^{z_j} = \prod_{j \in I} e^{h z_j d_j H_j} = \prod_{j \in I} e^{h z_j T_j} = e^{h \sum_{j \in I} z_j T_j} = e^{h T_\alpha} \quad \forall \alpha \in Q$$

Now this picture extends to any other lattice Γ in $\mathfrak{h}_\mathbb{Q}^* \equiv \mathbb{Q}Q$ containing Q . Indeed, mapping $K_\gamma \mapsto e^{h T_\gamma}$ — for all $\gamma \in \Gamma$ — provides a unique, well-defined monomorphism of Hopf algebras from $U_{q, \Gamma}(\mathfrak{h})$ to $U_h(\mathfrak{h})$. In other words, $U_{q, \Gamma}(\mathfrak{h})$ *canonically identifies with the \mathbb{k}_q -subalgebra of $U_h(\mathfrak{h})$ generated by $\{K_\gamma := e^{h T_\gamma} \mid \gamma \in \Gamma\}$.*

3.3.5. Polynomial QUEA's with larger torus. We aim now to introduce new polynomial QUEA's having a “larger Cartan subalgebra”, modeled on those of the form $U_{q, \Gamma}(\mathfrak{h})$ presented in §3.3.4 above.

To begin with, let Γ be any lattice in $\mathbb{Q}Q$ such that $Q \subseteq \Gamma$. Like we did for $U_q(\mathfrak{g})$, we define $U_{q, \Gamma}(\mathfrak{g})$ as being the unital \mathbb{k}_q -subalgebra of $U_h(\mathfrak{g})$ generated by $\{E_i, K_\gamma := e^{h T_\gamma}, F_i \mid i \in I, \gamma \in \Gamma\}$. From this definition and from the presentation of $U_h(\mathfrak{g})$ we deduce that $U_{q, \Gamma}(\mathfrak{g})$ can be presented as the associative, unital \mathbb{k}_q -algebra with generators E_i, K_γ and F_i ($i \in I := \{1, \dots, n\}, \gamma \in \Gamma$) satisfying

(3.12) and (3.13), together with the relations

$$\begin{aligned} K_{\gamma'} K_{\gamma''} &= K_{\gamma'+\gamma''} = K_{\gamma''} K_{\gamma'} \quad , & K_{+\gamma} K_{-\gamma} &= 1 = K_{-\gamma} K_{+\gamma} \\ K_{\gamma} E_j K_{\gamma}^{-1} &= q^{+(\gamma, \alpha_j)} E_j \quad , & K_{\gamma} F_j K_{\gamma}^{-1} &= q^{-(\gamma, \alpha_j)} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{K_{\alpha_i} - K_{\alpha_i}^{-1}}{q_i - q_i^{-1}} \end{aligned}$$

(for all $i, j \in I$, $\gamma, \gamma', \gamma'' \in \Gamma$), where $q_i := q^{d_i} = e^{h d_i}$ as before. The formulas for the coproduct, antipode and counit in $U_{\hbar}(\mathfrak{g})$ then give

$$\begin{aligned} \Delta(E_i) &= E_i \otimes 1 + K_{\alpha_i} \otimes E_i \quad , & \mathcal{S}(E_i) &= -K_{\alpha_i}^{-1} E_i \quad , & \epsilon(E_i) &= 0 \\ \Delta(K_{\gamma}^{\pm 1}) &= K_{\gamma}^{\pm 1} \otimes K_{\gamma}^{\pm 1} \quad , & \mathcal{S}(K_{\gamma}^{\pm 1}) &= K_{\gamma}^{\mp 1} \quad , & \epsilon(K_{\gamma}^{\pm 1}) &= 1 \\ \Delta(F_i) &= F_i \otimes K_{\alpha_i}^{-1} + 1 \otimes F_i \quad , & \mathcal{S}(F_i) &= -F_i K_{\alpha_i} \quad , & \epsilon(F_i) &= 0 \end{aligned}$$

(for all $i \in I$, $\gamma \in \Gamma$) so that $U_{q, \Gamma}(\mathfrak{g})$ is (again) a Hopf subalgebra inside $U_{\hbar}(\mathfrak{g})$.

With notation of §3.3.4, let $C := (c_{ij})_{i=1, \dots, n; j=1, \dots, n}$ be the matrix of change of \mathbb{Q} -basis (for $\mathbb{Q}Q$) from any basis of Γ to the basis $\{\alpha_1, \dots, \alpha_n\}$ of simple roots and set $c := |\det(C)| \in \mathbb{N}_+$. Then from the above presentation one easily sees that $U_{q, \Gamma}(\mathfrak{g})$ is actually well-defined over the subfield $\mathbb{k}(q^{1/c})$ of \mathbb{k}_q .

With much the same method we define also quantum Borel subalgebras with larger torus, modeled on the lattice Γ , hereafter denoted $U_{q, \Gamma}(\mathfrak{b}_+)$, resp. $U_{q, \Gamma}(\mathfrak{b}_-)$, simply dropping the F_i 's, resp. the E_i 's, from the set of generators.

Similarly, if we take as generators only the K_{γ} 's we get a Hopf \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{h})$ (hence of $U_{\hbar}(\mathfrak{g})$ as well) that is isomorphic to $U_{q, \Gamma}(\mathfrak{h})$ — cf. §3.3.4 above.

Note also that definitions give $U_q(\mathfrak{b}_{\pm}) = U_{q, \mathbb{Q}}(\mathfrak{b}_{\pm}) \subseteq U_{q, \Gamma}(\mathfrak{b}_{\pm}) \subseteq U_{\hbar}(\mathfrak{b}_{\pm})$. Moreover, all these algebraic objects are actually well-defined over $\mathbb{k}(q^{1/c})$ as well.

Now let Γ_+ and Γ_- be any two lattices in $\mathbb{Q}Q$ such that $Q \subseteq \Gamma_{\pm}$, and let $\Gamma_{\bullet} := \Gamma_+ \times \Gamma_-$. Then we define $U_{q, \Gamma_{\bullet}}(\mathfrak{g}_D)$ as being the unital \mathbb{k}_q -subalgebra of $U_{\hbar}(\mathfrak{g}_D)$ generated by $\{E_i, K_{\gamma_{\pm}} := e^{h T_{\gamma_{\pm}}}, F_i \mid i \in I, \gamma_{\pm} \in \Gamma_{\pm}\}$. From the presentation of $U_{\hbar}(\mathfrak{g}_D)$ we deduce that $U_{q, \Gamma_{\bullet}}(\mathfrak{g})$ can be presented as the associative, unital \mathbb{k}_q -algebra with generators $E_i, K_{\gamma_{\pm}}$ and F_i ($i \in I := \{1, \dots, n\}$, $\gamma_{\pm} \in \Gamma_{\pm}$) satisfying the relations (3.12) and (3.13), together with the following:

$$\begin{aligned} K_{\gamma_+} K_{\gamma_-} &= K_{\gamma_-} K_{\gamma_+} \\ K_{\gamma'_{\pm}} K_{\gamma''_{\pm}} &= K_{\gamma'_{\pm} + \gamma''_{\pm}} = K_{\gamma''_{\pm}} K_{\gamma'_{\pm}} \quad , & K_{+\gamma_{\pm}} K_{-\gamma_{\pm}} &= 1 = K_{-\gamma_{\pm}} K_{+\gamma_{\pm}} \\ K_{\gamma_{\pm}} E_j K_{\gamma_{\pm}}^{-1} &= q^{+(\gamma_{\pm}, \alpha_j)} E_j \quad , & K_{\gamma_{\pm}} F_j K_{\gamma_{\pm}}^{-1} &= q^{-(\gamma_{\pm}, \alpha_j)} F_j \\ E_i F_j - F_j E_i &= \delta_{ij} \frac{K_{\alpha_i^+} - K_{\alpha_i^-}^{-1}}{q_i - q_i^{-1}} \end{aligned}$$

(for all $i, j \in I$, $\gamma_{\pm}, \gamma'_{\pm}, \gamma''_{\pm} \in \Gamma_{\pm}$), where α_i^{\pm} is the copy of $\alpha_i (\in Q)$ inside Γ_{\pm} and $q_i := q^{d_i} = e^{h d_i}$ as usual. Moreover, coproduct, antipode and counit of $U_{\hbar}(\mathfrak{g}_D)$ on these generators read

$$\begin{aligned} \Delta(E_i) &= E_i \otimes 1 + K_{\alpha_i^+} \otimes E_i \quad , & \mathcal{S}(E_i) &= -K_{\alpha_i^+}^{-1} E_i \quad , & \epsilon(E_i) &= 0 \\ \Delta(K_{\gamma_{\pm}}^{\pm 1}) &= K_{\gamma_{\pm}}^{\pm 1} \otimes K_{\gamma_{\pm}}^{\pm 1} \quad , & \mathcal{S}(K_{\gamma_{\pm}}^{\pm 1}) &= K_{\gamma_{\pm}}^{\mp 1} \quad , & \epsilon(K_{\gamma_{\pm}}^{\pm 1}) &= 1 \\ \Delta(F_i) &= F_i \otimes K_{\alpha_i^-}^{-1} + 1 \otimes F_i \quad , & \mathcal{S}(F_i) &= -F_i K_{\alpha_i^-} \quad , & \epsilon(F_i) &= 0 \end{aligned}$$

(for all $i \in I$, $\gamma_{\pm} \in \Gamma_{\pm}$). Thus $U_{q, \Gamma_{\bullet}}(\mathfrak{g}_D)$ is also a Hopf subalgebra inside $U_{\hbar}(\mathfrak{g}_D)$.

Fix again two lattices Γ_+ and Γ_- in $\mathbb{Q}\mathbb{Q}$ such that $Q \subseteq \Gamma_\pm$. As $U_{q,\Gamma_\pm}(\mathfrak{b}_\pm) \subseteq U_{\hbar}(\mathfrak{b}_\pm)$, the skew-pairing $\eta : U_{\hbar}(\mathfrak{b}_+) \otimes_{\mathbb{k}[[\hbar]]} U_{\hbar}(\mathfrak{b}_-) \longrightarrow \mathbb{k}[[\hbar]]$ restricts to a similar skew-Hopf pairing $\eta : U_{q,\Gamma_+}(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_{q,\Gamma_-}(\mathfrak{b}_-) \longrightarrow \mathbb{k}_q$, which in turn extends the one on $U_q(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_q(\mathfrak{b}_-) = U_{q,Q}(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_{q,Q}(\mathfrak{b}_-)$; the formulas

$$\eta(K_{\gamma_+}, K_{\gamma_-}) = q^{-(\gamma_+, \gamma_-)} , \quad \eta(E_i, F_j) = \frac{\delta_{ij}}{q_i^{-1} - q_i} , \quad \eta(E_i, K_j) = 0 = \eta(K_i, F_j)$$

uniquely determines this pairing. Using the latter, we construct the corresponding Drinfeld double $D(U_{q,\Gamma_+}(\mathfrak{b}_+), U_{q,\Gamma_-}(\mathfrak{b}_-), \eta)$ as in §2.3. It follows by construction that this double coincides with the Hopf algebra $U_{q,\Gamma_\bullet}(\mathfrak{g}_D)$ considered right above.

Again by construction, we also have that the projection $\pi_{\mathfrak{g}} : U_{\hbar}(\mathfrak{g}_D) \longrightarrow U_{\hbar}(\mathfrak{g})$ yields by restriction a Hopf \mathbb{k}_q -algebra epimorphism $\pi_{\mathfrak{g}} : U_{q,\Gamma_\bullet}(\mathfrak{g}_D) \longrightarrow U_{q,\Gamma_*}(\mathfrak{g})$, where $\Gamma_* := \Gamma_+ + \Gamma_- (\subseteq \mathbb{Q}\mathbb{Q})$ whose explicit description is obvious.

An alternative method to construct these QUEA's with larger toral part goes as follows. Fix a lattice Γ in $\mathbb{Q}\mathbb{Q}$ such that $Q \subseteq \Gamma$, and consider the associated Hopf algebra $U_{q,\Gamma}(\mathfrak{h})$ as in §3.3.4 and the canonical embedding $U_q(\mathfrak{h}) = U_{q,Q}(\mathfrak{h}) \subseteq U_{q,\Gamma}(\mathfrak{h})$. The natural (adjoint) action of $U_q(\mathfrak{h})$ onto $U_q(\mathfrak{g})$ extends (uniquely) to a $U_{q,\Gamma}(\mathfrak{h})$ -action $\cdot : U_{q,\Gamma}(\mathfrak{h}) \times U_q(\mathfrak{g}) \longrightarrow U_q(\mathfrak{g})$ given (for all $i \in I$ and $\gamma \in \Gamma$) by

$$K_\gamma \cdot E_j = q^{+(\gamma, \alpha_j)} E_j , \quad K_\gamma \cdot F_j = q^{-(\gamma, \alpha_j)} F_j , \quad K_\gamma \cdot K_j = K_j$$

that makes $U_q(\mathfrak{g})$ into a $U_{q,\Gamma}(\mathfrak{h})$ -module algebra. Then we can consider the Hopf algebra $U_{q,\Gamma}(\mathfrak{h}) \times U_q(\mathfrak{g})$ given by the *smash product* of $U_{q,\Gamma}(\mathfrak{h})$ and $U_q(\mathfrak{g})$: the underlying vector space is just $U_{q,\Gamma}(\mathfrak{h}) \otimes_{\mathbb{k}_q} U_q(\mathfrak{g})$, the coalgebra structure is the one given by the tensor product of the corresponding coalgebras, and the product is given by the formula

$$(h \times x)(k \times y) = h k_{(1)} \times (\mathcal{S}(k_{(2)}) \cdot x) y \quad (3.16)$$

for all $h, k \in U_{q,\Gamma}(\mathfrak{h})$, $x, y \in U_q(\mathfrak{g})$. Since $U_{q,\Gamma}(\mathfrak{h})$ contains $U_q(\mathfrak{h})$ as a Hopf subalgebra, it follows that $U_{q,\Gamma}(\mathfrak{h})$ itself is a right $U_q(\mathfrak{h})$ -module Hopf algebra with respect to the adjoint action. Under these conditions, it is easy to see that the smash product $U_{q,\Gamma}(\mathfrak{h}) \times U_q(\mathfrak{g})$ maps onto a Hopf algebra structure on the \mathbb{k}_q -module $U_{q,\Gamma}(\mathfrak{h}) \otimes_{U_q(\mathfrak{h})} U_q(\mathfrak{g})$, which we can denote by $U_{q,\Gamma}(\mathfrak{h}) \times_{U_q(\mathfrak{h})} U_q(\mathfrak{g})$, see [Len, Theorem 2.8]. Finally, tracking the whole construction one easily sees that this Hopf algebra $U_{q,\Gamma}(\mathfrak{h}) \times_{U_q(\mathfrak{h})} U_q(\mathfrak{g})$ actually coincides with the Hopf algebra $U_{q,\Gamma}(\mathfrak{g})$ considered above.

With the same approach, one can also realize $U_{q,\Gamma_\pm}(\mathfrak{b}_\pm)$, resp. $U_{q,\Gamma}(\mathfrak{b}_-)$, as Hopf algebra structure on $U_{q,\Gamma_\pm}(\mathfrak{h}) \otimes_{U_q(\mathfrak{h})} U_q(\mathfrak{b}_\pm)$, resp. $U_{q,\Gamma_\bullet}(\mathfrak{h}) \otimes_{U_q(\mathfrak{h}) \otimes U_q(\mathfrak{h})} U_q(\mathfrak{g}_D)$, obtained as quotient of the smash product Hopf algebra $U_{q,\Gamma_\pm}(\mathfrak{h}) \times U_q(\mathfrak{b}_\pm)$, resp. $U_{q,\Gamma_\bullet}(\mathfrak{h}) \times U_q(\mathfrak{g}_D)$.

3.4. Comultiplication twistings of polynomial QUEA's (=TwQUEA's).

We introduce now the polynomial version of twisted QUEA's, or twisted polynomial QUEA's, just by matching what we did in §3.3 and §3.2.

3.4.1. Comultiplication twistings of $U_q(\mathfrak{g})$. Let $\mathbb{k}_q := \varprojlim_{m \in \mathbb{N}} \mathbb{k}(q^{1/m}) (\subseteq \mathbb{k}[[\hbar]])$ and $U_q(\mathfrak{g})$ be as in §3.3. Fix also an $(n \times n)$ -matrix $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{k}[[\hbar]])$ as in §3.2.1, but now we make the stronger assumption that $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{Q})$.

Via the recipe in §3.2.1, we pick the corresponding twisted formal QUEA $U_h^\Psi(\mathfrak{g})$. Then inside the latter we consider the unital \mathbb{k}_q -subalgebra $U_q(\mathfrak{g})$, generated by $\{E_i, K_i^{\pm 1} := e^{\pm h d_i H_i} = e^{\pm h T_i}, F_i \mid i \in I\}$, for which we have an explicit presentation. On the other hand, the new, twisted Hopf structure of $U_h^\Psi(\mathfrak{g})$ on the generators of $U_q(\mathfrak{g})$ — with notation of §2.1.2 and §3.3.4 — reads

$$\begin{aligned} \Delta^{(\Psi)}(E_\ell) &= E_\ell \otimes K_{+\psi_-(\alpha_\ell)} + K_{+(\text{id}+\psi_+)(\alpha_\ell)} \otimes E_\ell = E_\ell \otimes K_{+\zeta_\ell^-} + K_{+\alpha_\ell+\zeta_\ell^+} \otimes E_\ell \\ \Delta^{(\Psi)}(K_\ell^{\pm 1}) &= K_\ell^{\pm 1} \otimes K_\ell^{\pm 1} \\ \Delta^{(\Psi)}(F_\ell) &= F_\ell \otimes K_{-(\text{id}+\psi_-)(\alpha_\ell)} + K_{-\psi_+(\alpha_\ell)} \otimes F_\ell = F_\ell \otimes K_{-\alpha_\ell-\zeta_\ell^-} + K_{-\zeta_\ell^+} \otimes F_\ell \\ \mathcal{S}^{(\Psi)}(E_\ell) &= -K_{-(\text{id}+\psi_+)(\alpha_\ell)} E_\ell K_{-\psi_-(\alpha_\ell)} = -K_{-\alpha_\ell-\zeta_\ell^+} E_\ell K_{-\zeta_\ell^-}, \quad \epsilon(E_\ell) = 0 \\ \mathcal{S}^{(\Psi)}(K_\ell^{\pm 1}) &= K_\ell^{\mp 1}, \quad \epsilon(K_\ell^{\pm 1}) = 1 \\ \mathcal{S}^{(\Psi)}(F_\ell) &= -K_{+\psi_+(\alpha_\ell)} F_\ell K_{+(\text{id}+\psi_-)(\alpha_\ell)} = -K_{+\zeta_\ell^+} F_\ell K_{+\alpha_\ell+\zeta_\ell^-}, \quad \epsilon(F_\ell) = 0 \end{aligned}$$

(for all $\ell \in I$). This shows explicitly that

$U_q(\mathfrak{g})$ is a Hopf subalgebra (over \mathbb{k}_q) inside $U_h^\Psi(\mathfrak{g})$ if and only if $\psi_\pm(Q) \subseteq Q$.

In order to settle this point, we pick the sublattice in $\mathbb{Q}\mathbb{Q}$ given by $Q^\Psi := Q + \psi_+(Q) + \psi_-(Q)$, and for each lattice Γ in $\mathbb{Q}\mathbb{Q}$ containing Q^Ψ we consider the corresponding polynomial QUEA, namely $U_{q,\Gamma}(\mathfrak{g})$ as in §3.3.5. Then $U_{q,\Gamma}(\mathfrak{g})$ is naturally embedded inside $U_h(\mathfrak{g})$, and repeating the previous analysis we see that

$$U_{q,\Gamma}(\mathfrak{g}) \text{ is a Hopf subalgebra (over } \mathbb{k}_q \text{) inside } U_h^\Psi(\mathfrak{g}). \quad (3.17)$$

To sum up, the previous analysis allows us to give the following definition:

Definition 3.4.2. We denote by $U_{q,\Gamma}^\Psi(\mathfrak{g})$ the Hopf algebra defined in (3.17), whose Hopf structure is given by restriction from $U_h^\Psi(\mathfrak{g})$. We call any such Hopf algebra *twisted polynomial QUEA* — or *TwQUEA*, in short — saying it is obtained from $U_{q,\Gamma}(\mathfrak{g})$ by twisting (although, strictly speaking, this is not entirely correct).

Remark 3.4.3. The multiparameter quantum groups $U_q^\varphi(\mathfrak{g})$ introduced by Costantini and Varagnolo in [CV1, CV2, CV3] are a particular case of a twisted polynomial QUEA, obtained by taking $2\psi_- = -2\psi_+ = \varphi$. More precisely, they fix assumptions on φ — hence on ψ — that guarantee that it is enough to take, once and for all, the “larger torus” modeled on the lattice $\Gamma = P$.

3.4.4. Comultiplication twistings of $U_q(\mathfrak{g}_D)$. Let again $\mathbb{k}_q := \varprojlim_{m \in \mathbb{N}} \mathbb{k}(q^{1/m})$ and

$\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{Q})$ be as in §3.4.1. Following §3.2.2, we pick the twisted formal QUEA $U_h^\Psi(\mathfrak{g}_D)$: inside the latter, we pick the unital \mathbb{k}_q -subalgebra $U_q(\mathfrak{g}_D)$ of §3.3.3, for which we know an explicit presentation.

Since we are working with a double copy of \mathfrak{h}' , hence also with $\mathbb{Q}(Q \times Q) = \mathbb{Q}\mathbb{Q} \times \mathbb{Q}\mathbb{Q}$, we write Q_\pm to mean $Q_+ := Q \times \{0\}$ and $Q_- := \{0\} \times Q$, and also — for every $\gamma \in \mathbb{Q}\mathbb{Q}$ — we set $\gamma_+ := (\gamma, 0)$ and $\gamma_- := (0, \gamma)$ inside $\mathbb{Q}\mathbb{Q} \times \mathbb{Q}\mathbb{Q}$. Besides this, hereafter the notation is that of §2.1.2 and §3.3.4.

The twisted Hopf structure of $U_h^\Psi(\mathfrak{g}_D)$ on the generators of $U_q(\mathfrak{g}_D)$ yields

$$\begin{aligned}\Delta^{(\Psi)}(E_\ell) &= E_\ell \otimes K_{+\psi_-(\alpha_\ell^-)} + K_{+(\text{id}+\psi_+)(\alpha_\ell^+)} \otimes E_\ell = E_\ell \otimes K_{+\zeta_{\ell,-}^-} + K_{+\alpha_\ell^+ + \zeta_{\ell,+}^+} \otimes E_\ell \\ \Delta^{(\Psi)}(K_\ell^{\pm 1}) &= K_\ell^{\pm 1} \otimes K_\ell^{\pm 1}\end{aligned}$$

$$\Delta^{(\Psi)}(F_\ell) = F_\ell \otimes K_{-(\text{id}+\psi_-)(\alpha_\ell^-)} + K_{-\psi_+(\alpha_\ell^+)} \otimes F_\ell = F_\ell \otimes K_{-\alpha_\ell^- - \zeta_{\ell,-}^-} + K_{-\zeta_{\ell,+}^+} \otimes F_\ell$$

$$\begin{aligned}\mathcal{S}^{(\Psi)}(E_\ell) &= -K_{-(\text{id}+\psi_+)(\alpha_\ell^+)} E_\ell K_{-\psi_-(\alpha_\ell^-)} = -K_{-\alpha_\ell^+ - \zeta_{\ell,+}^+} E_\ell K_{-\zeta_{\ell,-}^-}, & \epsilon(E_\ell) &= 0 \\ \mathcal{S}^{(\Psi)}(K_\ell^{\pm 1}) &= K_\ell^{\mp 1}, & \epsilon(K_\ell^{\pm 1}) &= 1\end{aligned}$$

$$\mathcal{S}^{(\Psi)}(F_\ell) = -K_{+\psi_+(\alpha_\ell^+)} F_\ell K_{+(\text{id}+\psi_-)(\alpha_\ell^-)} = -K_{+\zeta_{\ell,+}^+} F_\ell K_{+\alpha_\ell^- + \zeta_{\ell,-}^-}, \quad \epsilon(F_\ell) = 0$$

for all $\ell \in I$. From the formulas above we see that

$U_q(\mathfrak{g}_D)$ is a Hopf subalgebra (over \mathbb{k}_q) in $U_h^\Psi(\mathfrak{g}_D)$ if and only if $\psi_\pm(Q_\pm) \subseteq Q_\pm$.

To fix this issue, we consider the sublattices $Q_{(\pm)}^\Psi := (\text{id} + \psi_\pm)(Q_\pm) + \psi_\mp(Q_\mp)$ inside $\mathbb{Q}\mathbb{Q} \times \mathbb{Q}\mathbb{Q}$ and their sum $Q_*^\Psi := Q_{(+)}^\Psi + Q_{(-)}^\Psi$; then for each lattice Γ_\circ inside $\mathbb{Q}\mathbb{Q} \times \mathbb{Q}\mathbb{Q}$ containing Q_*^Ψ we consider the associated polynomial QUEA $U_{q,\Gamma_\circ}(\mathfrak{g}_D)$ as in §3.3.5. By construction $U_{q,\Gamma_\circ}(\mathfrak{g}_D)$ sits inside $U_h(\mathfrak{g}_D)$, and repeating the previous analysis we find that

$$U_{q,\Gamma_\circ}(\mathfrak{g}_D) \text{ is a Hopf subalgebra (over } \mathbb{k}_q) \text{ inside } U_h^\Psi(\mathfrak{g}_D). \quad (3.18)$$

At the end of the day, we are allowed to give the following definition:

Definition 3.4.5. We denote by $U_{q,\Gamma_\circ}^\Psi(\mathfrak{g}_D)$ the Hopf algebra given in (3.18), whose Hopf structure is given by restriction from $U_h^\Psi(\mathfrak{g}_D)$. We call any such Hopf algebra *twisted polynomial QUEA* — or *TwQUEA*, in short — saying it is obtained from $U_{q,\Gamma_\circ}(\mathfrak{g}_D)$ by twisting (although, strictly speaking, this is not entirely correct).

Remark 3.4.6. It follows by construction that the epimorphism of twisted formal QUEA's $\pi_\mathfrak{g}^\Psi := \pi_\mathfrak{g} : U_h^\Psi(\mathfrak{g}_D) \longrightarrow U_h^\Psi(\mathfrak{g})$ (cf. §3.2.2) restricts to an epimorphism $\pi_\mathfrak{g}^\Psi := \pi_\mathfrak{g} : U_{q,\Gamma_\circ}^\Psi(\mathfrak{g}_D) \longrightarrow U_{q,\Gamma}^\Psi(\mathfrak{g})$ of twisted polynomial QUEA's, where Γ is the image of Γ_\circ for the projection of $\mathbb{Q}(Q \times Q)$ onto $\mathbb{Q}\mathbb{Q}$ mapping $\alpha_\pm \in Q_\pm$ onto $\alpha \in Q$.

3.4.7. Comultiplication twistings of $U_q(\mathfrak{b}_\pm)$. We still work with $\Psi := (\psi_{ij})_{i,j \in I} \in M_n(\mathbb{Q})$ as in §3.4.1. Using it, we define suitable “comultiplication twistings” of the quantum Borel algebras $U_q(\mathfrak{b}_\pm)$ as *Hopf subalgebras inside $U_q^\Psi(\mathfrak{g})$ and inside $U_q^\Psi(\mathfrak{g}_D)$* .

First we pick in $\mathbb{Q}\mathbb{Q}$ the sublattices $Q_\pm^\Psi := (\text{id} + \psi_\pm)(Q) + \psi_\mp(Q)$, and for any lattice Γ_\pm in $\mathbb{Q}\mathbb{Q}$ containing Q_\pm^Ψ we consider $\Gamma_* := \Gamma_+ + \Gamma_-$ and the corresponding $U_{q,\Gamma_*}(\mathfrak{g})$, as in §3.4.1; inside the latter, we consider the (polynomial) quantum Borel (or “Borel-like”) subalgebra $U_{q,\Gamma_\pm}(\mathfrak{b}_\pm)$. Now, the formulas for the twisted Hopf structure of $U_{q,\Gamma_*}(\mathfrak{g})$ show that the subalgebra $U_{q,\Gamma_\pm}(\mathfrak{b}_\pm)$ is also a *Hopf subalgebra* in $U_{q,\Gamma_*}^\Psi(\mathfrak{g})$. Therefore $U_{q,\Gamma_\pm}(\mathfrak{b}_\pm)$ with the twisted coproduct, antipode and counit is a new Hopf algebra, that we denote by $U_{q,\Gamma_\pm}^\Psi(\mathfrak{b}_\pm)$, and call (polynomial) “twisted” quantum Borel subalgebra.

Note also that both $U_{q,\Gamma_+}^\Psi(\mathfrak{b}_+)$ and $U_{q,\Gamma_-}^\Psi(\mathfrak{b}_-)$ are Hopf subalgebras in $U_{q,\Gamma_*}^\Psi(\mathfrak{g})$.

Second, we pick in $\mathbb{Q}(Q \times Q)$ the sublattices $Q_{(\pm)}^\Psi := (\text{id} + \psi_\pm)(Q_\pm) + \psi_\mp(Q_\mp)$, and for any lattice $\Gamma_{(\pm)}$ in $\mathbb{Q}(Q \times Q)$ containing $Q_{(\pm)}^\Psi$ we consider $\Gamma_{(*)} := \Gamma_{(+)} + \Gamma_{(-)}$

and the corresponding (double) quantum group $U_{q,\Gamma(*)}^\Psi(\mathfrak{g}_D)$; inside the latter, we fix the (polynomial) quantum Borel subalgebra $U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm)$. Then the explicit formulas for the (twisted) Hopf structure of $U_{q,\Gamma(*)}^\Psi(\mathfrak{g}_D)$ show that

$$U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm) \text{ is a Hopf subalgebra of } U_{q,\Gamma(*)}^\Psi(\mathfrak{g}_D).$$

However, a major drawback of both the subalgebras $U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm)$ — inside $U_{q,\Gamma(*)}^\Psi(\mathfrak{g})$ — and $U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm)$ — inside $U_{q,\Gamma(*)}^\Psi(\mathfrak{g}_D)$ — is that *they have too large a toral part* to be rightfully called “(polynomial) quantum Borel (sub)algebras”. We tackle and settle this problem in §3.4.8 hereafter.

3.4.8. Twisted generators for polynomial twisted QUEA’s. Let us consider in $\mathbb{Q}(Q \times Q)$ the elements $\tau_i^\pm := (\text{id} + \psi_\pm)(\alpha_i^\pm) - \psi_\mp(\alpha_i^\mp) = \alpha_i^\pm + \zeta_{i,\pm}^\pm - \zeta_{i,\mp}^\mp$ ($i \in I$) — cf. §2.1.2 — and the sublattices $Q_{[\pm]}^\Psi$ with \mathbb{Z} -basis $\{\tau_i^\pm \mid i \in I\}$.

Let $Q_{(\pm)}^\Psi := (\text{id} + \psi_\pm)(Q_\pm) + \psi_\mp(Q_\mp)$ and $\Gamma_{(\pm)}$ be as in §3.4.7 above; then $Q_{[\pm]}^\Psi$ is a sublattice in the lattice $Q_{(\pm)}^\Psi$, hence in $\Gamma_{(\pm)}$ as well. Inside the (twisted) Borel QUEA $U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm)$ we consider the elements

$$\begin{aligned} E_i^\Psi &:= K_{-\psi_-(\alpha_i^-)} E_i = K_{-\zeta_{i,-}^-} E_i, & K_{i,+}^\Psi &:= K_{(\text{id} + \psi_+)(\alpha_i^+) - \psi_-(\alpha_i^-)} = K_{\alpha_i^+ + \zeta_{i,+}^+ - \zeta_{i,-}^-} \\ F_i^\Psi &:= K_{+\psi_+(\alpha_i^+)} F_i = K_{+\zeta_{i,+}^+} F_i, & K_{i,-}^\Psi &:= K_{(\text{id} + \psi_-)(\alpha_i^-) - \psi_+(\alpha_i^+)} = K_{\alpha_i^- + \zeta_{i,-}^- - \zeta_{i,+}^+} \end{aligned}$$

(for all $i \in I$) — hereafter called “twisted generators” — and then *define* $\hat{U}_q^\Psi(\mathfrak{b}_+)$, resp. $\hat{U}_q^\Psi(\mathfrak{b}_-)$, as being the unital \mathbb{k}_q -subalgebra of $U_{q,\Gamma(\pm)}^\Psi(\mathfrak{b}_\pm)$ generated by the E_i^Ψ ’s and the $(K_{i,+}^\Psi)^{\pm 1}$ ’s, resp. by the F_i^Ψ ’s and the $(K_{i,-}^\Psi)^{\pm 1}$ ’s.

More in general, for any other sublattice M_\pm containing $Q_{[\pm]}^\Psi$ we *define* $\hat{U}_{q,M_\pm}^\Psi(\mathfrak{b}_+)$, resp. $\hat{U}_{q,M_\pm}^\Psi(\mathfrak{b}_-)$, as the unital \mathbb{k}_q -subalgebra of $U_{q,M_\pm + Q_{(\pm)}^\Psi}(\mathfrak{b}_\pm)$ generated by the E_i^Ψ ’s and the K_{y_+} ’s, resp. by the F_i^Ψ ’s and the K_{y_-} ’s, with $i \in I$ and $y_\pm \in M_\pm$.

The key fact is that the (twisted) Hopf structure of $U_{q,M_\pm + Q_{(\pm)}^\Psi}(\mathfrak{b}_\pm)$ yields

$$\begin{aligned} \Delta^{(\Psi)}(E_\ell^\Psi) &= E_\ell^\Psi \otimes 1 + K_{+\tau_\ell^+} \otimes E_\ell^\Psi, & \mathcal{S}^{(\Psi)}(E_\ell^\Psi) &= -K_{-\tau_\ell^+} E_\ell^\Psi, & \epsilon(E_\ell^\Psi) &= 0 \\ \Delta^{(\Psi)}(K_{y_\pm}) &= K_{y_\pm} \otimes K_{y_\pm}, & \mathcal{S}^{(\Psi)}(K_{y_\pm}) &= K_{y_\pm}^{-1} = K_{-y_\pm}, & \epsilon(K_{y_\pm}) &= 1 \\ \Delta^{(\Psi)}(F_\ell^\Psi) &= F_\ell^\Psi \otimes K_{-\tau_\ell^-} + 1 \otimes F_\ell^\Psi, & \mathcal{S}^{(\Psi)}(F_\ell^\Psi) &= -F_\ell^\Psi K_{+\tau_\ell^-}, & \epsilon(F_\ell^\Psi) &= 0 \end{aligned}$$

for all $\ell \in I$. Altogether, these formulas show that

$$\hat{U}_{q,M_\pm}^\Psi(\mathfrak{b}_\pm) \text{ is a Hopf subalgebra of } U_{q,M_\pm + Q_{(\pm)}^\Psi}(\mathfrak{b}_\pm).$$

It is clear from definitions that if $M_\pm \supseteq Q_{(\pm)}^\Psi$ then we have $\hat{U}_{q,M_\pm}^\Psi(\mathfrak{b}_\pm) = U_{q,M_\pm}^\Psi(\mathfrak{b}_\pm)$ — cf. §3.4.7. Thus $\hat{U}_{q,M_+}^\Psi(\mathfrak{b}_+) = U_{q,M_+}^\Psi(\mathfrak{b}_+)$ can be generated by either set of generators $\{E_i, K_{y_+}\}_{i \in I, y_+ \in M_+}$ or $\{E_i^\Psi, K_{y_+}\}_{i \in I, y_+ \in M_+}$, while $\hat{U}_{q,M_-}^\Psi(\mathfrak{b}_-) = U_{q,M_-}^\Psi(\mathfrak{b}_-)$ can be generated by either $\{F_i, K_{y_-}\}_{i \in I, y_- \in M_-}$ or $\{F_i^\Psi, K_{y_-}\}_{i \in I, y_- \in M_-}$.

An entirely similar remark applies if we use the E_i^Ψ ’s, K_{y_\pm} ’s and F_i^Ψ ’s altogether as algebra generators of $U_{q,\Gamma(*)}^\Psi(\mathfrak{g}_D)$ for $\Gamma(*) := \Gamma_{(+)} + \Gamma_{(-)}$ with $\Gamma_{(\pm)} \supseteq Q_{(\pm)}^\Psi$.

Of course we can repeat this analysis for quantum Borel subalgebras inside $U_{q,\Gamma}^\Psi(\mathfrak{g})$ — for suitable lattices Γ in $\mathbb{Q}Q$ — and introduce “twisted” algebra generators for them, that will also be, altogether, (twisted) algebra generators of $U_{q,\Gamma}^\Psi(\mathfrak{g})$ itself.

Now, when we use the “twisted” generators, the above shows that the formulas for the (twisted) coproduct, antipode and counit look exactly like those for the generators of the *untwisted* Borel subalgebras. In other words, for these quantum Borel subalgebras *twisting the algebra generators (as above) we end up with untwisted formulas for the coalgebra and antipodal structure*. This remark applies to *both* cases: twisted Borel subalgebras (and their generators) inside a double quantum group $U_{q, \Gamma_0}^\Psi(\mathfrak{g}_D)$ and embedded inside a (“single”) quantum group $U_{q, \Gamma}^\Psi(\mathfrak{g})$.

In all these new presentations, the new (twisted) generators of our Borel algebras (inside $U_{q, \Gamma_0}^\Psi(\mathfrak{g}_D)$ and inside $U_{q, \Gamma}^\Psi(\mathfrak{g})$) enjoy new “twisted” relations. The striking fact is that in these presentations the “twisted” relations happen to present a well precise form, commonly formalized via the notion of “multiparameter quantum group”: we shall investigate this in detail in the forthcoming sections.

3.4.9. Comultiplication twistings for polynomial TwQUEA’s. With notation as above, let $\Psi \in M_n(\mathbb{Q})$ be as in §3.4.1, and let $U_{q, \Gamma}^\Psi(\mathfrak{g})$ be the corresponding polynomial TwQUEA (for some lattice Γ) over $\mathfrak{g}'(A)$. From its very construction, it follows that the comultiplication twisting procedure can be iterated: namely, for $\Psi' \in M_n(\mathbb{Q})$ we can define the Hopf algebra $(U_{q, \Gamma}^\Psi(\mathfrak{g}))^{\Psi'}$ just reproducing the construction of $U_{q, \Gamma}^\Psi(\mathfrak{g})$ but with Ψ' replacing Ψ and $U_{q, \Gamma}^\Psi(\mathfrak{g})$ replacing $U_{q, \Gamma}(\mathfrak{g})$.

It is then clear that $(U_{q, \Gamma}^\Psi(\mathfrak{g}))^{\Psi'} = U_{q, \Gamma}^{\Psi+\Psi'}(\mathfrak{g})$, as this equality holds for the TwQUEA $U_h^\Psi(\mathfrak{g})$. Therefore, all these new Hopf algebras will still be called TwQUEA’s.

The same construction is possible, and similar remarks apply, for TwQUEA’s associated with $\mathfrak{g}'_D(A)$, $\mathfrak{b}'_+(A)$ and $\mathfrak{b}'_-(A)$.

4. MULTIPARAMETER QUANTUM GROUPS

We introduce now the multiparameter QUEA $U_{\mathbf{q}}(\mathfrak{g}_D)$, or *MpQUEA* for short, associated with a suitable matrix of parameters. Hereafter, \mathbb{F} will be a field of characteristic zero, and $\mathbb{F}^\times := \mathbb{F} \setminus \{0\}$.

4.1. Defining multiparameter QUEA’s (=MpQUEA’s).

We fix a multiparameter matrix $\mathbf{q} := (q_{ij})_{i, j \in I} \in M_n(\mathbb{F})$, with $i, j \in I$, $i \neq j$, and a generalized, symmetrisable Cartan matrix $A := (a_{ij})_{i, j \in I}$, with the additional compatibility assumption that $q_{ii}^k \neq 1$ for $k = 1, \dots, \max(1 - a_{ij}, 2)$. Moreover, when \mathbf{q} is of Cartan type we assume that its associated (generalized, symmetrisable) Cartan matrix — cf. §2.2.1 — coincides with the matrix A mentioned above, and we fix scalars $q_i \in \mathbb{F}$ ($i \in I$) as in §2.2.1. Finally, like in §2.1, we denote by $\mathfrak{g} = \mathfrak{g}(A)$ the unique Kac-Moody algebra associated with A , by \mathfrak{g}' its derived subalgebra, etc.

Definition 4.1.1. (cf. [HPR]) We denote by $U_{\mathbf{q}}(\mathfrak{g}_D)$ the unital associative algebra over \mathbb{F} with generators $E_i, F_i, K_i^{\pm 1}, L_i^{\pm 1}$ (for all $i \in I$) and relations

$$\begin{aligned}
 (a) \quad & K_i^{\pm 1} L_j^{\pm 1} = L_j^{\pm 1} K_i^{\pm 1} \quad , \quad K_i^{\pm 1} K_i^{\mp 1} = 1 = L_i^{\pm 1} L_i^{\mp 1} \\
 (b) \quad & K_i^{\pm 1} K_j^{\pm 1} = K_j^{\pm 1} K_i^{\pm 1} \quad , \quad L_i^{\pm 1} L_j^{\pm 1} = L_j^{\pm 1} L_i^{\pm 1} \\
 (c) \quad & K_i E_j K_i^{-1} = q_{ij} E_j \quad , \quad L_i E_j L_i^{-1} = q_{ji}^{-1} E_j \\
 (d) \quad & K_i F_j K_i^{-1} = q_{ij}^{-1} F_j \quad , \quad L_i F_j L_i^{-1} = q_{ji} F_j
 \end{aligned}$$

$$\begin{aligned}
 (e) \quad & [E_i, F_j] = \delta_{i,j} q_{ii} \frac{K_i - L_i}{q_{ii} - 1} \\
 (f) \quad & \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k}_{q_{ii}} q_{ij}^k E_i^{1-a_{ij}-k} E_j E_i^k = 0 \quad (i \neq j) \\
 (g) \quad & \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k}_{q_{ii}} q_{ij}^k F_i^{1-a_{ij}-k} F_j F_i^k = 0 \quad (i \neq j)
 \end{aligned}$$

Moreover, $U_{\mathbf{q}}(\mathfrak{g}_D)$ is a Hopf algebra with coproduct, counit and antipode determined for all $i, j \in I$ by

$$\begin{aligned}
 \Delta(E_i) &= E_i \otimes 1 + K_i \otimes E_i, & \epsilon(E_i) &= 0, & \mathcal{S}(E_i) &= -K_i^{-1} E_i \\
 \Delta(F_i) &= F_i \otimes L_i + 1 \otimes F_i, & \epsilon(F_i) &= 0, & \mathcal{S}(F_i) &= -F_i L_i^{-1} \\
 \Delta(K_i^{\pm 1}) &= K_i^{\pm 1} \otimes K_i^{\pm 1}, & \epsilon(K_i^{\pm 1}) &= 1, & \mathcal{S}(K_i^{\pm 1}) &= K_i^{\mp 1} \\
 \Delta(L_i^{\pm 1}) &= L_i^{\pm 1} \otimes L_i^{\pm 1}, & \epsilon(L_i^{\pm 1}) &= 1, & \mathcal{S}(L_i^{\pm 1}) &= L_i^{\mp 1}
 \end{aligned}$$

Finally, for later use we introduce also, for every $\lambda = \sum_{i \in I} \lambda_i \alpha_i \in Q$, the notation $K_\lambda := \prod_{i \in I} K_i^{\lambda_i}$ and $L_\lambda := \prod_{i \in I} L_i^{\lambda_i}$.

Remark 4.1.2. Assume that $q \in \mathbb{F}^\times$ is not a root of unity and fix the q -canonical multiparameter $\check{\mathbf{q}} := (\check{q}_{ij} = q^{d_i a_{ij}})_{i,j \in I}$ like in (2.5) above; then we can define the corresponding MpQUEA as above, now denoted $U_{\check{\mathbf{q}}}(\mathfrak{g}_D)$. The celebrated one-parameter quantum group $U_q(\mathfrak{g})$ defined by Jimbo and Lusztig is (up to a minimal change of generators in its presentation, irrelevant for what follows) nothing but the quotient of our $U_{\check{\mathbf{q}}}(\mathfrak{g}_D)$ by the (Hopf) ideal generated by $\{L_i - K_i^{-1} \mid i = 1, \dots, n\}$.

As a matter of fact, *most constructions usually carried on for $U_q(\mathfrak{g})$ actually makes sense and apply the same to $U_{\check{\mathbf{q}}}(\mathfrak{g}_D)$ as well.*

We introduce then the so-called “quantum Borel / nilpotent / Cartan subalgebras” of any MpQUEA, say $U_{\mathbf{q}}(\mathfrak{g}_D)$, as follows:

Definition 4.1.3. Given $\mathbf{q} := (q_{ij})_{i,j \in I}$ and $U_{\mathbf{q}}(\mathfrak{g}_D)$ as in §4.1, we define $U_{\mathbf{q}}^0 := U_{\mathbf{q}}(\mathfrak{h}_D)$, $U_{\mathbf{q}}^{+,0}$, $U_{\mathbf{q}}^{-,0}$, $U_{\mathbf{q}}^- := U_{\mathbf{q}}(\mathfrak{n}_-)$, $U_{\mathbf{q}}^+ := U_{\mathbf{q}}(\mathfrak{n}_+)$, $U_{\mathbf{q}}^{\leq} := U_{\mathbf{q}}(\mathfrak{b}_-)$ and $U_{\mathbf{q}}^{\geq} := U_{\mathbf{q}}(\mathfrak{b}_+)$ to be the \mathbb{k} -subalgebras of $U_{\mathbf{q}}(\mathfrak{g}_D)$ respectively generated as

$$\begin{aligned}
 U_{\mathbf{q}}^0 &:= \langle K_i^{\pm 1}, L_i^{\pm 1} \rangle_{i \in I}, & U_{\mathbf{q}}^{+,0} &:= \langle K_i^{\pm 1} \rangle_{i \in I}, & U_{\mathbf{q}}^{-,0} &:= \langle L_i^{\pm 1} \rangle_{i \in I} \\
 U_{\mathbf{q}}^- &:= \langle F_i \rangle_{i \in I}, & U_{\mathbf{q}}^{\leq} &:= \langle F_i, L_i^{\pm 1} \rangle_{i \in I}, & U_{\mathbf{q}}^{\geq} &:= \langle K_i^{\pm 1}, E_i \rangle_{i \in I}, & U_{\mathbf{q}}^+ &:= \langle E_i \rangle_{i \in I}
 \end{aligned}$$

We shall refer to $U_{\mathbf{q}}^{\leq}$ and $U_{\mathbf{q}}^{\geq}$ as to the *positive* and *negative* multiparameter quantum Borel (sub)algebras, and $U_{\mathbf{q}}^0$, $U_{\mathbf{q}}^{+,0}$ and $U_{\mathbf{q}}^{-,0}$ as to the *global*, *positive* and *negative* multiparameter Cartan (sub)algebras.

For later use, we also define \mathcal{G}_n to be the free Abelian group of rank $n := |I|$, that we will write in multiplicative notation, and we denote by $U_{\mathbf{q}}(\mathfrak{h})$ the group algebra of \mathcal{G}_n over \mathbb{F} — which, in spite of notation, is independent of \mathbf{q} . Letting $\{G_i\}_{i \in I}$ be a \mathbb{Z} -basis of \mathcal{G}_n , we have natural Hopf algebra isomorphisms of $U_{\mathbf{q}}(\mathfrak{h})$ with $U_{\mathbf{q}}^{+,0}$ and with $U_{\mathbf{q}}^{-,0}$ given by $G_i^{\pm 1} \mapsto K_i^{\pm 1}$ and $G_i^{\pm 1} \mapsto L_i^{\pm 1}$, respectively.

Remarks 4.1.4.

(a) Let $\mathbf{q} = (q_{ij})_{1 \leq i, j \leq n}$ be a multiparameter matrix of Cartan type. Following [He, Section 4], we define the following braided vector spaces:

(i) V_E with \mathbb{F} -basis $\{E_1, \dots, E_n\}$ and braiding given by

$$c(E_i \otimes E_j) := q_{ij} E_j \otimes E_i \quad \text{for all } 1 \leq i, j \leq n ,$$

(ii) V_F with \mathbb{F} -basis $\{F'_1 = L_1^{-1} F_1, \dots, F'_n = L_n^{-1} F_n\}$ and braiding given by

$$c(F'_i \otimes F'_j) := q_{ji} F'_j \otimes F'_i \quad \text{for all } 1 \leq i, j \leq n .$$

Then we have the corresponding Nichols algebras — of diagonal type — $\mathcal{B}(V_E)$ and $\mathcal{B}(V_F)$, as well as their bosonizations over the group $\Gamma := \mathbb{Z}^n$:

$$\mathcal{H}(V_E) := \mathcal{B}(V_E) \# \mathbb{F}[\Gamma] \quad \text{and} \quad \mathcal{H}(V_F) := \mathcal{B}(V_F) \# \mathbb{F}[\Gamma]$$

Directly from definitions, one has canonical identifications $U_{\mathbf{q}}^- \cong \mathcal{B}(V_F)$ and $U_{\mathbf{q}}^+ \cong \mathcal{B}(V_E)$ (as $U_{\mathbf{q}}^{\pm, 0}$ -comodule algebras), and $U_{\mathbf{q}}^{\leq} \cong \mathcal{H}(V_F)$ and $U_{\mathbf{q}}^{\geq} \cong \mathcal{H}(V_E)$ as Hopf algebras, $U_{\mathbf{q}}^0 \cong \mathbb{F}[\Gamma \times \Gamma]$ as Hopf algebras, etc. For more details on the relation with Nichols algebras and bosonization see [An], [Gar], [He].

(b) It is also known, see for example [AA], that the multiparameter quantum group $U_{\mathbf{q}}(\mathfrak{g}_D)$ can be realized as a Drinfeld double of $U_{\mathbf{q}}^{\leq} \cong \mathcal{H}(V_F)$ and $U_{\mathbf{q}}^{\geq} \cong \mathcal{H}(V_E)$ using the Hopf pairing given in Proposition 4.1.5 below. Thus, in the end, $U_{\mathbf{q}}(\mathfrak{g}_D)$ is a Drinfeld double of bosonizations of Nichols algebras of diagonal type.

From [He, Proposition 4.3] — see also [HPR, Theorem 20] and [AY, Proposition 2.4] — we recall the following:

Proposition 4.1.5.

With the assumptions above, assume in addition that $q_{ii} \neq 1$ for all indices $i \in I$. Then there exists a unique skew-Hopf pairing $\eta : U_{\mathbf{q}}^{\geq} \otimes_{\mathbb{F}} U_{\mathbf{q}}^{\leq} \longrightarrow \mathbb{F}$ such that

$$\eta(K_i, L_j) = q_{ij} , \quad \eta(E_i, F_j) = \delta_{i,j} \frac{-q_{ii}}{q_{ii} - 1} , \quad \eta(E_i, L_j) = 0 = \eta(K_i, F_j)$$

for all $1 \leq i, j \leq n$. It enjoys the following property: for every $E \in U_{\mathbf{q}}^+$, $F \in U_{\mathbf{q}}^-$, and every Laurent monomials K in the K_i 's and L in the L_i 's, we have

$$\eta(EK, FL) = \eta(E, F) \eta(K, L) \quad \square$$

For later use, we also single out the following result, whose proof is trivial:

Lemma 4.1.6. *Keep notation as above. There exists a unique Hopf \mathbb{F} -algebra epimorphism $p : U_{\mathbf{q}}(\mathfrak{g}_D) \longrightarrow U_{\mathbf{q}}(\mathfrak{h})$ given by*

$$p(E_i) = 0 = p(F_i) , \quad p(K_i^{\pm 1}) = G_i^{\pm 1} = p(L_i^{\pm 1}) \quad \forall i \in I . \quad \square$$

4.2. Multiplication twistings of MpQUEA's.

We want to perform on the Hopf algebras $U_{\mathbf{q}}(\mathfrak{g}_D)$ a multiplication twisting process, via special types of 2-cocycles, like in §2.4.1, following [AST], [DT] and [Mo].

4.2.1. Special 2–cocycles of $U_{\mathbf{q}}(\mathfrak{g}_D)$. Let $\mathbf{q} := (q_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ be a multiparameter, and let $\chi : \mathcal{G}_n \times \mathcal{G}_n \longrightarrow \mathbb{F}^\times$ be any bicharacter of the Abelian group \mathcal{G}_n . Then χ is automatically a 2–cocycle of \mathcal{G}_n ; as such, it induces (by \mathbb{F} –linear extension) a normalized 2–cocycle of the group algebra of \mathcal{G}_n over \mathbb{F} , that is $U_{\mathbf{q}}(\mathfrak{h})$, cf. Definition 4.1.3: we denote this 2–cocycle of $U_{\mathbf{q}}(\mathfrak{h})$ by χ again. Now, composing $\chi : U_{\mathbf{q}}(\mathfrak{h}) \otimes U_{\mathbf{q}}(\mathfrak{h}) \longrightarrow \mathbb{F}$ with $p^{\otimes 2}$ — where $p : U_{\mathbf{q}}(\mathfrak{g}_D) \twoheadrightarrow U_{\mathbf{q}}(\mathfrak{h})$ is the Hopf \mathbb{F} –algebra epimorphism in Lemma 4.1.6 — we get a (normalized) bimultiplicative Hopf 2–cocycle $\sigma_\chi := \chi \circ p^{\otimes 2}$ of $U_{\mathbf{q}}(\mathfrak{g}_D)$. This leads us to the following definition:

Definition 4.2.2. We denote by $\tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ the set of all (normalized, bimultiplicative) Hopf 2–cocycles of $U_{\mathbf{q}}(\mathfrak{g}_D)$ of the form σ_χ (as defined above) for some bicharacter χ of the Abelian group \mathcal{G}_n . We call these “toral” 2–cocycles of $U_{\mathbf{q}}(\mathfrak{g}_D)$.

Note, in particular, that $\tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ is independent of the multiparameter \mathbf{q} .

By construction, for every $\sigma_\chi \in \tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ one has

$$\begin{aligned} \sigma_\chi(E_i, Y) &= 0 \quad , \quad \sigma_\chi(X, F_i) = 0 \quad \forall Y, X \in U_{\mathbf{q}}(\mathfrak{g}_D), \quad i, j \in I \\ \sigma_\chi(Z'_i, Z''_j) &= \chi(G_i, G_j) \quad \forall Z'_i, Z''_j \in \{K_i, L_i\}_{i \in I}, \quad i, j \in I \end{aligned}$$

in particular, the 2–cocycle σ_χ is determined uniquely by the values of χ on the Abelian group \mathcal{G}_n . Conversely, each (normalized) bimultiplicative Hopf 2–cocycle σ of $U_{\mathbf{q}}(\mathfrak{g}_D)$ such that

$$\begin{aligned} \sigma(E_i, Y) &= 0 \quad , \quad \sigma(X, F_i) = 0 \quad \forall Y, X \in U_{\mathbf{q}}(\mathfrak{g}_D), \quad i, j \in I \\ \sigma(K_i, K_j) &= \sigma(K_i, L_j) = \sigma(L_i, K_j) = \sigma(L_i, L_j) \quad \forall i, j \in I \end{aligned} \quad (4.1)$$

is necessarily of the form $\sigma = \sigma_\chi$, as it is uniquely determined by its values on the torus, hence σ belongs to $\tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$. The corresponding bicharacter χ is then defined by the conditions (for all $i, j \in I$)

$$\chi(G_i, G_j) = \sigma(K_i, K_j) \quad (= \sigma(K_i, L_j) = \sigma(L_i, K_j) = \sigma(L_i, L_j))$$

Therefore, (4.1) are the conditions that characterize intrinsically those (normalized, bimultiplicative) Hopf 2–cocycles of $U_{\mathbf{q}}(\mathfrak{g}_D)$ that belong to $\tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$.

4.2.3. Deforming MpQUEAs via 2–cocycles. For any multiparameter $\mathbf{q} := (q_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ we fix the notation $K_\lambda := \prod_{i \in I} K_i^{\lambda_i}$ and $L_\lambda := \prod_{i \in I} L_i^{\lambda_i}$ for every $\lambda = \sum_{i \in I} \lambda_i \alpha_i \in Q$; similarly, we shall also write $q_{\mu\nu} := \prod_{i,j \in I} q_{ij}^{\mu_i \nu_j}$ for each

$\mu = \sum_{i \in I} \mu_i \alpha_i$ and $\nu = \sum_{i \in I} \nu_i \alpha_i$ in Q .

If \mathbf{q} is of Cartan type, we fix a special element $q := q_{j_0} \in \mathbb{F}^\times$ as explained in §2.2.1 (so that $q_{ii} = q^{2d_i}$ for all $i \in I$); accordingly, we have also a well-defined multiparameter of canonical type, namely $\tilde{\mathbf{q}} := (q^{d_i a_{ij}})_{i,j \in I}$.

Finally, $U_{\mathbf{q}}(\mathfrak{g}_D)$ will be the MpQUEA associated with \mathbf{q} as in Definition 4.1.1; similarly, we have also $U_{\tilde{\mathbf{q}}}(\mathfrak{g}_D)$ when \mathbf{q} is of Cartan type.

The key result that we shall rely upon in the sequel is the following:

Theorem 4.2.4. (cf. [LHR, Theorem 4.5] and [HPR, Theorem 28])

(a) If any two multiparameters are \sim –equivalent (notation as in §2.2.2), then their associated MpQUEA’s are multiplicative twistings of each other. In detail, if

$\mathbf{q}' := (q'_{ij})_{i,j \in I}, \mathbf{q}'' := (q''_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ are multiparameters such that $\mathbf{q}' \sim \mathbf{q}''$, then there exists $\sigma \in \tilde{\mathcal{Z}}_2(U_{\mathbf{q}'}(\mathfrak{g}_D), \mathbb{F})$ such that $U_{\mathbf{q}''}(\mathfrak{g}_D) \cong (U_{\mathbf{q}'}(\mathfrak{g}_D))_\sigma$.

(b) Any MpQUEA with multiparameter of Cartan type is a multiplicative twisting of a MpQUEA with canonical multiparameter. In detail, if $\mathbf{q} := (q_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ is a multiparameter of Cartan type, with associated $q \in \mathbb{F}^\times$ and Cartan matrix $A = (a_{ij})_{i,j \in I}$, and corresponding multiparameter of canonical type $\check{\mathbf{q}} := (q^{d_i a_{ij}})_{i,j \in I}$, then there exists $\sigma \in \tilde{\mathcal{Z}}_2(U_{\check{\mathbf{q}}}(\mathfrak{g}_D), \mathbb{F})$ such that $U_{\mathbf{q}}(\mathfrak{g}_D) \cong (U_{\check{\mathbf{q}}}(\mathfrak{g}_D))_\sigma$.

(c) Similar results hold true for quantum Borel (sub)MpQUEA's as well.

Proof. For claim (a), by Lemma 2.2.3(a) there exists $\nu = (\nu_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ such that $\mathbf{q}'' = \nu \cdot \mathbf{q}'$, that is, $q''_{ij} = \nu_{ij} q'_{ij} \nu_{ji}^{-1}$ for all $i, j \in I$. Clearly there exists a unique bicharacter χ_ν of the Abelian group \mathcal{G}_n defined by $\chi_\nu(G_i, G_j) := \nu_{ij}$ (for $i, j \in I$): then, like in §4.2.1, this χ_ν defines a unique $\sigma = \sigma_{\chi_\nu} \in \tilde{\mathcal{Z}}_2(U_{\mathbf{q}'}(\mathfrak{g}_D), \mathbb{F})$. Through a rather straightforward calculation — much like in [LHR] and [HPR] — one may see that $U_{\mathbf{q}''}(\mathfrak{g}_D) \cong (U_{\mathbf{q}'}(\mathfrak{g}_D))_\sigma$, as claimed; for the sake of completeness, we include below most of the computations. As in §2.4.1, we write $a \cdot_\sigma b := m_\sigma(a, b)$.

First, we show that the generators $E_i, F_i, K_i^{\pm 1}, L_i^{\pm 1}$ (for all $i \in I$) of $U_{\mathbf{q}'}(\mathfrak{g}_D)$ do satisfy the relations that define $U_{\mathbf{q}''}(\mathfrak{g}_D)$ when computed with this new product.

In general, the product of any two group-like elements remains unchanged under a multiplicative twist by a 2-cocycle, for

$$g \cdot_\sigma h = \sigma(g, h) g h \sigma^{-1}(g, h) = \epsilon(g) \epsilon(h) g h = g h \quad \text{for all } g, h \in G(U_{\mathbf{q}'}(\mathfrak{g}_D)) .$$

Since the elements $K_i^{\pm 1}, L_i^{\pm 1}$ are group-like, for all $i \in I$, we have

$$\begin{aligned} (a) \quad & K_i^{\pm 1} \cdot_\sigma L_j^{\pm 1} = L_j^{\pm 1} \cdot_\sigma K_i^{\pm 1} \quad , \quad K_i^{\pm 1} \cdot_\sigma K_i^{\mp 1} = 1 = L_i^{\pm 1} \cdot_\sigma L_i^{\mp 1} \\ (b) \quad & K_i^{\pm 1} \cdot_\sigma K_j^{\pm 1} = K_j^{\pm 1} \cdot_\sigma K_i^{\pm 1} \quad , \quad L_i^{\pm 1} \cdot_\sigma L_j^{\pm 1} = L_j^{\pm 1} \cdot_\sigma L_i^{\pm 1} \end{aligned}$$

In order to compute the remaining relations, one has to keep in mind formulas involving the coproduct and powers of the generators, for instance

$$\begin{aligned} \Delta(E_i^k) &= \sum_{\ell=0}^k \binom{k}{\ell}_{q'_{ii}} E_i^{k-\ell} K_i^\ell \otimes E_i^\ell \quad , \quad \Delta(F_i^k) = \sum_{\ell=0}^k \binom{k}{\ell}_{q'_{ii}} F_i^{k-\ell} \otimes F_i^\ell L_i^{k-\ell} \\ \Delta^{(2)}(E_i^k) &= \sum_{\ell=0}^k \sum_{j=0}^{\ell} \binom{k}{\ell}_{q'_{ii}} \binom{\ell}{j}_{q'_{ii}} E_i^{k-\ell} K_i^\ell \otimes E_i^{\ell-j} K_i^j \otimes E_i^j \\ \Delta^{(2)}(F_i^k) &= \sum_{\ell=0}^k \sum_{j=0}^{\ell} \binom{k}{\ell}_{q'_{ii}} \binom{\ell}{j}_{q'_{ii}} F_i^{k-\ell-j} \otimes F_i^j L_i^{k-\ell-j} \otimes F_i^\ell L_i^{k-\ell} \\ \Delta^{(2)}(E_i^k E_s) &= \sum_{\ell=0}^k \sum_{j=0}^{\ell} \binom{k}{\ell}_{q'_{ii}} \binom{\ell}{j}_{q'_{ii}} \left(E_i^{k-\ell} K_i^\ell E_s \otimes E_i^{\ell-j} K_i^j \otimes E_i^j + \right. \\ &\quad \left. + E_i^{k-\ell} K_i^\ell K_s \otimes E_i^{\ell-j} K_i^j E_s \otimes E_i^j + E_i^{k-\ell} K_i^\ell K_s \otimes E_i^{\ell-j} K_i^j K_s \otimes E_i^j E_s \right) \end{aligned}$$

for all $i, s \in I$ and all $k \geq 0$. Here $\Delta^{(2)} := (\Delta \otimes \text{id}) \circ \Delta = (\text{id} \otimes \Delta) \circ \Delta$. So, for the relations in (c) we have

$$\begin{aligned}
 K_i \cdot_{\sigma} E_j &= \sigma(K_i, E_j) K_i \sigma^{-1}(K_i, 1) + \sigma(K_i, K_j) K_i E_j \sigma^{-1}(K_i, 1) + \\
 &\quad + \sigma(K_i, K_j) K_i K_j \sigma^{-1}(K_i, E_j) = \sigma(K_i, K_j) K_i E_j = \nu_{ij} K_i E_j \\
 (K_i E_j) \cdot_{\sigma} K_i^{-1} &= \sigma(K_i E_j, K_i^{-1}) K_i K_i^{-1} \sigma^{-1}(K_i, K_i^{-1}) + \\
 &\quad + \sigma(K_i K_j, K_i^{-1}) K_i E_j K_i^{-1} \sigma^{-1}(K_i, K_i^{-1}) + \\
 &\quad + \sigma(K_i K_j, K_i^{-1}) K_i K_j K_i^{-1} \sigma^{-1}(K_i E_j, K_i^{-1}) = \\
 &= \sigma(K_i K_j, K_i^{-1}) K_i E_j K_i^{-1} \sigma^{-1}(K_i, K_i^{-1}) = \\
 &= \sigma(K_j, K_i^{-1}) K_i E_j K_i^{-1} = \nu_{ji}^{-1} K_i E_j K_i^{-1} = \nu_{ji}^{-1} q'_{ij} E_j
 \end{aligned}$$

for all $i, j \in I$. Therefore,

$$K_i \cdot_{\sigma} E_j \cdot_{\sigma} K_i^{-1} = \nu_{ij} (K_i E_j) \cdot_{\sigma} K_i^{-1} = \nu_{ij} q'_{ij} \nu_{ji}^{-1} E_j = q''_{ij} E_j$$

Similarly, we have

$$\begin{aligned}
 L_i \cdot_{\sigma} E_j \cdot_{\sigma} L_i^{-1} &= (L_i \cdot_{\sigma} E_j) \cdot_{\sigma} L_i^{-1} = \nu_{ij} (L_i E_j) \cdot L_i^{-1} = \nu_{ij} \nu_{ji}^{-1} L_i E_j L_i^{-1} = \\
 &= \nu_{ij} (q'_{ji})^{-1} \nu_{ji}^{-1} E_j = (\nu_{ji} q'_{ji} \nu_{ij}^{-1})^{-1} E_j = (q''_{ji})^{-1} E_j
 \end{aligned}$$

The computation for relation (d) is entirely similar: for all $i, j \in I$ we have

$$\begin{aligned}
 K_i \cdot_{\sigma} F_j &= \sigma(K_i, F_j) K_i L_j \sigma^{-1}(K_i, L_j) + \sigma(K_i, 1) K_i F_j \sigma^{-1}(K_i, L_j) + \\
 &\quad + \sigma(K_i, 1) K_i \sigma^{-1}(K_i, F_j) = \nu_{ij}^{-1} K_i F_j \\
 (K_i F_j) \cdot_{\sigma} K_i^{-1} &= \sigma(K_i F_j, K_i^{-1}) K_i L_j K_i^{-1} \sigma^{-1}(K_i L_j, K_i^{-1}) + \\
 &\quad + \sigma(K_i, K_i^{-1}) K_i F_j K_i^{-1} \sigma^{-1}(K_i L_j, K_i^{-1}) + \\
 &\quad + \sigma(K_i, K_i^{-1}) K_i K_i^{-1} \sigma^{-1}(K_i F_j, K_i^{-1}) \\
 &= \sigma(K_i, K_i^{-1}) K_i F_j K_i^{-1} \sigma^{-1}(K_i L_j, K_i^{-1}) = \\
 &= \sigma(L_j, K_i) K_i F_j K_i^{-1} = \nu_{ji} K_i F_j K_i^{-1} = \nu_{ji} (q'_{ij})^{-1} F_j
 \end{aligned}$$

whence, we obtain

$$K_i \cdot_{\sigma} F_j \cdot_{\sigma} K_i^{-1} = \nu_{ij}^{-1} (K_i F_j) \cdot_{\sigma} K_i^{-1} = \nu_{ij}^{-1} (q'_{ij})^{-1} \nu_{ji} F_j = (\nu_{ij} q'_{ij} \nu_{ji}^{-1})^{-1} F_j = (q''_{ij})^{-1} F_j$$

Also, we have that $L_i \cdot_{\sigma} F_j \cdot_{\sigma} L_i^{-1} = q''_{ji} F_j$.

As to relation (e), it is enough to note that $E_i \cdot_{\sigma} F_j = E_i F_j$, $F_j \cdot_{\sigma} E_i = F_j E_i$ and $q''_{ii} = q''_{ii}$ for all $i, j \in I$. Thus, for instance,

$$\begin{aligned}
 E_i \cdot_{\sigma} F_j &= \sigma(E_i, F_j) L_j \sigma^{-1}(1, L_j) + \sigma(K_i, F_j) E_i L_j \sigma^{-1}(1, L_j) + \\
 &\quad + \sigma(K_i, F_j) K_i L_j \sigma^{-1}(E_i, L_j) + \sigma(E_i, 1) F_j \sigma^{-1}(1, L_j) + \\
 &\quad + \sigma(K_i, 1) E_i F_j \sigma^{-1}(1, L_j) + \sigma(K_i, 1) K_i F_j \sigma^{-1}(E_i, L_j) + \\
 &\quad + \sigma(E_i, 1) \sigma^{-1}(1, F_j) + \sigma(K_i, 1) E_i \sigma^{-1}(1, F_j) + \\
 &\quad + \sigma(K_i, 1) K_i \sigma^{-1}(E_i, F_j) = E_i F_j
 \end{aligned}$$

hence $E_i \cdot_{\sigma} F_j - F_j \cdot_{\sigma} E_i = [E_i, F_j] = \delta_{i,j} q'_{ii} \frac{K_i - L_i}{q'_{ii} - 1} = \delta_{i,j} q''_{ii} \frac{K_i - L_i}{q''_{ii} - 1}$ for all $i, j \in I$.

Finally, we check that the assertion holds for the quantum Serre relations; we prove only (f) , since (g) is completely analogous. We have to verify the equality

$$(f) \quad \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k} (q''_{ii})^{\binom{k}{2}} (q''_{ij})^k E_i^{\cdot\sigma(1-a_{ij}-k)} \cdot_{\sigma} E_j \cdot_{\sigma} E_i^{\cdot\sigma k} = 0 \quad (i \neq j)$$

where $E_i^{\cdot\sigma k}$ denotes the k -th power with respect to the twisted product. A routine check yields that $E_i^{\cdot\sigma k} = \sigma(K_i, K_i) \cdots \sigma(K_i^{k-1}, K_i) E_i^k = \nu_{ii}^{\binom{k-1}{2}} E_i$ for all $k \geq 0$. Thus for all $i \neq j$ the left-hand side of (f) above, that we denote by \clubsuit , now reads

$$\begin{aligned} \clubsuit &= \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k} (q''_{ii})^{\binom{k}{2}} (q''_{ij})^k E_i^{\cdot\sigma(1-a_{ij}-k)} \cdot_{\sigma} E_j \cdot_{\sigma} E_i^{\cdot\sigma k} = \\ &= \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k} (q'_{ii})^{\binom{k}{2}} \nu_{ij}^k (q'_{ij})^k \nu_{ji}^{-k} \nu_{ii}^{\binom{-a_{ij}-k}{2}} \nu_{ii}^{\binom{k-1}{2}} E_i^{(1-a_{ij}-k)} \cdot_{\sigma} E_j \cdot_{\sigma} E_i^k = \\ &= \nu_{ii}^{\frac{a_{ij}(a_{ij}-1)}{2}} \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k} (q'_{ii})^{\binom{k}{2}} (q'_{ij})^k \nu_{ij}^k \nu_{ji}^{-k} \nu_{ii}^{a_{ij}k+k^2-k} E_i^{(1-a_{ij}-k)} \cdot_{\sigma} E_j \cdot_{\sigma} E_i^k \end{aligned}$$

But since $E_i^{(1-a_{ij}-k)} \cdot_{\sigma} E_j = \nu_{ij}^{(1-a_{ij}-k)} E_i^{(1-a_{ij}-k)} E_j$ and

$$\begin{aligned} (E_i^{(1-a_{ij}-k)} E_j) \cdot_{\sigma} E_i^k &= \\ &= \sum_{\ell=0}^{1-a_{ij}-k} \sum_{r=0}^{\ell} \sum_{s=0}^k \sum_{t=0}^s \sigma(E_i^{(1-a_{ij}-k)-\ell} K_i^{\ell} K_j, E_i^{k-s} K_i^s) E_i^{\ell-r} K_i^r E_j E_i^{s-t} K_i^t \sigma^{-1}(E_i^r, E_i^t) = \\ &= \sigma(K_i^{(1-a_{ij}-k)} K_j, K_i^k) E_i^{(1-a_{ij}-k)} E_j E_i^k = \nu_{ii}^{(1-a_{ij}-k)k} \nu_{ji}^k E_i^{(1-a_{ij}-k)} E_j E_i^k \end{aligned}$$

we have $\nu_{ij}^k \nu_{ji}^{-k} \nu_{ii}^{a_{ij}k+k^2-k} E_i^{(1-a_{ij}-k)} \cdot_{\sigma} E_j \cdot_{\sigma} E_i^k = \nu_{ij}^{1-a_{ij}} E_i^{(1-a_{ij}-k)} E_j E_i^k$ and so

$$\clubsuit = \nu_{ii}^{\frac{a_{ij}(a_{ij}-1)}{2}} \nu_{ij}^{1-a_{ij}} \sum_{k=0}^{1-a_{ij}} (-1)^k \binom{1-a_{ij}}{k} (q'_{ii})^{\binom{k}{2}} (q'_{ij})^k E_i^{(1-a_{ij}-k)} E_j E_i^k = 0$$

This proves that there is a Hopf algebra epimorphism $p : U_{\mathbf{q}''}(\mathfrak{g}_D) \longrightarrow (U_{\mathbf{q}'}(\mathfrak{g}_D))_{\sigma}$ which is defined as the identity on the generators. Using the same relation among the multiparameters \mathbf{q}' and \mathbf{q}'' but written as $q'_{ij} = \nu_{ij}^{-1} q''_{ij} \nu_{ji}$, one may define a multiplicative 2-cocycle $\sigma_{\nu^{-1}}$ on $U_{\mathbf{q}''}(\mathfrak{g}_D)$ and, via the same computations, one gets an epimorphism $t : U_{\mathbf{q}'}(\mathfrak{g}_D) \longrightarrow (U_{\mathbf{q}''}(\mathfrak{g}_D))_{\sigma_{\nu^{-1}}}$. Also, one may consider the multiplicative 2-cocycle σ_{ν} on $(U_{\mathbf{q}''}(\mathfrak{g}_D))_{\sigma_{\nu^{-1}}}$ and performing the twisting procedure one again yields $\left((U_{\mathbf{q}''}(\mathfrak{g}_D))_{\sigma_{\nu^{-1}}} \right)_{\sigma_{\nu}} = (U_{\mathbf{q}''}(\mathfrak{g}_D))_{\sigma_{\nu^{-1} * \sigma_{\nu}}} = U_{\mathbf{q}''}(\mathfrak{g}_D)$. On the other hand, since t is an epimorphism one may push forward the cocycle deformation by σ and obtain an epimorphism $t_{\sigma} : (U_{\mathbf{q}'}(\mathfrak{g}_D))_{\sigma} \longrightarrow \left((U_{\mathbf{q}''}(\mathfrak{g}_D))_{\sigma_{\nu^{-1}}} \right)_{\sigma_{\nu}} = U_{\mathbf{q}''}(\mathfrak{g}_D)$. Composing the latter with p , one obtains a sequence of epimorphisms

$$U_{\mathbf{q}''}(\mathfrak{g}_D) \xrightarrow{p} (U_{\mathbf{q}'}(\mathfrak{g}_D))_{\sigma} \xrightarrow{t_{\sigma}} U_{\mathbf{q}''}(\mathfrak{g}_D)$$

whose composition acts as the identity on the generators, so that $t_{\sigma} \circ p = \text{id}$. Similarly, one proves that $p \circ t_{\sigma} = \text{id}$, so that p is actually an isomorphism.

Claim (b) is a special case of (a) , with $\mathbf{q}' := \check{\mathbf{q}}$ and $\mathbf{q}'' := \mathbf{q}$.

Claim (c) is treated like the previous ones. \square

Remark 4.2.5. Our (sketch of) proof of Theorem 4.2.4 above mimics that of Theorem 4.5 in [LHR]: however, in the latter the very existence of a 2–cocycle as required is, to the best of our understanding, not well proved. The proof of Theorem 28 in [HPR] is along the same lines, but the 2–cocycle used there is that of Lemma 2.2.3(c) — hence needs the existence of suitable square roots.

Corollary 4.2.6. *If $\mathbf{q} = (q_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ is a multiparameter, $\boldsymbol{\nu} = (\nu_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ and $\sigma_{\chi_{\boldsymbol{\nu}}} \in \tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ is the unique element of $\tilde{\mathcal{Z}}_2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ such that $\sigma_{\chi_{\boldsymbol{\nu}}}(K_i, K_j) = \nu_{ij}$ for $i, j \in I$ (as in the proof of Theorem 4.2.4 above), then*

$$U_{\boldsymbol{\nu}, \mathbf{q}}(\mathfrak{g}_D) \cong (U_{\mathbf{q}}(\mathfrak{g}_D))_{\sigma_{\chi_{\boldsymbol{\nu}}}}$$

An entirely similar claim holds true for Borel (sub)MpQUEA’s as well.

Proof. By Lemma 2.2.3(a) we know that $\mathbf{q} \sim \boldsymbol{\nu}, \mathbf{q}$; then Theorem 4.2.4(a) above applies, and the claim follows at once. \square

We still need some extra notation:

Definition 4.2.7. For any $\mathbf{q} = (q_{ij})_{i,j \in I}$, $\mathbf{q}' = (q'_{ij})_{i,j \in I} \in M_n(\mathbb{F})$, we say that $\mathbf{q} \approx \mathbf{q}'$ if and only if there exists a permutation γ of I such that

$$q'_{ij} = q_{\gamma(i)\gamma(j)} \quad \forall i, j \in I \quad \text{or} \quad q'_{ij} = q_{\gamma(j)\gamma(i)}^{-1} \quad \forall i, j \in I.$$

Proposition 4.2.8. *Let $\mathbf{q} = (q_{ij})_{i,j \in I}$, $\mathbf{q}' = (q'_{ij})_{i,j \in I} \in M_n(\mathbb{F})$ be multiparameters. Then $U_{\mathbf{q}}(\mathfrak{g}_D) \cong U_{\mathbf{q}'}(\mathfrak{g}_D) \iff \mathbf{q} \approx \mathbf{q}'$.*

Proof. We denote by $E_i, F_i, K_i^{\pm 1}$ and $L_i^{\pm 1}$ (with $i \in I$) the generators of $U_{\mathbf{q}}(\mathfrak{g}_D)$ and by $E'_j, F'_j, K'_j{}^{\pm 1}$ and $L'_j{}^{\pm 1}$ (with $j \in I$) the generators of $U_{\mathbf{q}'}(\mathfrak{g}_D)$.

Assume first that there is a Hopf algebra isomorphism $\varphi : U_{\mathbf{q}}(\mathfrak{g}_D) \longrightarrow U_{\mathbf{q}'}(\mathfrak{g}_D)$. Note that both Hopf algebras $U_{\mathbf{q}}(\mathfrak{g}_D)$ and $U_{\mathbf{q}'}(\mathfrak{g}_D)$ are pointed, as they are generated by group-likes and skew-primitives: then, in particular, their coradicals coincide with the subalgebras $U_{\mathbf{q}}(\mathfrak{h}_D)$ and $U_{\mathbf{q}'}(\mathfrak{h}_D)$, respectively, which are actually isomorphic. Moreover, by [Gar, Theorem 3.8], they belong to a special family of pointed Hopf algebras introduced in [ARS]; then, using the coradical filtration one can write $\text{gr}(U_{\mathbf{q}}(\mathfrak{g}_D)) = R \# U_{\mathbf{q}}(\mathfrak{h}_D)$ and $\text{gr}(U_{\mathbf{q}'}(\mathfrak{g}_D)) = R' \# U_{\mathbf{q}'}(\mathfrak{h}_D)$, where R and R' are braided Hopf algebras in the category of Yetter-Drinfeld modules over $U_{\mathbf{q}}(\mathfrak{h}_D)$ and $U_{\mathbf{q}'}(\mathfrak{h}_D)$, respectively. These are given by the subalgebra of coinvariants, i.e., $R = \{x \in \text{gr}(U_{\mathbf{q}}(\mathfrak{g}_D)) \mid (\text{id} \otimes \pi)\Delta(x) = x \otimes 1\}$, where $\pi : \text{gr}(U_{\mathbf{q}}(\mathfrak{g}_D)) \longrightarrow U_{\mathbf{q}}(\mathfrak{h}_D)$ is the canonical projection. As $\varphi(U_{\mathbf{q}}(\mathfrak{h}_D)) = U_{\mathbf{q}'}(\mathfrak{h}_D)$, the map φ induces a braided Hopf algebra isomorphism $\tilde{\varphi} : \text{gr}(U_{\mathbf{q}}(\mathfrak{g}_D)) \longrightarrow \text{gr}(U_{\mathbf{q}'}(\mathfrak{g}_D))$ such that $\tilde{\varphi}(R) = R'$. In particular, $\tilde{\varphi}$ maps primitive element in R to primitive elements in R' .

Now take $i \in I$. As (the image of the element in the graded algebra) E_i is primitive in R , we have that $\tilde{\varphi}(E_i)$ is primitive in R' as well. As the latter is linearly spanned by the elements E'_i and $L'_j{}^{-1}F'_j$ with $i, j \in I$ — these are the linear generators of V_E and V_F in Remarks 4.1.4 above — this implies that $\varphi(K_i) = K'_j$ or $\varphi(K_i) = L'_j{}^{-1}$ for some $j \in I$. In particular, one has that $\varphi(E_i) = c_{i,j}E'_j$ or $\varphi(E_i) = d_{i,j}L'_j{}^{-1}F'_j$, respectively, for some $c_{i,j}, d_{i,j} \in \mathbb{F}$. Thus, in the end, φ induces

a bijective map $\gamma : I \longrightarrow I$ such that $\varphi(K_i) = K'_{\gamma(i)}$ and $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$ or $\varphi(K_i) = L'^{-1}_{\gamma(i)}$ and $\varphi(E_i) = d_{i,\gamma(i)}L'^{-1}_{\gamma(i)}F'_{\gamma(i)}$.

Claim 1: $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$ for all $i \in I$ or $\varphi(E_i) = d_{i,\gamma(i)}L'^{-1}_{\gamma(i)}F'_{\gamma(i)}$ for all $i \in I$.

Indeed, suppose that $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$ and $\varphi(E_j) = d_{j,\gamma(j)}L'^{-1}_{\gamma(j)}F'_{\gamma(j)}$ for some $i \neq j \in I$; as the Cartan matrix A is indecomposable — hence the Dynkin diagram is connected — we can assume in addition that i and j are adjacent in the Dynkin diagram, i.e., $a_{ij} \neq 0$. Now, $\varphi(K_i) = K'_{\gamma(i)}$ and $\varphi(K_j) = L'^{-1}_{\gamma(j)}$; since

$$q_{ij} \varphi(E_j) = \varphi(K_i E_j K_i^{-1}) = d_{j,\gamma(j)} K'_{\gamma(i)} L'^{-1}_{\gamma(j)} F'_{\gamma(j)} K'^{-1}_{\gamma(i)} = q'^{-1}_{\gamma(i)\gamma(j)} \varphi(E_j)$$

it follows that $q_{ij} = q'^{-1}_{\gamma(i)\gamma(j)}$. On the other hand,

$$q_{ji} \varphi(E_i) = \varphi(K_j E_i K_j^{-1}) = c_{i,\gamma(i)} L'^{-1}_{\gamma(j)} E'_{\gamma(i)} L'_{\gamma(j)} = q'_{\gamma(i)\gamma(j)} \varphi(E_i)$$

which implies that $q_{ji} = q'_{\gamma(i)\gamma(j)}$. Therefore $q_{ii}^{a_{ij}} = q_{ij} q_{ji} = 1$ and $a_{ij} = 0$, as the q_{ii} 's are not a_{ij} -th roots of unity: but $a_{ij} \neq 0$ by assumption, a contradiction.

Claim 2: Let $i \in I$. If $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$, respectively $\varphi(E_i) = d_{i,\gamma(i)}L'^{-1}_{\gamma(i)}F'_{\gamma(i)}$, then $\varphi(F_i) = r_{i,\gamma(i)}F'_{\gamma(i)}$, respectively $\varphi(F_i) = s_{i,\gamma(i)}K'^{-1}_{\gamma(i)}E'_{\gamma(i)}$, for some $r_{i,\gamma(i)}, s_{i,\gamma(i)} \in \mathbb{F}^\times$.

Indeed, since $[E_i, F_i] = q_{ii} \frac{K_i - L_i}{q_{ii} - 1} \in U_{\mathbf{q}}(\mathfrak{h}_D)$, and $\varphi(U_{\mathbf{q}}(\mathfrak{h}_D)) = U_{\mathbf{q}'}(\mathfrak{h}_D)$, we must have $\varphi([E_i, F_i]) = [\varphi(E_i), \varphi(F_i)] \in U_{\mathbf{q}'}(\mathfrak{h}_D)$. Hence, if $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$ and $\varphi(F_i) = s_{i,\gamma(i)}K'^{-1}_{\gamma(i)}E'_{\gamma(i)}$, we would have that $[\varphi(E_i), \varphi(F_i)] \in U_{\mathbf{q}'}^{\geq} \cap U_{\mathbf{q}'}(\mathfrak{h}_D)$ is zero or the degree w.r.t. the PBW-element $E'_{\gamma(i)}$ is 2. Thus $[\varphi(E_i), \varphi(F_i)] = 0$, which in turn implies that $\varphi(K_i) = \varphi(L_i)$, against the injectivity of φ . Therefore we must have that $\varphi(F_i) = r_{i,\gamma(i)}F'_{\gamma(i)}$ and $\varphi(L_i) = L'_{\gamma(i)}$.

The remaining case is analogous.

The claims above imply that φ is given by a Dynkin diagram automorphism possibly composed with a ‘‘Chevalley involution’’. Namely, if $\varphi(E_i) = c_{i,\gamma(i)}E'_{\gamma(i)}$ for all $i \in I$, then $q'_{ij} = q_{\gamma(i)\gamma(j)}$ for all $i, j \in I$ and γ is an automorphism of the generalized Dynkin diagram $D(\mathbf{q})$. If instead $\varphi(E_i) = d_{i,\gamma(i)}L'^{-1}_{\gamma(i)}F'_{\gamma(i)}$ for all $i \in I$, then $q'_{ij} = q_{\gamma(j)\gamma(i)}^{-1}$ for all $i, j \in I$, and γ is an automorphism of $D(\mathbf{q})$ composed with a Chevalley involution. Altogether, this implies that $\mathbf{q} \approx \mathbf{q}'$, q.e.d.

Conversely, assume now that $\mathbf{q} \approx \mathbf{q}'$. If $q'_{ij} = q_{\gamma(i)\gamma(j)}$ for all $i, j \in I$, then the algebra map $\varphi : U_{\mathbf{q}}(\mathfrak{g}_D) \longrightarrow U_{\mathbf{q}'}(\mathfrak{g}_D)$ given for all $i \in I$ by

$$\varphi(E_i) := E'_{\gamma(i)}, \quad \varphi(F_i) := F'_{\gamma(i)}, \quad \varphi(K_i^{\pm 1}) := K'_{\gamma(i)}^{\pm 1}, \quad \varphi(L_i^{\pm 1}) := L'_{\gamma(i)}^{\pm 1}$$

is well-defined and yields a Hopf algebra isomorphism. Similarly, if $q'_{ij} = q_{\gamma(j)\gamma(i)}^{-1}$ for all $i, j \in I$, then the algebra map $\varphi : U_{\mathbf{q}}(\mathfrak{g}_D) \longrightarrow U_{\mathbf{q}'}(\mathfrak{g}_D)$ given for $i \in I$ by

$$\varphi(E_i) := L'^{-1}_{\gamma(i)} F'_{\gamma(i)}, \quad \varphi(F_i) := K'^{-1}_{\gamma(i)} E'_{\gamma(i)}, \quad \varphi(K_i^{\pm 1}) := L'^{\mp 1}_{\gamma(i)}, \quad \varphi(L_i^{\pm 1}) := K'^{\mp 1}_{\gamma(i)}$$

is well-defined and yields a Hopf algebra isomorphism. \square

4.3. Multiparameter quantum groups with larger torus.

The MpQUEA's $U_{\mathbf{q}}(\mathfrak{g}_D)$ we considered so far have a toral part (i.e., the subalgebra $U_{\mathbf{q}}^0$ generated by the $K_i^{\pm 1}$'s and the $L_j^{\pm 1}$'s) that is nothing but the group algebra of a double copy of the root lattice Q of \mathfrak{g} , much like in the one-parameter case (but for the duplication of Q , say). Now, in that (uniparameter) case, one also considers MpQUEA's with a larger toral part, namely the group algebra of any intermediate lattice between Q and P ; similarly, we can introduce MpQUEA's whose toral part is the group algebra of any lattice $\Gamma_\ell \times \Gamma_r$ with $Q \subseteq \Gamma_\ell$ and $Q \subseteq \Gamma_r$.

4.3.1. Larger tori for MpQUEA's. The definition of the “toral parts” of our MpQUEA's $U_{\mathbf{q}}(\mathfrak{g}_D)$ — cf. Definition 4.1.3 — is actually independent of the multiparameter \mathbf{q} . We will use this fact to define “larger toral MpQUEA's” as toral parts of some larger MpQUEA's, as we did in §3.3.5. This requires some compatibility constraints on \mathbf{q} ; for later use, we fix now some more preliminary facts.

Let Γ be any sublattice of $\mathbb{Q}Q$ of rank n with $Q \leq \Gamma$. For any basis $\{\gamma_1, \dots, \gamma_n\}$ of Γ , let $C := (c_{ij})_{i,j=1,\dots,n}^{j=1,\dots,n}$ $\in \mathbb{Z}^{n \times n}$ be defined by $\alpha_i = \sum_{j=1}^n c_{ij} \gamma_j$ for every $i \in I = \{1, \dots, n\}$. Write $c := |\det(C)| \in \mathbb{N}_+$ and $C^{-1} = (c'_{ij})_{i,j=1,\dots,n}^{j=1,\dots,n}$ for the inverse matrix to C ; in particular, we have that $c''_{ij} := c \cdot c'_{ij} \in \mathbb{Z}$ for all $i, j \in I$.

For any such Γ , let $U_{\mathbf{q},\Gamma}^{+,0} = \mathbb{F}\Gamma$ be the group algebra with generators $K_{\gamma_i}^{\pm 1}$ (for $i \in I$). If we set $K_{\alpha_i} := \prod_{j \in I} K_{\gamma_j}^{c_{ij}}$ for all $i \in I$, then the \mathbb{F} -subalgebra of $U_{\mathbf{q},\Gamma}^{+,0}$ generated by the $K_{\alpha_i}^{\pm 1}$'s is an isomorphic copy of $U_{\mathbf{q},Q}^{+,0}$. In the obvious symmetric way we define also the “negative counterpart” $U_{\mathbf{q},\Gamma}^{-,0}$ generated by the $L_{\gamma_i}^{\pm 1}$'s.

Finally, given any two sublattices Γ_{\pm} of rank n in $\mathbb{Q}Q$ containing Q , letting $\Gamma_{\bullet} := \Gamma_+ \times \Gamma_-$ we define $U_{\mathbf{q},\Gamma_{\bullet}}^0 := U_{\mathbf{q},\Gamma_+}^{+,0} \otimes_{\mathbb{F}} U_{\mathbf{q},\Gamma_-}^{-,0}$; in this case, the basis elements for Γ_{\pm} will be denoted by γ_i^{\pm} ($i \in I$).

4.3.2. MpQUEA's with larger torus. Let Γ_+ and Γ_- be any two lattices in $\mathbb{Q}Q$ such that $Q \leq \Gamma_{\pm}$; then set $\Gamma_{\bullet} := \Gamma_+ \times \Gamma_-$. For these lattices Γ_{\pm} , we have matrices of integers $C_{\pm} = (c_{ij}^{\pm})_{i,j \in I}$ and $C_{\pm}^{-1} = (c_{ij}^{\pm, \prime})_{i,j \in I}$, and also $c_{\pm} := |\det(C_{\pm})|$.

For the rest of this subsection, we assume that for every $i, j \in I$, the field \mathbb{F} contains a c_{\pm} -th root of q_{ij} , hereafter denoted by $q_{ij}^{1/c_{\pm}}$; moreover, if the multiparameter $\mathbf{q} = (q_{ij})_{i,j \in I}$ is of Cartan type we assume that the multiparameter $\mathbf{q}^{1/c_{\pm}} := (q_{ij}^{1/c_{\pm}})_{i,j \in I}$ is of Cartan type as well. In the sequel, these assumptions will always tacitly be taken whenever we consider MpQUEA's with larger tori.

The natural (adjoint) action of $U_{\mathbf{q}}^0$ onto $U_{\mathbf{q}}$ extends (uniquely) to a $U_{\mathbf{q},\Gamma_{\bullet}}^0$ -action $\cdot : U_{\mathbf{q},\Gamma_{\bullet}}^0 \times U_{\mathbf{q}} \rightarrow U_{\mathbf{q}}$, given by

$$\begin{aligned} K_{\gamma_i^+} \cdot E_j &= q_{ij}^{\Gamma_+} E_j, & L_{\gamma_i^-} \cdot E_j &= (q_{ji}^{\Gamma_-})^{-1} E_j \\ K_{\gamma_i^+} \cdot F_j &= (q_{ij}^{\Gamma_+})^{-1} F_j, & L_{\gamma_i^-} \cdot F_j &= q_{ji}^{\Gamma_-} F_j \\ K_{\gamma_i^+} \cdot K_{\alpha_j} &= K_{\alpha_j}, & L_{\gamma_i^-} \cdot K_{\alpha_j} &= K_{\alpha_j} \end{aligned}$$

— where $q_{rs}^{\Gamma_+} := \prod_{k \in I} (q_{ks})^{c_{rk}^{\Gamma_+}}$ and $q_{ae}^{\Gamma_-} := \prod_{k \in I} (q_{ak})^{c_{ek}^{\Gamma_-}}$ — that makes $U_{\mathbf{q}}$ into a $U_{\mathbf{q},\Gamma_{\bullet}}^0$ -module algebra. This allows us to consider the *smash product* Hopf

algebra $U_{\mathbf{q}, \Gamma_\bullet}^0 \times U_{\mathbf{q}}$; its product is given by the formula (3.16). As $U_{\mathbf{q}, \Gamma_\bullet}^0$ is a right $U_{\mathbf{q}}$ -module Hopf algebra with respect to the adjoint action, one may consider the vector space $U_{\mathbf{q}, \Gamma_\bullet}^0 \otimes_{U_{\mathbf{q}}} U_{\mathbf{q}}$. Moreover, the smash product $U_{\mathbf{q}, \Gamma_\bullet}^0 \times U_{\mathbf{q}}$ maps onto a

Hopf algebra structure on it, which hereafter we denote by $U_{\mathbf{q}, \Gamma_\bullet}^0 \times_{U_{\mathbf{q}}} U_{\mathbf{q}}$ — see, e.g., [Len, Theorem 2.8]. We define then

$$U_{\mathbf{q}, \Gamma_\bullet}(\mathfrak{g}_D) \equiv U_{\mathbf{q}, \Gamma_\bullet} := U_{\mathbf{q}, \Gamma_\bullet}^0 \times_{U_{\mathbf{q}}} U_{\mathbf{q}} = U_{\mathbf{q}, \Gamma_\bullet}^0 \times_{U_{\mathbf{q}}} U_{\mathbf{q}}(\mathfrak{g}_D) \quad (4.2)$$

Since the coalgebra structure is the one given by the tensor product, to give a presentation by generators and relations like that for $U_{\mathbf{q}}(\mathfrak{g}_D)$, one has to describe only the algebra structure. For this, one has to replace the generators $K_i^{\pm 1} = K_{\pm \alpha_i}$ and $L_i^{\pm 1} = L_{\pm \alpha_i}$ with the generators $\mathcal{K}_i^{\pm 1} = K_{\pm \gamma_i^+}$ and $\mathcal{L}_i^{\pm 1} = L_{\pm \gamma_i^-}$; replace relations (c) and (d) of Definition 4.1.1 with the following, generalized relations:

$$\begin{aligned} (c') \quad & K_{\gamma_i^+} E_j K_{\gamma_i^+}^{-1} = q_{ij}^{\Gamma_+} E_j \quad , \quad L_{\gamma_i^-} E_j L_{\gamma_i^-}^{-1} = (q_{ji}^{\Gamma_-})^{-1} E_j \\ (d') \quad & K_{\gamma_i^+} F_j K_{\gamma_i^+}^{-1} = (q_{ij}^{\Gamma_+})^{-1} F_j \quad , \quad L_{\gamma_i^-} F_j L_{\gamma_i^-}^{-1} = q_{ji}^{\Gamma_-} F_j \end{aligned}$$

finally, in relation (e) replace the elements K_i and L_i by K_{α_i} and L_{α_i} , respectively, and leave the quantum Serre relations (f) and (g) unchanged.

With much the same approach, one defines also the “(multiparameter) quantum subgroups” of $U_{\mathbf{q}, \Gamma_\bullet}(\mathfrak{g}_D)$ akin to those of $U_{\mathbf{q}}(\mathfrak{g}_D)$ (cf. Definition 4.1.3), that we denote by adding a subscript Γ_\bullet , namely $U_{\mathbf{q}, \Gamma_\bullet}^+$, $U_{\mathbf{q}, \Gamma_\bullet}^-$, $U_{\mathbf{q}, \Gamma_\bullet}^{\geq}$, $U_{\mathbf{q}, \Gamma_\bullet}^{\leq}$, $U_{\mathbf{q}, \Gamma_\bullet}^{+,0}$ and $U_{\mathbf{q}, \Gamma_\bullet}^{-,0}$.

4.3.3. Duality among MpQUEA’s with larger torus. Let again Γ_\pm be two lattices of rank n in $\mathbb{Q}Q$ containing Q , and set $\Gamma_\bullet := \Gamma_+ \times \Gamma_-$. We repeat our assumptions (cf. §4.3.2). We fix bases $\{\gamma_s^\pm\}_{s \in I}$ of Γ_\pm , the matrices $C_\pm = (c_{ij}^\pm)_{i,j \in I}$, $C_\pm^{-1} = (c_{ij}^{\pm'})_{i,j \in I}$, and write $c_\pm := |\det(C_\pm)|$ and $c_{ij}^{\pm''} := c_\pm \cdot c_{ij}^{\pm'}$ ($i, j \in I$). In addition, we assume that \mathbb{F} contain a $(c_+ c_-)$ -th root of q_{ij} , say $q_{ij}^{1/(c_+ c_-)}$, and that overall the multiparameter $\mathbf{q}^{1/(c_+ c_-)} := \left(q_{ij}^{1/(c_+ c_-)} \right)_{i,j \in I}$ be of Cartan type.

It is straightforward to check that the Hopf skew-pairing $\eta : U_{\mathbf{q}}^{\geq} \otimes_{\mathbb{F}} U_{\mathbf{q}}^{\leq} \longrightarrow \mathbb{F}$ in Proposition 4.1.5 — cf. Remarks 4.1.4(a) too — actually extends to a similar skew-pairing $U_{\mathbf{q}, \Gamma_+}^{\geq} \otimes_{\mathbb{F}} U_{\mathbf{q}, \Gamma_-}^{\leq} \longrightarrow \mathbb{F}$ given for all $i, j \in I$ by

$$\begin{aligned} \eta(E_i \otimes L_{\gamma_j^-}) &= 0 & \eta(K_{\gamma_i^+} \otimes F_j) &= 0 \\ \eta(E_i \otimes F_j) &= -\delta_{ij} \frac{q_{ii}}{q_{ii}-1} & \eta(K_{\gamma_i^+} \otimes L_{\gamma_j^-}) &= \prod_{h,k \in I} \left(q_{hk}^{1/(c_+ c_-)} \right)^{c_{ih}^+ c_{jk}^-} \end{aligned}$$

One easily sees that, using such a Hopf pairing η between (suitably chosen) quantum Borel subgroups $U_{\mathbf{q}, \Gamma_+}^{\geq}$ and $U_{\mathbf{q}, \Gamma_-}^{\leq}$, every MpQUEA with larger torus $U_{\mathbf{q}, \Gamma_\bullet}(\mathfrak{g}_D)$ can be realized as a Drinfeld double (of those quantum Borel subgroups), thus generalizing what happens for MpQUEA’s with “standard” torus — see Remarks 4.1.4(b).

Remark 4.3.4. From the description above, it follows that one clearly can suitably extend both Theorem 4.2.4 and Corollary 4.2.6 to the case of MpQUEA’s with larger torus: Namely, let Γ_\pm be two lattices of rank n in $\mathbb{Q}Q$ containing Q , and

set $\Gamma_\bullet := \Gamma_+ \times \Gamma_-$. In addition, assume that \mathbb{F} contain a $(c_1 c_2)$ -th root of q_{ij} with $c_i \in \{c_+, c_-\}$, say $q_{ij}^{1/(c_1 c_2)}$, and that overall the multiparameter $\mathbf{q}^{1/(c_1 c_2)} := \left(q_{ij}^{1/(c_1 c_2)} \right)_{i,j \in I}$ be of Cartan type. Pick the standard MpQUEA $U_{\check{\mathbf{q}}, \Gamma_\bullet}(\mathfrak{g}_D)$ with larger torus Γ_\bullet and canonical multiparameter $\check{\mathbf{q}}$. Then there exist a normalized 2-cocycle of the Hopf algebra $U_{\check{\mathbf{q}}, \Gamma_\bullet}(\mathfrak{g}_D)$ and a Hopf algebra isomorphism such that

$$U_{\mathbf{q}, \Gamma_\bullet}(\mathfrak{g}_D) \cong \left(U_{\check{\mathbf{q}}, \Gamma_\bullet}(\mathfrak{g}_D) \right)_\sigma$$

5. TWISTED QUEA'S VS. MULTIPARAMETER QUEA'S

In this section we show that, *assuming the Cartan matrix A to be of finite type*, all TwQUEA's — as considered in §3.4 — are actually (isomorphic to) MpQUEA's — for multiparameters in a special, yet quite general, subclass of integral type. Conversely, any MpQUEA for such a multiparameter is (isomorphic to) a TwQUEA for a single, specific twist element. In short — and up to sticking to twists of type (3.5) or (3.9) and to multiparameters of integral type — we have

every TwQUEA is (isomorphic to) a MpQUEA, and viceversa.

In all this section, we fix as ground field $\mathbb{F} := \mathbb{k}_q := \varprojlim_{m \in \mathbb{N}} \mathbb{k}(q^{1/m})$, with \mathbb{k} being a field of characteristic zero and q an indeterminate. As before, by q^r for any $r \in \mathbb{Q}$ we always mean $q^r = q^{a/d} := (q^{1/d})^a$ if $r = a/d$ with $a \in \mathbb{Z}$ and $d \in \mathbb{N}_+$. In addition, we assume that our (indecomposable) Cartan matrix A is of finite type.

5.1. Twists vs. “rational” multiparameters.

We introduce now a special subclass of integral multiparameters (cf. §2.2.1) in $\mathbb{F} := \mathbb{k}_q$; then we link them with twist elements (as in (3.5) or (3.9) alike).

5.1.1. q -rational multiparameters. We fix now some more notation.

First of all, we remark that for any matrix $C \in M_n(\mathbb{Q})$ it is uniquely defined $d_C := \min \{ d \in \mathbb{N} \mid dC \in M_n(\mathbb{Z}) \}$. Second, for any given $R \in M_n(\mathbb{Q})$ we adopt notation $\underline{q}^R := (q^{r_{ij}})_{i,j \in I}$. We say that a multiparameter $\mathbf{q} := (q_{ij})_{i,j \in I}$ in $\mathbb{F} := \mathbb{k}_q$ is of q -rational type $R := (r_{ij})_{i,j \in I}$ if $R \in M_n(\mathbb{Q})$ and \mathbf{q} is of Cartan type with $\mathbf{q} = \underline{q}^R$, i.e., $q_{ij} = q^{r_{ij}}$ for all $i, j \in I$. In other words, writing $r_{ij} = b_{ij}/d_R$ for some $b_{ij} \in \mathbb{Z}$ ($i, j \in I$), we have that $\mathbf{q} := (q_{ij})_{i,j \in I}$ is of q -rational type $R := (r_{ij})_{i,j \in I}$ if and only if it is of integral type $(q^{1/d_R}, B := (b_{ij})_{i,j \in I})$.

In the sequel, we call q -Mp $_{\mathbb{Q}}$ the set of all multiparameters of q -rational type (or simply “ q -rational multiparameters”) — for any possible $R := (r_{ij})_{i,j \in I}$ — in \mathbb{k}_q .

5.1.2. The links {twists} \leftrightarrow {multiparameters} \leftrightarrow {2-cocycles}. We begin defining two maps from $M_n(\mathbb{Q})$ to itself given by

$$\Psi \mapsto \vartheta(\Psi) := d_\Psi^{-1} D A + A^T (\Psi^T - \Psi) A \quad (5.1)$$

and

$$R \mapsto \xi(R) := 2^{-1} A^{-T} (d_R^{-1} D A - R) A^{-1} \quad (5.2)$$

A moment's check shows that the following hold:

- (a) $Im(\vartheta) = \{ R \in M_n(\mathbb{Q}) \mid R + R^T = d_R^{-1} 2DA \} =: \mathfrak{o}_{2DA}(\mathbb{Q})$
- (b) $Im(\xi) = \{ \Psi \in M_n(\mathbb{Q}) \mid \Psi + \Psi^T = 0 \} =: \mathfrak{so}_n(\mathbb{Q})$
- (c) $\xi \circ \vartheta = \text{id}_{\mathfrak{so}_n(\mathbb{Q})}$, $\vartheta \circ \xi = \text{id}_{\mathfrak{o}_{2DA}(\mathbb{Q})}$

As a consequence, $\vartheta' := \vartheta|_{\mathfrak{so}_n(\mathbb{Q})}$ and $\xi' := \xi|_{\mathfrak{o}_{2DA}(\mathbb{Q})}$ yield mutually inverse bijections $\mathfrak{so}_n(\mathbb{Q}) \xrightarrow{\vartheta'} \mathfrak{o}_{2DA}(\mathbb{Q})$ and $\mathfrak{o}_{2DA}(\mathbb{Q}) \xleftarrow{\xi'} \mathfrak{so}_n(\mathbb{Q})$. Note also that, being defined over *antisymmetric* matrices, ϑ' is described by the following modified version of (5.1):

$$\Psi \mapsto \vartheta'(\Psi) := d_\Psi^{-1}DA + A^T(\Psi^T - \Psi)A = (d_\Psi^{-1}D - 2A^T\Psi)A \quad (5.3)$$

In addition, there exists also a natural bijection $\mathfrak{o}_{2DA}(\mathbb{Q}) \xleftarrow{\chi} q\text{-Mp}_{\mathbb{Q}}$ given by $R \mapsto \chi(R) := \underline{q^R}$. Using it, we can define the maps

$$\mathfrak{so}_n(\mathbb{Q}) \xleftarrow{\chi \circ \vartheta'} q\text{-Mp}_{\mathbb{Q}} \quad , \quad \Psi \mapsto \mathbf{q}_\Psi := \underline{q^{\vartheta'(\Psi)}} \quad (5.4)$$

and

$$q\text{-Mp}_{\mathbb{Q}} \xleftarrow{\xi' \circ \chi^{-1}} \mathfrak{so}_n(\mathbb{Q}) \quad , \quad \mathbf{q} = \underline{q^R} \mapsto \Psi_{\mathbf{q}} := \xi'(R) \quad (5.5)$$

that are inverse to each other, hence are bijections.

As a matter of notation, in the following when $\Psi \in \mathfrak{so}_n(\mathbb{Q})$ and $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ are such that $\mathbf{q} = (\chi \circ \vartheta')(\Psi)$ and $\Psi = (\xi' \circ \chi^{-1})(\mathbf{q})$ we shall write in short $\Psi \leftrightarrow \mathbf{q}$.

Finally, by Definition 4.2.2, every multiparameter $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ uniquely defines a corresponding 2-cocycle $\sigma = \sigma_{\mathbf{q}}$ on $U_{\mathbf{q}}(\mathfrak{g}_D)$: this yields a map $\mathbf{q} \mapsto \sigma_{\mathbf{q}}$ which is injective, hence $q\text{-Mp}_{\mathbb{Q}}$ is in bijection with the subset $\widehat{\mathcal{Z}}_{\mathbb{Q}}^2$ of all $\sigma_{\mathbf{q}}$'s inside the set $\mathcal{Z}^2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ of all 2-cocycles. Composing this bijection with the bijection $\Psi \leftrightarrow \mathbf{q}$ we eventually get a third bijection between $\mathfrak{so}_n(\mathbb{Q})$ — encoding “rational twists” — and $\widehat{\mathcal{Z}}_{\mathbb{Q}}^2$ — encoding “rational 2-cocycles”. We shall shortly denote this bijection by $\Psi \leftrightarrow \sigma$: explicitly, it is described by

$$\mathfrak{so}_n(\mathbb{Q}) \xleftarrow{\quad} \widehat{\mathcal{Z}}_{\mathbb{Q}}^2 \quad , \quad \Psi \mapsto \sigma \quad \text{with} \quad \sigma(K_{\alpha_i}, K_{\alpha_j}) := q^{-(DA+A^T\Psi A)_{ij}} \quad (5.6)$$

On the other hand, we shall also consider yet another bijection, namely

$$\mathfrak{so}_n(\mathbb{Q}) \xleftarrow{\quad} \widetilde{\mathcal{Z}}_{\mathbb{Q}}^2 \quad , \quad \Psi \mapsto \sigma \quad , \quad \text{with} \quad \sigma(K_{\alpha_i}, K_{\alpha_j}) := q^{-(A^T\Psi A)_{ij}} \quad (5.7)$$

where $\widetilde{\mathcal{Z}}_{\mathbb{Q}}^2$ is nothing but the subset of all 2-cocycles in $\mathcal{Z}^2(U_{\mathbf{q}}(\mathfrak{g}_D), \mathbb{F})$ of the form given in (5.7); the inverse of this map is clearly given by

$$\widetilde{\mathcal{Z}}_{\mathbb{Q}}^2 \xleftarrow{\quad} \mathfrak{so}_n(\mathbb{Q}) \quad , \quad \sigma \mapsto \Psi_\sigma := -A^{-T}S_\sigma A^{-1} \quad (5.8)$$

where $S_\sigma := (s_{ij})_{i,j \in I}$ is uniquely defined by $\sigma(K_{\alpha_i}, K_{\alpha_j}) := q^{s_{ij}}$.

In the sequel, we shall denote this last bijective correspondence by $\Psi \longleftrightarrow \sigma$.

5.2. TwQUEA's vs. MpQUEA's: duality.

Roughly speaking — that is, up to technicalities such as dealing with finite-dimensional objects, or dealing with Hopf algebras in categories with a well-behaving notion of “dual Hopf algebra”, etc. — the two notions of “twist element” and of “2-cocycle” are, by definition, dual to each other (in Hopf-theoretical sense). As a

consequence the two procedures of “comultiplication twisting” and of “multiplication twisting” are dual to each other as well (see Proposition 2.4.4 for a formalization).

Beyond this, we can prove the following: when the Hopf algebras H and K are opposite (polynomial) quantum Borel subgroups in duality, the link between Ψ and σ is ruled precisely by the bijection $\Psi \longleftrightarrow \sigma$. In short, we can claim that

TwQEA's and MpQEA's of (opposite) Borel type are dual to each other and in this duality the correspondence twists \rightleftharpoons 2-cocycles is given by the bijection $\Psi \longleftrightarrow \sigma$. The precise statement is the following:

Theorem 5.2.1. *Let $U_{q,\Gamma_{\pm}}(\mathfrak{b}_{\pm})$ be opposite Borel quantum subgroups and let also $\eta : U_{q,\Gamma_+}(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} U_{q,\Gamma_-}(\mathfrak{b}_-) \longrightarrow \mathbb{k}_q$ be a skew-Hopf pairing as in §3.3.5. Given $\Psi \in \mathfrak{so}_n(\mathbb{Q})$, let $\Gamma'_{\pm} := (\text{id} + \psi_+)(\Gamma_{\pm}) - \psi_-(\Gamma_{\mp}) + Q_{\pm}$. For every $\sigma \in \mathcal{Z}^2(U_{q,\Gamma'_-}(\mathfrak{b}_-), \mathbb{k}_q)$, consider the corresponding twistings $U_{q,\Gamma'_+}^{\Psi}(\mathfrak{b}_+)$ and $(U_{q,\Gamma'_-}(\mathfrak{b}_-))_{\sigma}$. Then the extended linear map $\eta : U_{q,\Gamma'_+}^{\Psi}(\mathfrak{b}_+) \otimes_{\mathbb{k}_q} (U_{q,\Gamma'_-}(\mathfrak{b}_-))_{\sigma} \longrightarrow \mathbb{k}_q$ given in §3.3.5 (see also §4.3.3) with respect to $\Gamma_{\bullet} = \Gamma'_+ \times \Gamma'_-$ is again a skew-Hopf pairing with respect to the new, deformed Hopf structures if and only if $\Psi \longleftrightarrow \sigma$. In other words, the deformed coproduct $\Delta^{(\Psi)}$ on $U_{q,\Gamma'_+}(\mathfrak{b}_+)$ and the deformed product \cdot_{σ} on $U_{q,\Gamma'_-}(\mathfrak{b}_-)$ are dual to each other (via η) if and only if $\Psi \longleftrightarrow \sigma$.*

A symmetric, parallel result holds true when switching the roles of Ψ and σ from left to right and viceversa.

Proof. This is a sheer matter of computation. Indeed, let us consider for instance the element $\Delta^{(\Psi)}(E_j)$: by construction, if we consider the standard Q -grading on $U_{q,\Gamma'_-}(\mathfrak{b}_-)$ then for $Y \in U_{q,\Gamma'_-}(\mathfrak{b}_-)^{\times 2}$ we see that $\eta(E_j, Y) \neq 0$ can only occur with elements which actually belong to the \mathbb{k}_q -span of elements of the form $K_{\gamma_-} \cdot_{\sigma} F_j$ or $F_j \cdot_{\sigma} K_{\gamma_-}$ with $\gamma_- \in \Gamma'_-$; in fact, to simplify the notation we can assume $\gamma_- = \alpha_i^-$. Now write $\Delta^{(\Psi)}(E_j) = (E_j)_{(1)}^{\Psi} \otimes (E_j)_{(2)}^{\Psi} = E_j \otimes K_{+\psi_-(\alpha_j^-)} + K_{+(\text{id}+\psi_+)(\alpha_j^+)} \otimes E_j$. Then, a direct computation gives

$$\begin{aligned} (\eta \otimes \eta)(\Delta^{(\Psi)}(E_j), K_{\alpha_i^-} \otimes F_j) &= \eta((E_j)_{(1)}^{\Psi}, K_{\alpha_i^-}) \eta((E_j)_{(2)}^{\Psi}, F_j) = \\ &= \eta(E_j, K_{\alpha_i^-}) \eta(K_{+\psi_-(\alpha_j^-)}, F_j) + \eta(K_{+(\text{id}+\psi_+)(\alpha_j^+)}, K_{\alpha_i^-}) \eta(E_j, F_j) = \\ &= \eta(K_{+(\text{id}+\psi_+)(\alpha_j^+)}, K_{\alpha_i^-}) \eta(E_j, F_j) = q^{-(\text{id}+\psi_+)(\alpha_j^+), \alpha_i^-} \eta(E_j, F_j) \end{aligned}$$

On the other hand, for the deformed product \cdot_{σ} in $(U_{q,\Gamma'_-}(\mathfrak{b}_-))_{\sigma}$ we have

$$\begin{aligned} K_{\alpha_i^-} \cdot_{\sigma} F_j &= \sigma((K_{\alpha_i^-})_{(1)}, (F_j)_{(1)}) K_{\alpha_i^-} (F_j)_{(2)} \sigma^{-1}((K_{\alpha_i^-})_{(3)}, (F_j)_{(3)}) = \\ &= \sigma(K_{\alpha_i^-}, (F_j)_{(1)}) \sigma^{-1}(K_{\alpha_i^-}, (F_j)_{(3)}) K_{\alpha_i^-} (F_j)_{(2)} = \sigma^{-1}(K_{\alpha_i^-}, K_{\alpha_j^+}^{-1}) K_{\alpha_i^-} F_j \end{aligned}$$

so that

$$\begin{aligned} \eta(E_j, K_{\alpha_i^-} \cdot_{\sigma} F_j) &= \sigma^{-1}(K_{\alpha_i^-}, K_{\alpha_j^+}^{-1}) \eta(E_j, K_{\alpha_i^-} F_j) = \\ &= \sigma^{-1}(K_{\alpha_i^-}, K_{\alpha_j^+}^{-1}) \eta(K_{\alpha_j^+}, K_{\alpha_i^-}) \eta(E_j, F_j) = \\ &= q^{-d_i a_{ij}} \sigma(K_{\alpha_i^-}, K_{\alpha_j^+}) \eta(E_j, F_j) \end{aligned}$$

Comparing with the above, this means that we have

$$(\eta \otimes \eta)(\Delta^{(\Psi)}(E_j), K_{\alpha_i^-} \otimes F_j) = \eta(E_j, K_{\alpha_i^-} \cdot_{\sigma} F_j) \quad (5.9)$$

if and only if

$$\sigma(K_{\alpha_i^-}, K_{\alpha_j^+}) = q^{d_i a_{ij} - ((\text{id} + \psi_+)(\alpha_j^+), \alpha_i^-)} = q^{-(A^T \Psi A)_{ij}} \quad (5.10)$$

and the last condition means that $\Psi \longleftrightarrow \sigma$. Similar computations show that

$$\eta(\Delta^{(\Psi)}(E_j), F_j \otimes K_{\alpha_i^-}) = \eta(E_j, F_j \cdot_{\sigma} K_{\alpha_i^-}) \quad (5.11)$$

if and only if (5.10) holds, again, that is if and only if $\Psi \longleftrightarrow \sigma$.

Furthermore, notice that $\Delta^{(\Psi)}(K_{\gamma_+}) = \Delta(K_{\gamma_+})$ and $K_{\gamma'_-} \cdot_{\sigma} K_{\gamma''_-} = K_{\gamma'_-} K_{\gamma''_-}$ for all $\gamma'_{\pm}, \gamma''_{\pm} \in \Gamma_{\pm}$, so that we automatically have

$$\eta(\Delta^{(\Psi)}(K_{\gamma_+}), K_{\gamma'_-} \otimes K_{\gamma''_-}) = \eta(K_{\gamma_+}, K_{\gamma'_-} \cdot_{\sigma} K_{\gamma''_-}) \quad (5.12)$$

Now, conditions (5.9), (5.11) and (5.12) altogether are the conditions for $\Delta^{(\Psi)}$ and \cdot_{σ} to be *dual to each other* via η — so that the pairing η itself be a *Hopf skew-pairing* w.r.t. the new, deformed structures. The above proves that the sole necessary and sufficient condition for all this is (5.10), i.e., that $\Psi \longleftrightarrow \sigma$, as claimed.

The same argument proves also the last part of the claim, when the roles of Ψ and σ are interchanged. \square

5.3. TwQUEA's vs. MpQUEA's: correspondence.

We shall presently prove the following striking fact: *the classes of TwQUEA's and of MpQUEA's associated with "rational" data actually coincide*. More precisely, if we consider a (rational) antisymmetric twist Ψ and a q -rational multiparameter \mathbf{q} such that $\Psi \longleftrightarrow \mathbf{q}$, then any TwQUEA (over \mathfrak{b}_{\pm} , \mathfrak{g} or \mathfrak{g}_D) with twist Ψ and any MpQUEA (over \mathfrak{b}_{\pm} , \mathfrak{g} or \mathfrak{g}_D) with multiparameter \mathbf{q} are isomorphic to each other.

In other words, as each MpQUEA is a 2-cocycle deformation of the canonical quantum group, this result can be read as follows: *every comultiplication twisting by a (rational) twist of a canonical quantum group is a multiplication twisting by a (rational) 2-cocycle, and viceversa, with the correspondence twist \rightleftarrows 2-cocycle ruled by $\Psi \longleftrightarrow \sigma$.*

The precise statement of our main result — formulated here for "double" quantum groups — reads as follows:

Theorem 5.3.1. *Let $\Psi \in \mathfrak{so}_n(\mathbb{Q})$ and $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ be such that $\Psi \longleftrightarrow \mathbf{q}$. Let M_{\pm} be any lattice in $\mathbb{Q}Q$ containing Q , with $\{\mu_i^{\pm}\}_{i \in I}$ any \mathbb{Z} -basis of it, let M_{\pm}^{Ψ} be the sublattice of $\mathbb{Q}Q^{\times 2}$ with \mathbb{Z} -basis $\{\varpi_i^{\pm} := \mu_i^{\pm} + \psi_{\pm}(\mu_i^{\pm}) - \psi_{\mp}(\mu_i^{\mp})\}_{i \in I}$ and consider in $\mathbb{Q}Q^{\times 2}$ also the lattices $M_{\bullet} := M_+ \times M_-$ and $M_{*}^{\Psi} := M_+^{\Psi} + M_-^{\Psi}$.*

Let $U_{\mathbf{q}, M_{\bullet}}(\mathfrak{g}_D)$ be the MpQUEA associated with the lattice M_{\bullet} (inside $\mathbb{Q}Q^{\times 2}$), and $\hat{U}_{\mathbf{q}, M_{}^{\Psi}}^{\Psi}(\mathfrak{g}_D)$ be the TwQUEA associated with M_{*}^{Ψ} (cf. §3.4.8). Then there exists a Hopf algebra isomorphism $U_{\mathbf{q}, M_{\bullet}}(\mathfrak{g}_D) \cong \hat{U}_{\mathbf{q}, M_{*}^{\Psi}}^{\Psi}(\mathfrak{g}_D)$ given (for $i \in I$) by*

$$\begin{aligned} E_i &\mapsto q_i E_i^{\Psi} := q_i K_{-\psi_-(\alpha_i^-)} E_i, & K_{\mu_i^+} &\mapsto K_{+(\mu_i^+ + \psi_+(\mu_i^+) - \psi_-(\mu_i^-))} = K_{+\varpi_i^+} \\ L_{\mu_i^-} &\mapsto K_{-(\mu_i^- + \psi_-(\mu_i^-) - \psi_+(\mu_i^+))} = K_{-\varpi_i^-}, & F_i &\mapsto F_i^{\Psi} := K_{+\psi_+(\alpha_i^+)} F_i \end{aligned}$$

In other words, letting $\sigma = \sigma_{\mathbf{q}}$ be the (rational) 2-cocycle corresponding to \mathbf{q} so that $\Psi \longleftrightarrow \sigma$ (as in §5.1.2), we have a Hopf algebra isomorphism

$$(U_{\mathbf{q}, M_{\bullet}}(\mathfrak{g}_D))_{\sigma} \cong \hat{U}_{\mathbf{q}, M_{*}^{\Psi}}^{\Psi}(\mathfrak{g}_D)$$

given by the same formulas as above.

Proof. Define an algebra map $\Phi : U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D) \longrightarrow \hat{U}_{q, M_\bullet^\Psi}(\mathfrak{g}_D)$ on the generators of $U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D)$ as above. We only have to prove that such a Φ is well-defined, for then it is clearly surjective. Actually, Φ is well-defined indeed, since the defining relations of $U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D)$ — see Definition 4.1.1 and §4.3.2 — are all mapped to zero: this follows straightforward calculations, so we provide only some of them as guidelines.

It is clear that Φ “respects” the commutation relations (a) and (b) in Definition 4.1.1, so now we go for the other ones. As in §3.4.8, write

$$K_{i,+}^\Psi := K_{(\text{id} + \psi_+)(\alpha_i^+) - \psi_-(\alpha_i^-)} = K_{\alpha_i^+ + \zeta_{i,+}^+ - \zeta_{i,-}^-}$$

and

$$K_{i,-}^\Psi := K_{-(\text{id} + \psi_-)(\alpha_i^-) + \psi_+(\alpha_i^+)} = K_{-\alpha_i^- - \zeta_{i,-}^- + \zeta_{i,+}^+}$$

Recall that, through the correspondence $\Psi \longleftrightarrow \mathbf{q}$, we have

$$q_{ij} = q^{\vartheta'(\Psi)_{ij}} = q^{(\alpha_i + \psi_+(\alpha_i) - \psi_-(\alpha_i), \alpha_j)} \quad \forall i, j \in I.$$

Now, by definition we have $K_{+\varpi_i^+}^{\pm 1} K_{-\varpi_j^-}^{\pm 1} = K_{-\varpi_j^-}^{\pm 1} K_{+\varpi_i^+}^{\pm 1}$ for all $i, j \in I$. Moreover,

$$\begin{aligned} \Phi(K_i E_j K_i^{-1}) &= q_j (K_{i,+}^\Psi) E_j^\Psi (K_{i,+}^\Psi)^{-1} = \\ &= q_j K_{\alpha_i^+ + \psi_+(\alpha_i^+) - \psi_-(\alpha_i^-)} K_{-\psi_-(\alpha_j^-)} E_j K_{\alpha_i^+ + \psi_+(\alpha_i^+) - \psi_-(\alpha_i^-)}^{-1} = \\ &= q^{(\alpha_i + \psi_+(\alpha_i) - \psi_-(\alpha_i), \alpha_j)} q_j K_{-\psi_-(\alpha_j^-)} E_j = \\ &= q^{\vartheta'(\Psi)_{ij}} q_j E_j^\Psi = q_{ij} q_j E_j^\Psi = q_{ij} \Phi(E_j) \end{aligned}$$

$$\begin{aligned} \Phi(L_i E_j L_i^{-1}) &= q_j (K_{i,-}^\Psi) E_j^\Psi (K_{i,-}^\Psi)^{-1} = \\ &= q_j K_{-\alpha_i^- - \psi_-(\alpha_i^-) + \psi_+(\alpha_i^+)} K_{-\psi_-(\alpha_j^-)} E_j K_{-\alpha_i^- - \psi_-(\alpha_i^-) + \psi_+(\alpha_i^+)}^{-1} = \\ &= q^{(-\alpha_i - \psi_-(\alpha_i) + \psi_+(\alpha_i), \alpha_j)} q_j K_{-\psi_-(\alpha_j^-)} E_j \stackrel{\circledast}{=} \\ &\stackrel{\circledast}{=} q^{-(\alpha_j + \psi_+(\alpha_j) - \psi_-(\alpha_j), \alpha_i)} q_j K_{-\psi_-(\alpha_j^-)} E_j = \\ &= q^{-\vartheta'(\Psi)_{ji}} q_j E_j^\Psi = q_{ji}^{-1} q_j E_j^\Psi = q_{ji}^{-1} \Phi(E_j) \end{aligned}$$

where the equality $\stackrel{\circledast}{=}$ follows from Lemma 2.1.3. This proves that Φ “respects” the commutation relations (c) in Definition 4.1.1; similar computations prove the same for relations (d) as well. To check the relations (e), we first observe that

$$(\psi_+(\alpha_j), \alpha_i) = \sum_{k,\ell} \psi_{k,\ell} a_{\ell,j} a_{k,i} = (\alpha_j, \psi_-(\alpha_i)) \quad \text{for all } i, j \in I.$$

Then, for all $i, j \in I$ we have

$$\begin{aligned} [\Phi(E_i), \Phi(F_j)] &= q_i [E_i^\Psi, F_j^\Psi] = q_i [K_{-\psi_-(\alpha_i^-)} E_i, K_{+\psi_+(\alpha_j^+)} F_j] \\ &= q_i (K_{-\psi_-(\alpha_i^-)} E_i K_{+\psi_+(\alpha_j^+)} F_j - K_{+\psi_+(\alpha_j^+)} F_j K_{-\psi_-(\alpha_i^-)} E_i) \\ &= q_i K_{-\psi_-(\alpha_i^-)} K_{+\psi_+(\alpha_j^+)} (q^{-(\alpha_i, \psi_+(\alpha_j))} E_i F_j - q^{-(\alpha_j, \psi_-(\alpha_i))} F_j E_i) \end{aligned}$$

$$\begin{aligned}
&= q_i q^{-(\alpha_i, \psi_+(\alpha_j))} K_{-\psi_-(\alpha_i^-)} K_{+\psi_+(\alpha_j^+)} \delta_{ij} \frac{K_{\alpha_i^+} - K_{\alpha_i^-}^{-1}}{q_i - q_i^{-1}} \\
&= \delta_{ij} q_i q^{-(A^T \Psi A)_{ii}} \frac{K_{\alpha_i^+ + \psi_+(\alpha_i^+) - \psi_-(\alpha_i^-)} - K_{-\alpha_i^- + \psi_+(\alpha_i^+) - \psi_-(\alpha_i^-)}}{q_i - q_i^{-1}} \\
&= \delta_{ij} q_i^2 \frac{K_{i,+}^\Psi - K_{i,-}^\Psi}{q_i^2 - 1} = \Phi([E_i, F_j])
\end{aligned}$$

Finally, for the quantum Serre relations (f) and (g), we use that for all $m \in \mathbb{N}$ we have formal identities involving q -numbers and $q^{\frac{1}{2}}$ -numbers, namely

$$(m)_q = \frac{q^m - 1}{q - 1} = q^{\frac{m-1}{2}} \frac{q^{\frac{m}{2}} - q^{-\frac{m}{2}}}{q^{\frac{1}{2}} - q^{-\frac{1}{2}}} = q^{\frac{m-1}{2}} [m]_{q^{\frac{1}{2}}} , \quad \binom{m}{k}_q = q^{\frac{-k^2 + km}{2}} \left[\begin{matrix} m \\ k \end{matrix} \right]_{q^{\frac{1}{2}}}$$

and the identities $((\psi_+ - \psi_-)(\alpha_i), \alpha_j) = -(\alpha_i, (\psi_+ - \psi_-)(\alpha_j))$, $(\psi_+(\alpha_i), \alpha_j) = (\alpha_i, \psi_-(\alpha_j))$ and $(\psi_-(\alpha_i), \alpha_i) = 0$ for all $i, j \in I$.

It is not hard to see that Φ is also a *Hopf algebra map* too: for example, we have

$$\begin{aligned}
\Delta^{(\Psi)}(\Phi(E_i)) &= q_i \Delta^{(\Psi)}(E_i^\Psi) = q_i \Delta^{(\Psi)}(K_{-\psi_-(\alpha_i^-)} E_i) = q_i \Delta^{(\Psi)}(K_{-\psi_-(\alpha_i^-)}) \Delta^{(\Psi)}(E_i) \\
&= q_i K_{-\psi_-(\alpha_i^-)} E_i \otimes 1 + K_{\alpha_i^+ + \psi_+(\alpha_i^+) - \psi_-(\alpha_i^-)} \otimes q_i K_{-\psi_-(\alpha_i^-)} E_i = \\
&= q_i E_i^\Psi \otimes 1 + K_{i,+}^\Psi \otimes q_i E_i^\Psi = (\Phi \otimes \Phi)(\Delta(E_i))
\end{aligned}$$

and $\epsilon^{(\Psi)}(\Phi(E_i)) = q_i \epsilon^{(\Psi)}(E_i^\Psi) = \epsilon(q_i K_{-\psi_-(\alpha_i^-)} E_i) = 0 = \epsilon(E_i)$ for all $1 \leq i \leq n$.

Now define an algebra map $\Phi' : \hat{U}_{q, M_*^\Psi}(\mathfrak{g}_D) \longrightarrow U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D)$ by $\Phi'(E_i^\Psi) := q_i^{-1} E_i$, $\Phi'(F_i^\Psi) := F_i$, $\Phi'((K_{+\varpi_i^+})^{\pm 1}) := K_{\mu_i^+}^{\pm 1}$ and $\Phi'((K_{-\varpi_i^-})^{\pm 1}) := L_{\mu_i^-}^{\pm 1}$, for all $1 \leq i \leq n$. By means of calculations quite similar to the previous ones, one proves that such a Φ' is well-defined, and definitions clearly yield $\Phi' \circ \Phi = \text{id}_{U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D)}$ and $\Phi \circ \Phi' = \text{id}_{\hat{U}_{q, M_*^\Psi}(\mathfrak{g}_D)}$. So in the end $\hat{U}_{q, M_*^\Psi}(\mathfrak{g}_D) \cong U_{\mathbf{q}, M_\bullet}(\mathfrak{g}_D)$ as Hopf algebras. \square

Similar arguments as those used in the proof above lead to a simpler version of Theorem 5.3.1 for Borel MpQUEA's and Borel TwQUEA's:

Proposition 5.3.2. *Let $\Psi \in \mathfrak{so}_n(\mathbb{Q})$ and $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ be such that $\Psi \leftrightarrow \mathbf{q}$. Let M be any lattice in $\mathbb{Q}Q$ containing Q , with $\{\mu_i\}_{i \in I}$ any \mathbb{Z} -basis of it, and let $M_{\pm}^{(\Psi)}$ be the sublattice of $\mathbb{Q}Q \times \mathbb{Q}Q$ with \mathbb{Z} -basis $\{\varpi_i^{\pm} := \mu_i^{\pm} + \psi_{\pm}(\mu_i^{\pm}) - \psi_{\mp}(\mu_i^{\mp})\}_{i \in I}$.*

Let $U_{\mathbf{q}, M}(\mathfrak{b}_{\pm})$ be the (positive/negative) Borel MpQUEA associated with M , and $\hat{U}_{q, M_{\pm}^{(\Psi)}}(\mathfrak{b}_{\pm})$ be the (positive/negative) Borel TwQUEA associated with $M_{\pm}^{(\Psi)}$ (cf. §3.4.8). Then there exist Hopf algebra isomorphisms

$$U_{\mathbf{q}, M}(\mathfrak{b}_+) \cong \hat{U}_{q, M_+^{(\Psi)}}(\mathfrak{b}_+) , \quad U_{\mathbf{q}, M}(\mathfrak{b}_-) \cong \hat{U}_{q, M_-^{(\Psi)}}(\mathfrak{b}_-)$$

respectively given by

$$\begin{aligned}
E_i &\mapsto E_i^\Psi := K_{-\psi_-(\alpha_i^-)} E_i , & K_{\mu_i} &\mapsto K_{+\mu_i}^\Psi := K_{+(\mu_i^+ + \psi_+(\mu_i^+) - \psi_-(\mu_i^-))} = K_{+\varpi_i^+} \\
L_{\mu_i} &\mapsto K_{-\mu_i}^\Psi := K_{-(\mu_i^- + \psi_-(\mu_i^-) - \psi_+(\mu_i^+))} = K_{-\varpi_i^-} , & F_i &\mapsto F_i^\Psi := K_{+\psi_+(\alpha_i^+)} F_i
\end{aligned}$$

In other words, letting $\sigma = \sigma_{\mathbf{q}}$ be the (rational) 2-cocycle corresponding to \mathbf{q} as in §5.1.2, so that $\Psi \leftrightarrow \sigma$, we have Hopf algebra isomorphisms

$$(U_{q,M}(\mathfrak{b}_+))_{\sigma} \cong \hat{U}_{q,M_+}^{\Psi}(\mathfrak{b}_+) \quad , \quad (U_{q,M}(\mathfrak{b}_-))_{\sigma} \cong \hat{U}_{q,M_-}^{\Psi}(\mathfrak{b}_-)$$

given by the same formulas as above. In particular, when $M \supseteq Q_{\pm} + \psi_{\pm}(Q_{\pm}) + \psi_{\mp}(Q_{\mp})$ we have $\hat{U}_{q,M_{\pm}}^{\Psi}(\mathfrak{b}_{\pm}) = U_{q,M}^{\Psi}(\mathfrak{b}_{\pm})$, hence the isomorphisms above read

$$(U_{q,M}(\mathfrak{b}_{\pm}))_{\sigma} \cong U_{\mathbf{q},M}(\mathfrak{b}_{\pm}) \cong U_{q,M}^{\Psi}(\mathfrak{b}_{\pm})$$

A similar claim holds true as well for twisted quantum Borel subalgebras and multiparameter quantum Borel subalgebras inside the MpQUEA's $U_{q,M}^{\Psi}(\mathfrak{g}_D)$. \square

Finally, we shall presently see that any TwQUEA of type $U_{q,\Gamma}^{\Psi}(\mathfrak{g})$ can be seen as a (still to define) “ q -rational MpQUEA over \mathfrak{g} ”. We begin by defining the latter.

Definition 5.3.3. Let $\Psi \in \mathfrak{so}_n(\mathbb{Q})$ and $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ be such that $\Psi \leftrightarrow \mathbf{q}$, and let ψ_+ be the map associated with Ψ (cf §2.1.2). Let M_{\pm} be lattices in $\mathbb{Q}Q$ containing Q such that $\psi_+(M_{\pm}) \subseteq M_{\pm}$; denote by $\{\mu_i^{\pm}\}_{i \in I}$ any \mathbb{Z} -basis of it. Let \mathcal{I}_{Ψ} be the two-sided ideal of $U_{\mathbf{q},M_{\bullet}}(\mathfrak{g}_D)$ generated by the elements

$$K_{\mu_i^+} L_{\mu_i^-} - K_{2\psi_+(\mu_i^+)} L_{2\psi_-(\mu_i^-)} \quad (i \in I) .$$

Then we denote by $U_{\mathbf{q},M_{\bullet}}(\mathfrak{g})$ — that we loosely call “MpQUEA over \mathfrak{g} ” — the quotient Hopf algebra

$$U_{\mathbf{q},M_{\bullet}}(\mathfrak{g}) := U_{\mathbf{q},M_{\bullet}}(\mathfrak{g}_D) / \mathcal{I}_{\Psi}$$

Remark 5.3.4. If M is any lattice in $\mathbb{Q}Q$ containing Q and such that $\psi_+(M) \subseteq M$, then for $M_+ := M =: M_-$ we write $M_{\bullet} := M_+ \times M_-$. In particular, $M = M_+ + M_- = M^{\Psi} = M + \psi_+(M) + \psi_-(M)$.

Now, under the assumption $\psi_+(M) \subseteq M$ for the lattice in $\mathbb{Q}Q$, we have that $\hat{U}_{q,M}^{\Psi}(\mathfrak{g}) = U_{q,M}^{\Psi}(\mathfrak{g})$. In fact, $\hat{U}_{q,M}^{\Psi}(\mathfrak{g}_D)$ is the Hopf subalgebra of $U_{q,M}^{\Psi}(\mathfrak{g}_D)$ generated by the elements $E_i^{\Psi} := K_{-\psi_-(\alpha_i^-)} E_i$, $F_i^{\Psi} := K_{+\psi_+(\alpha_i^+)} F_i$, $K_{\varpi_i^+} := K_{\mu_i^+ + \psi_+(\mu_i^+) - \psi_-(\mu_i^-)}$ and $K_{\varpi_i^-} := K_{\mu_i^- + \psi_-(\mu_i^-) - \psi_+(\mu_i^+)}$ for all $i \in I$. Besides, by Remark 3.4.6, $\hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ is the image of the Hopf algebra epimorphism $\hat{\pi}_{\mathfrak{g}}^{\Psi} : \hat{U}_{q,M}^{\Psi}(\mathfrak{g}_D) \rightarrow \hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ given by $\hat{\pi}_{\mathfrak{g}}^{\Psi}(E_i^{\Psi}) := E_i^{\Psi}$, $\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^+}) := K_{\mu_i + 2\psi_+(\mu_i)}$, $\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^-}) := K_{\mu_i - 2\psi_+(\mu_i)}$ and $\hat{\pi}_{\mathfrak{g}}^{\Psi}(F_i^{\Psi}) := F_i^{\Psi}$ for all $i \in I$; therefore $K_{\mu_i}^2 = \hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^+} K_{\varpi_i^-}) \in \hat{U}_{q,M}^{\Psi}(\mathfrak{g})$. This implies that $K_{2\psi_+(\mu_i)} \in \hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ and consequently $K_{\mu_i} \in \hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ for all $i \in I$; hence, the toral part (the group of group-like elements) of both Hopf algebras coincide. Since the other generators of $\hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ differ only by a group-like element, the two Hopf algebras $\hat{U}_{q,M}^{\Psi}(\mathfrak{g})$ and $U_{q,M}^{\Psi}(\mathfrak{g})$ must coincide.

We are now ready for the result comparing TwQUEA's and MpQUEA's “over \mathfrak{g} ”:

Theorem 5.3.5. Let $\Psi \in \mathfrak{so}_n(\mathbb{Q})$ and $\mathbf{q} \in q\text{-Mp}_{\mathbb{Q}}$ be such that $\Psi \leftrightarrow \mathbf{q}$. Let also M be any lattice in $\mathbb{Q}Q$ containing Q and such that $\psi_+(M) \subseteq M$. Then there exists a Hopf algebra isomorphism

$$U_{\mathbf{q},M_{\bullet}}(\mathfrak{g}) \cong U_{q,M}^{\Psi}(\mathfrak{g})$$

given by the formulas $K_{\mu_i} \longleftrightarrow K_{+(\text{id}+\psi_+-\psi_-)(\mu_i)}$, $L_{\mu_i} \longleftrightarrow K_{-(\text{id}+\psi_--\psi_+)(\mu_i)}$, $E_i \longleftrightarrow q_i K_{-\psi_-(\alpha_i)} E_i$ and $F_i \longleftrightarrow K_{+\psi_+(\alpha_i)} F_i$, for all $i \in I$.

In other words, letting $\sigma = \sigma_{\mathbf{q}}$ be the (rational) 2-cocycle corresponding to \mathbf{q} as in §5.1.2, so that $\Psi \rightsquigarrow \sigma$, we have a Hopf algebra isomorphism

$$(U_{\mathbf{q},M\bullet}(\mathfrak{g}))_{\sigma} \cong U_{\mathbf{q},M}^{\Psi}(\mathfrak{g})$$

given by the same formulas as above.

Remark: as we assume Ψ to be antisymmetric, which is not restrictive, we can re-cast the formulas above as $K_{\mu_i} \longleftrightarrow K_{+(\text{id}+2\psi_+)(\mu_i)}$, $L_{\mu_i} \longleftrightarrow K_{-(\text{id}+2\psi_-)(\mu_i)}$, $E_i \longleftrightarrow q_i K_{+\psi_+(\alpha_i)} E_i$ and $F_i \longleftrightarrow K_{-\psi_-(\alpha_i)} F_i$, for all $i \in I$.

Proof. Let $\{\mu_i\}_{i \in I}$ be any \mathbb{Z} -basis of M and denote by $\{\mu_i^{\pm}\}_{i \in I}$ the corresponding \mathbb{Z} -basis of M_{\pm} . From Remark 3.4.6 we know that $\hat{U}_{\mathbf{q},M}^{\Psi}(\mathfrak{g})$ is the image of the Hopf algebra epimorphism $\hat{\pi}_{\mathfrak{g}}^{\Psi} : \hat{U}_{\mathbf{q},M}^{\Psi}(\mathfrak{g}_D) \longrightarrow \hat{U}_{\mathbf{q},M}^{\Psi}(\mathfrak{g})$, where the latter Hopf algebra equals $U_{\mathbf{q},M}^{\Psi}(\mathfrak{g})$ by Remark 5.3.4.

On the other hand, thanks to Theorem 5.3.1 we have a Hopf algebra isomorphism $\Phi : U_{\mathbf{q},M\bullet}(\mathfrak{g}_D) \xrightarrow{\sim} \hat{U}_{\mathbf{q},M}^{\Psi}(\mathfrak{g}_D)$. Hence, under the condition that $(\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(\mathcal{I}_{\Psi}) = 0$, there exists a surjective Hopf algebra morphism $\varphi : U_{\mathbf{q},M\bullet}(\mathfrak{g}) \longrightarrow U_{\mathbf{q},M}^{\Psi}(\mathfrak{g})$ defined on the generators by

$$\begin{aligned} \varphi(K_{\mu_i}) &:= K_{+(\text{id}+\psi_+-\psi_-)(\mu_i)} = K_{+(\text{id}+2\psi_+)(\mu_i)} , & \varphi(E_i) &:= q_i K_{-\psi_-(\alpha_i)} E_i \\ \varphi(L_{\mu_i}) &:= K_{-(\text{id}+\psi_--\psi_+)(\mu_i)} = K_{-(\text{id}+2\psi_-)(\mu_i)} , & \varphi(F_i) &:= K_{+\psi_+(\alpha_i)} F_i \end{aligned}$$

for all $i \in I$, where we denote by K_{μ_i} , L_{μ_i} , E_i and F_i the generators of $U_{\mathbf{q},M\bullet}(\mathfrak{g})$.

Now, the fact that $(\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(\mathcal{I}_{\Psi}) = 0$ follows by direct computation. In fact, since by assumption M contains Q and $\psi_+(M)$, and Ψ is antisymmetric, we have that $(\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(K_{\varpi_i^+}) = \hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{+\varpi_i^+}) = K_{\mu_i+2\psi_+(\mu_i)}$ and $(\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(L_{\mu_i^-}) = \hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{-\varpi_i^-}) = K_{-(\text{id}+2\psi_-)(\mu_i)}$. Therefore, we have also

$$(\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(K_{\mu_i^+} L_{\mu_i^-}) = (\hat{\pi}_{\mathfrak{g}}^{\Psi} \circ \Phi)(K_{2\psi_+(\mu_i^+)} L_{2\psi_-(\mu_i^-)})$$

Conversely, let $\hat{\varphi} : U_{\mathbf{q},M}^{\Psi}(\mathfrak{g}) \longrightarrow U_{\mathbf{q},M\bullet}(\mathfrak{g})$ be the algebra morphism given by

$$\hat{\varphi}(K_{\mu_i}) := K_{\mu_i^+-\psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} , \quad \hat{\varphi}(E_i^{\Psi}) := q_i^{-1} E_i , \quad \hat{\varphi}(F_i^{\Psi}) := F_i$$

for all $i \in I$, where $E_i^{\Psi} := q_i K_{-\psi_-(\alpha_i)} E_i$ and $F_i^{\Psi} := K_{+\psi_+(\alpha_i)} F_i$ are the images in $U_{\mathbf{q},M\bullet}(\mathfrak{g})$ of the same name elements in $U_{\mathbf{q},M}^{\Psi}(\mathfrak{g}_D)$. With this definition we have

$$\begin{aligned} \hat{\varphi}(\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^+})) &= \hat{\varphi}(K_{(\text{id}+2\psi_+)(\mu_i)}) = K_{+\mu_i^+} \\ \hat{\varphi}(\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^-})) &= \hat{\varphi}(K_{(\text{id}+2\psi_-)(\mu_i)}) = L_{-\mu_i^-} \end{aligned} \tag{5.13}$$

Indeed, since $K_{\mu_j^2} = \hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_j^+} K_{\varpi_j^-})$ for all $j \in I$, setting $\psi_+(\mu_i) := \sum_{j \in I} m_{ji} \mu_j$ yields $\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^+}) = K_{\mu_i+2\psi_+(\mu_i)} = K_{\mu_i} \prod_{j \in I} (K_{\mu_j^2})^{m_{ji}}$. Therefore

$$\begin{aligned} \hat{\varphi}(\hat{\pi}_{\mathfrak{g}}^{\Psi}(K_{\varpi_i^+})) &= \hat{\varphi}\left(K_{\mu_i} \prod_{j \in I} (K_{\mu_j^2})^{m_{ji}}\right) = \\ &= K_{\mu_i^+-\psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} \cdot \prod_{j \in I} K_{\mu_j^+-\psi_+(\mu_j^+)}^{2m_{ji}} \cdot \prod_{j \in I} L_{\psi_+(\mu_j^-)}^{2m_{ji}} = \end{aligned}$$

$$\begin{aligned}
 &= K_{\mu_i^+ - \psi_+(\mu_i^+) + 2\psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} \cdot \prod_{j \in I} K_{-\psi_+(\mu_j^+)}^{2m_{ji}} \cdot \prod_{j \in I} L_{\psi_+(\mu_j^-)}^{2m_{ji}} = \\
 &= K_{\mu_i^+ + \psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} \cdot \prod_{j \in I} (K_{-2\psi_+(\mu_j^+)} L_{-2\psi_-(\mu_j^-)})^{m_{ji}} = \\
 &= K_{\mu_i^+ + \psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} \cdot \prod_{j \in I} (K_{\mu_j^+} L_{\mu_j^-})^{-m_{ji}} = \\
 &= K_{\mu_i^+ + \psi_+(\mu_i^+)} L_{\psi_+(\mu_i^-)} K_{-\psi_+(\mu_i^+)} L_{-\psi_+(\mu_i^-)} = K_{\mu_i^+}
 \end{aligned}$$

Similarly, one may check the equalities in (5.13). In particular, $\hat{\varphi}$ is surjective.

Since $Q \subseteq M$, there exist $c_{ji} \in \mathbb{Z}$ such that $\alpha_i = \sum_{j \in I} c_{ji} \mu_j$ for all $i \in I$. Then $K_{\alpha_i^+} = \prod_{j \in J} K_{\mu_j^+}^{c_{ji}}$ and $K_{i,+}^\Psi = \prod_{j \in I} K_{\varpi_j^+}^{c_{ji}}$. This implies that $\tilde{\varphi}(\hat{\pi}_{\mathfrak{g}}^\Psi(K_{i,+}^\Psi)) = \hat{\varphi}(\hat{\pi}_{\mathfrak{g}}^\Psi(\prod_{j \in I} K_{\varpi_j^+}^{c_{ji}})) = \prod_{j \in I} K_{\mu_j^+}^{c_{ji}} = K_{\alpha_i^+}$. Similarly, one sees that $\hat{\varphi}(\hat{\pi}_{\mathfrak{g}}^\Psi(K_{i,-}^\Psi)) = L_{-\alpha_i^-}$, and from this it is easy to verify that $\hat{\varphi}$ is indeed a Hopf algebra morphism.

Finally, since $\hat{\varphi} \circ \varphi$ and $\varphi \circ \hat{\varphi}$ are the identity on the generators, we conclude that $U_{\mathfrak{q}, M_\bullet}(\mathfrak{g}) \cong U_{\mathfrak{q}, M}^\Psi(\mathfrak{g})$, via the formulas given in the claim, q.e.d. \square

Remark 5.3.6. A final remark is in order. In the previous results — Proposition 5.3.2, Theorem 5.3.1 and Theorem 5.3.5 — we took from scratch $\Psi \in \mathfrak{so}_n(\mathbb{Q})$, i.e., our “twisting datum” Ψ was antisymmetric. However, we can also start with any twisting matrix $\Psi \in M_n(\mathbb{Q})$: then *those results read the same as soon as we replace $\Psi \leftrightarrow \mathfrak{q}$ with $(\chi \circ \vartheta)(\Psi) = \mathfrak{q}$ (notation of §5.1.2)*. Notice then that one has

$$\mathfrak{q} := (\chi \circ \vartheta)(\Psi) = (\chi \circ \vartheta)(\Psi_a)$$

where $\Psi_a := 2^{-1}(\Psi - \Psi^T)$ is the antisymmetric part of Ψ . Eventually, the outcome of this discussion, in short, is the following:

The (polynomial) TwQUEA’s built out of any matrix $\Psi \in M_n(\mathbb{Q})$ are exactly the same as those obtained just from antisymmetric matrices $\Psi \in \mathfrak{so}_n(\mathbb{Q})$.

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